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# Towards a Global Evolutionary Picture of Protoplanetary Disks

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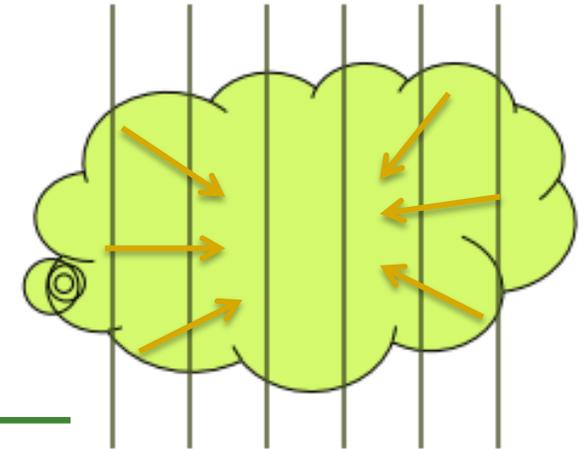
# Prelude: initial condition

- Protostellar cores threaded by strong B field

(e.g. Crutcher, 12, Hull+13)

- ~100 AU disk formation in class 0 phase

(e.g., Tobin+12,15,Murillo+13,Ohashi+14,Harsono+14,Yen+15)

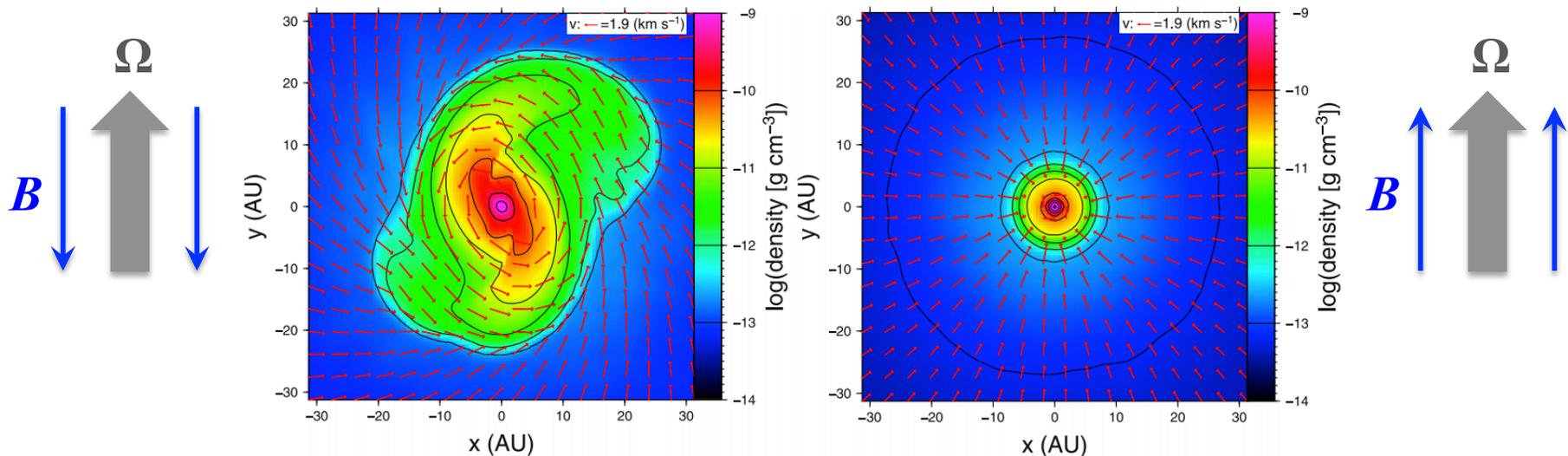


- Magnetic breaking catastrophe? (Mellon & Li 08)

Problem largely solved by turbulent collapse/misaligned field/non-ideal MHD

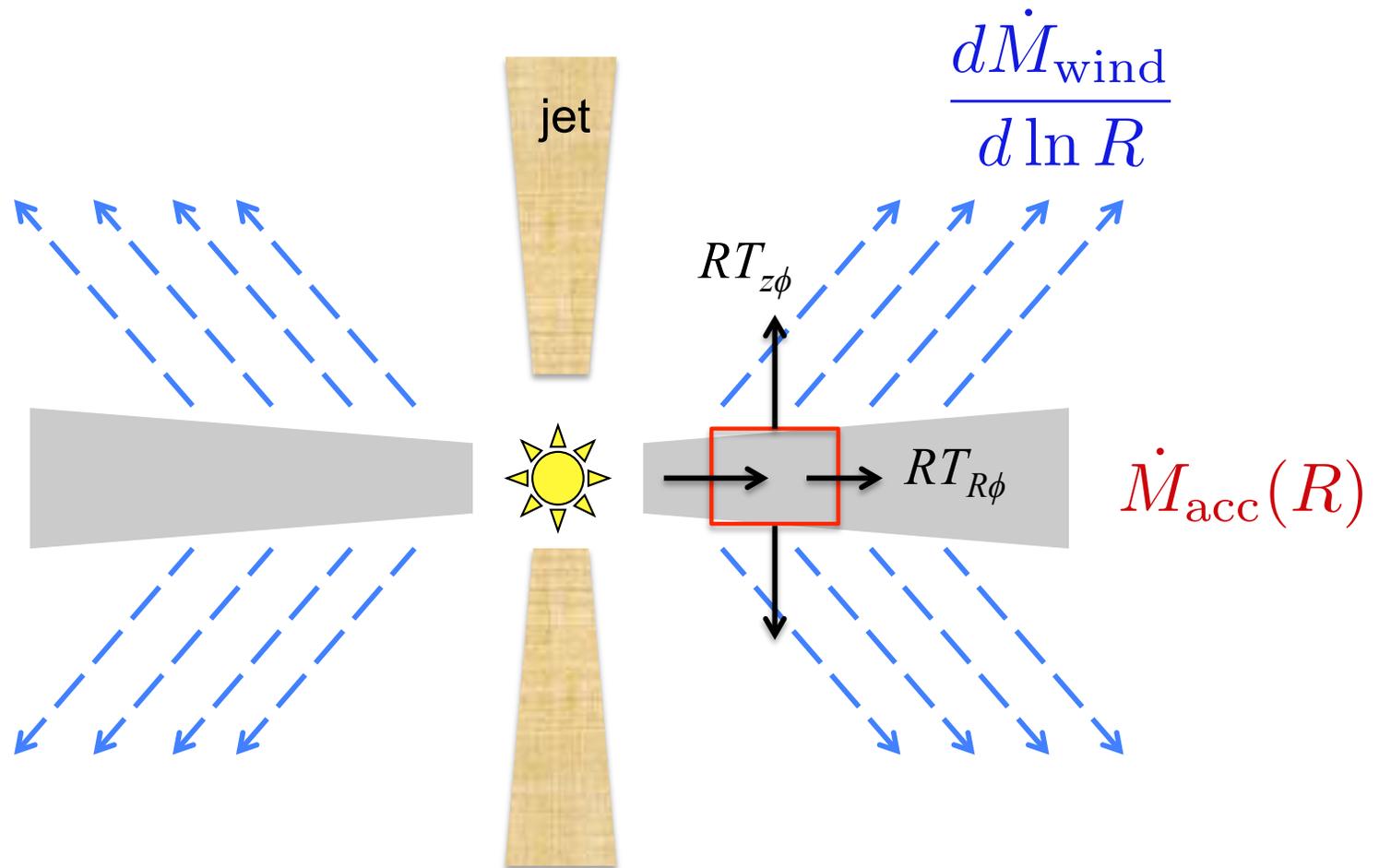
(e.g. Machida+10,Santos-Lima+11, Seifried+12,13, Joos+12,13, Li+13, Tomida+13,15)

- Hall-effect induced bimodality (Tsukamoto+15):



# What drives global disk evolution

- Angular momentum transport (radially/vertically)
- Mass loss (via disk wind)



# Radial transport of angular momentum

$$T_{R\phi} = \overline{\rho v_R v_\phi} - \frac{\overline{B_R B_\phi}}{4\pi} + \frac{\overline{g_R g_\phi}}{4\pi G} \equiv \alpha P$$

hydro turbulence
Maxwell stress
self-gravity

Prescribe radial profile of  $\alpha \Rightarrow$  disk evolution is straightforward.

Hydro processes	GI, VSI (GSF), baroclinic/ convective, zombie vortex... spiral shocks	see <a href="#">G. Lesur's</a> talk for more details about these processes  see <a href="#">E. Borobyov's</a> talk for early disk evolution with GI
MHD processes	turbulent stress (due to the MRI) laminar stress (i.e., spiral field)	

$\alpha$ -disk model OK only when radial transport of AM dominates

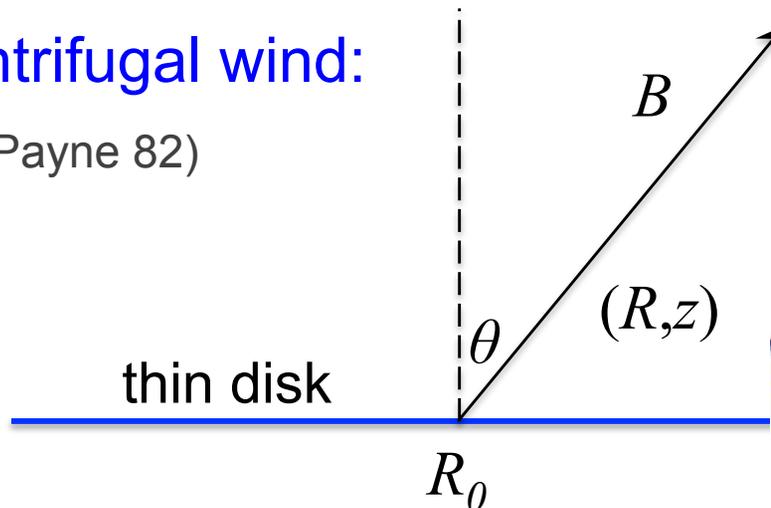
# Vertical transport: magnetized disk wind

$$T_{z\phi} = \cancel{\rho v_z \Delta v_\phi} - \frac{B_z B_\phi}{4\pi}$$

Requires external B field threading disk.

## Magneto-centrifugal wind:

(Blandford & Payne 82)

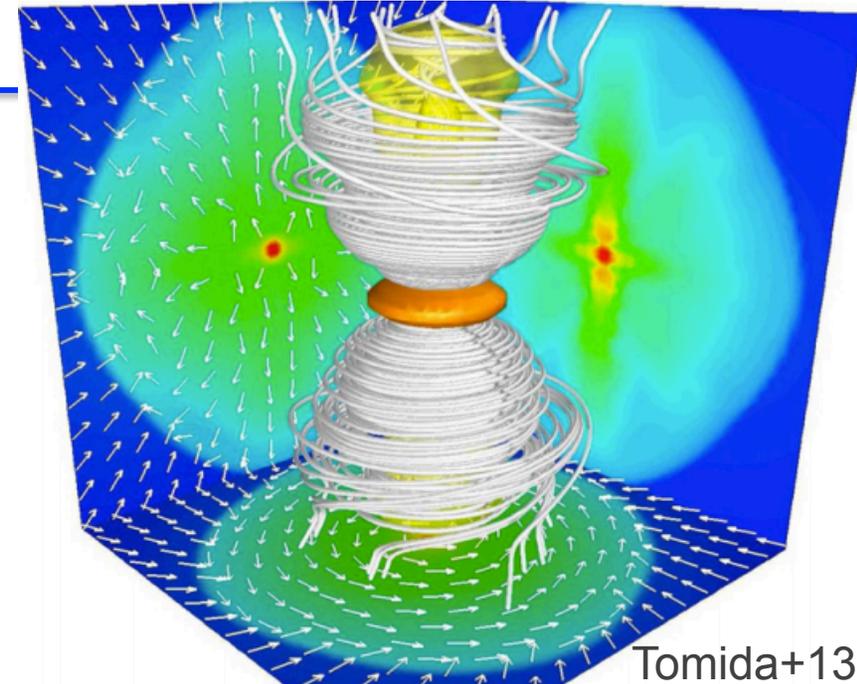


strong poloidal field enforces corotation.

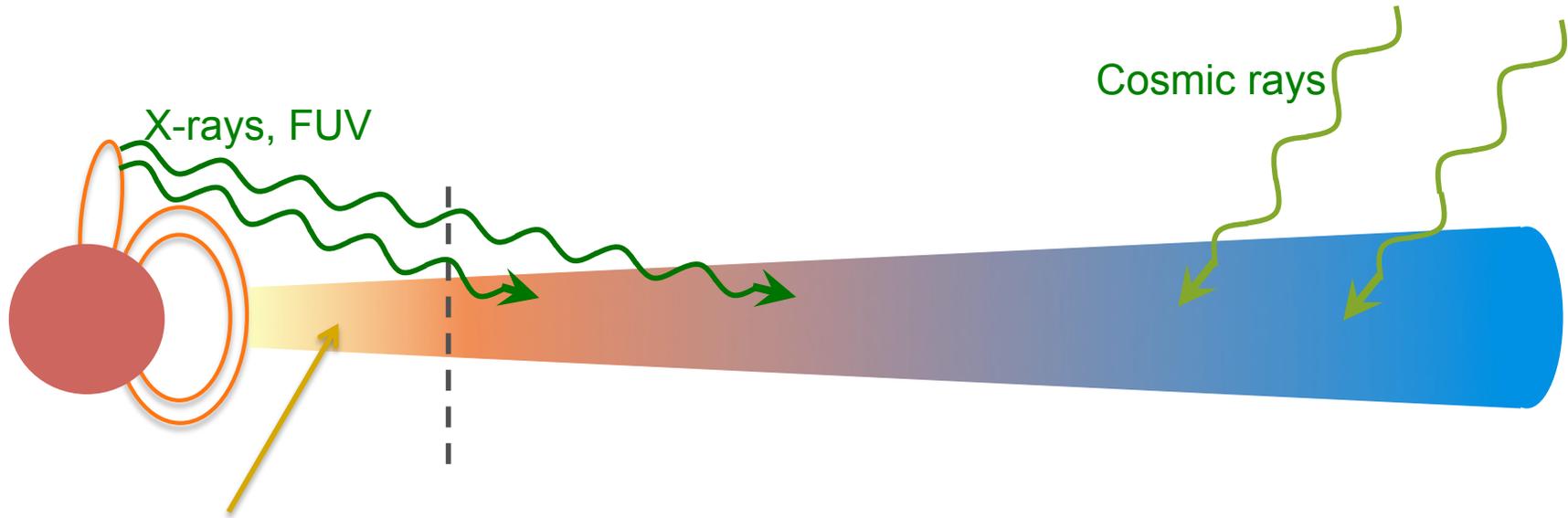
## “Magnetic tower”:

(e.g. Lynden-Bell 03)

Weak poloidal field shear-amplified & twisted, outflow driven by magnetic pressure gradient (of toroidal field).



# Disk microphysics: source of ionization



**Thermal ionization zone:** ionization of Na/K above  $\sim 800$  K (Desch & Turner 15)

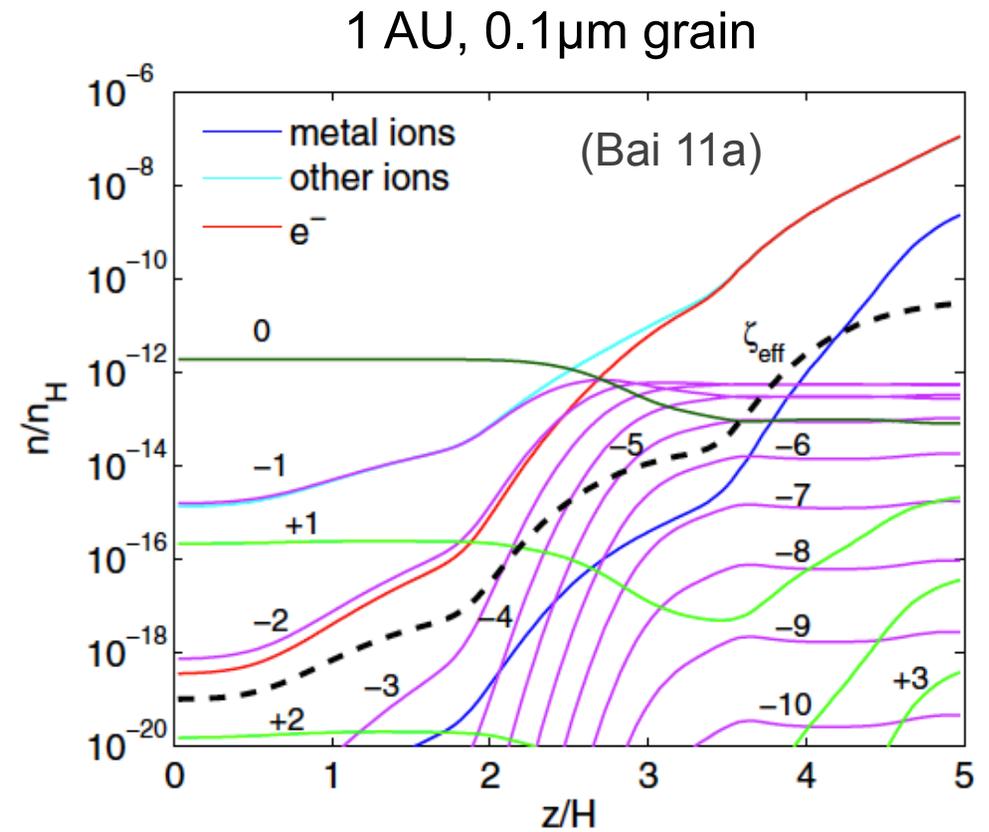
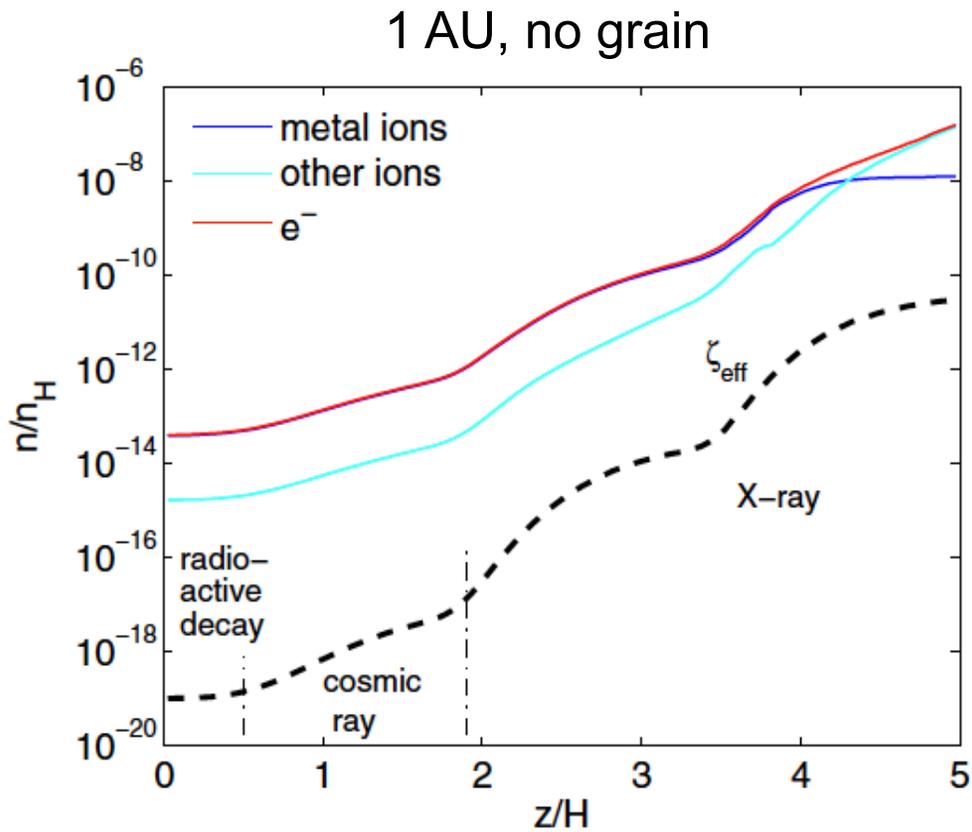
**Cosmic-ray ionization:** large penetration depth, can be affected by wind  
(Umebayashi & Nakano 81, Cleeves+13)

**Stellar X-ray ionization:** modest penetration depth (Igea & Glassgold 99,  
Ercolano & Glassgold 13)

**Stellar Far UV ionization:** fully ionize C, S with small penetration depth  
(Perez-Becker & Chiang 11)

**Radioactive decay:** time-dependent,  $\sim 10^{-19} \text{ s}^{-1}$  (Umebayashi & Nakano 09).

# Importance of grain size and abundance



- Ionization level strongly depends on grain size/abundance as long as grain abundance  $>$  ionization fraction.
- Grains can become the dominant charge carrier.

# Disk microphysics: non-ideal MHD effects

Generalized Ohm's law is in tensor form:  $\mathbf{J} = \boldsymbol{\sigma}(\mathbf{B}, \mathbf{E})\mathbf{E}$

Induction equation (grain-free):

$$\frac{\partial \mathbf{B}}{\partial t} = \nabla \times (\mathbf{v} \times \mathbf{B}) - \nabla \times \left[ \frac{4\pi\eta}{c} \mathbf{J} + \frac{\mathbf{J} \times \mathbf{B}}{en_e} - \frac{(\mathbf{J} \times \mathbf{B}) \times \mathbf{B}}{c\gamma\rho\rho_i} \right]$$

**inductive**

**Ohmic**

**Hall**

**AD**

midplane region  
of the inner disk

$$\sim \frac{n}{n_e}$$

$$\sim \frac{n}{n_e} \frac{B}{\rho}$$

$$\sim \frac{n}{n_e} \frac{B^2}{\rho^2}$$

Intermediate layer @ inner disk

Midplane towards outer disk

inner disk surface  
and outer disk

Further complications:

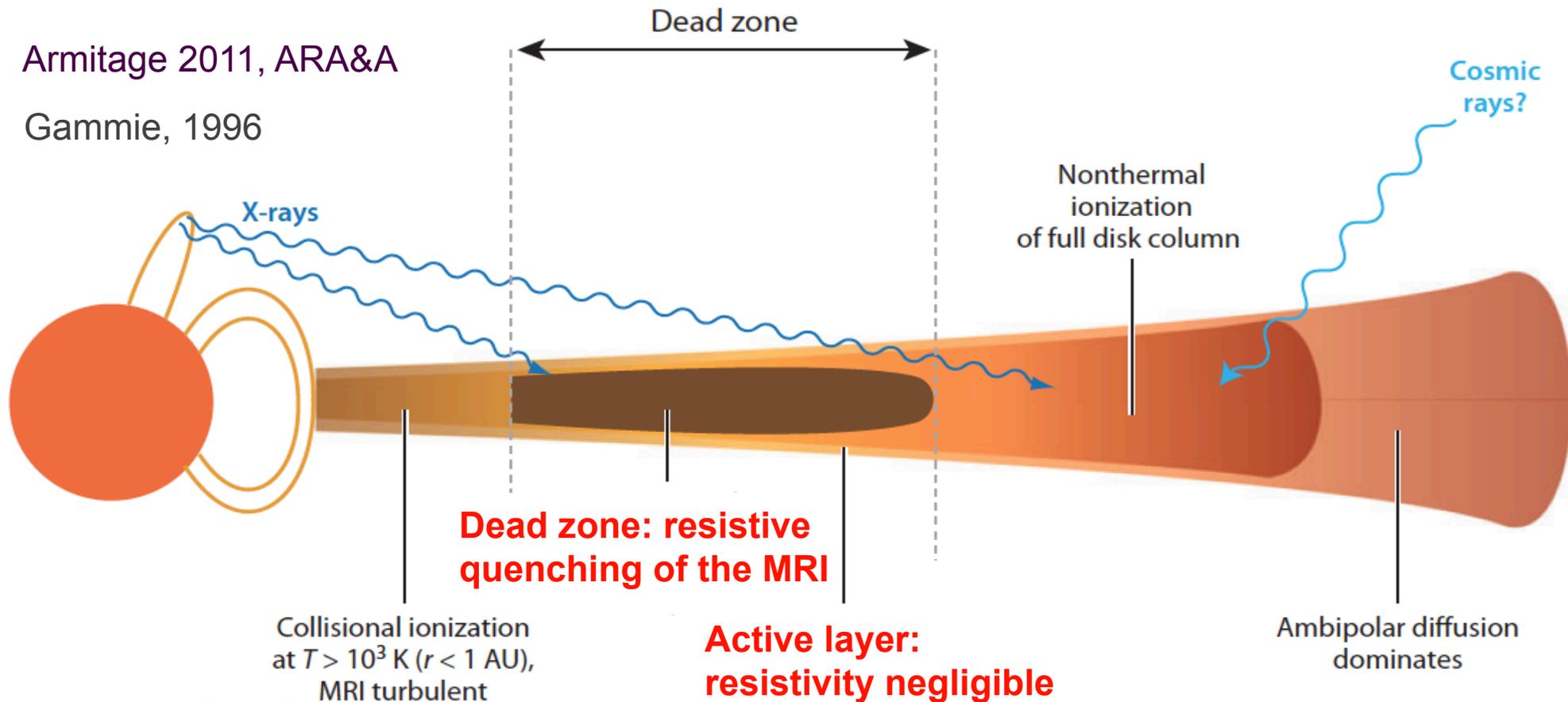
Non-linear Ohm's law (Okuzumi+15)

Complex dependence of diffusivities on  $\mathbf{B}$  (Wardle 07, Bai 11b, Xu & Bai, in prep)

# Ohmic resistivity only

Armitage 2011, ARA&A

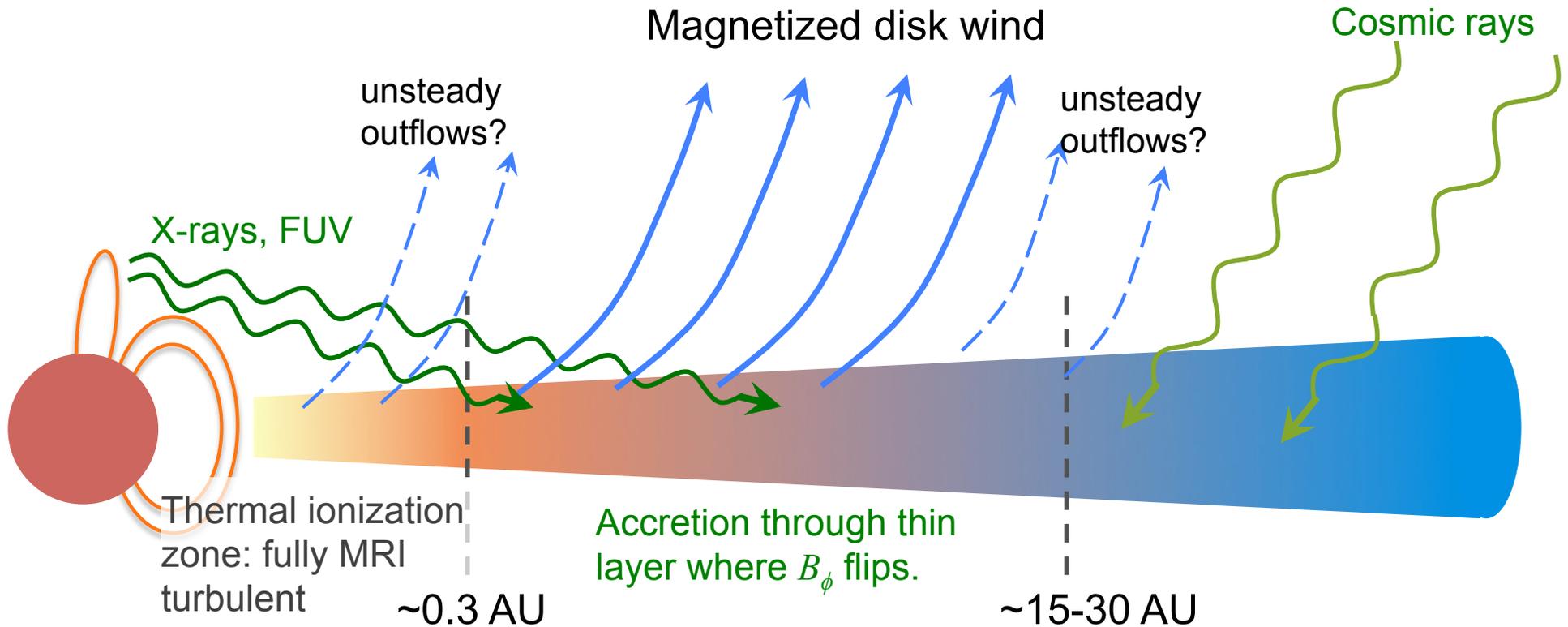
Gammie, 1996



- Semi-analytical studies already indicated that MRI is insufficient to drive rapid accretion when including the effect of ambipolar diffusion (Bai & Stone, 2011, Bai, 2011a,b, Perez-Becker & Chiang, 2011a,b).

# Ohmic + ambipolar diffusion

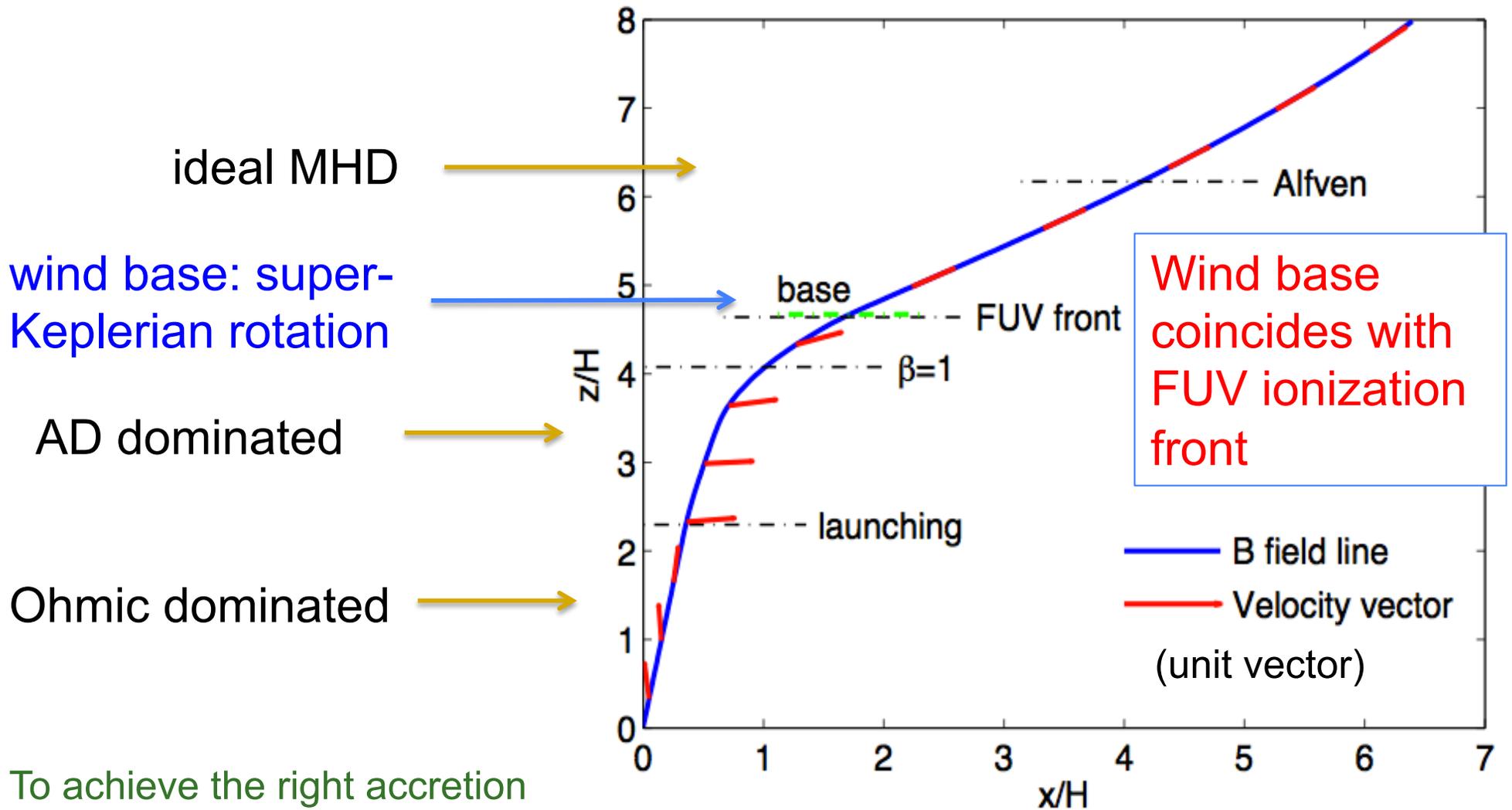
(Bai & Stone 13b, Bai, 2013, Simon, Bai+ 13a,b, Gressel+15)



MRI suppressed by Ohmic+AD, disk largely laminar.

MRI damped by AD

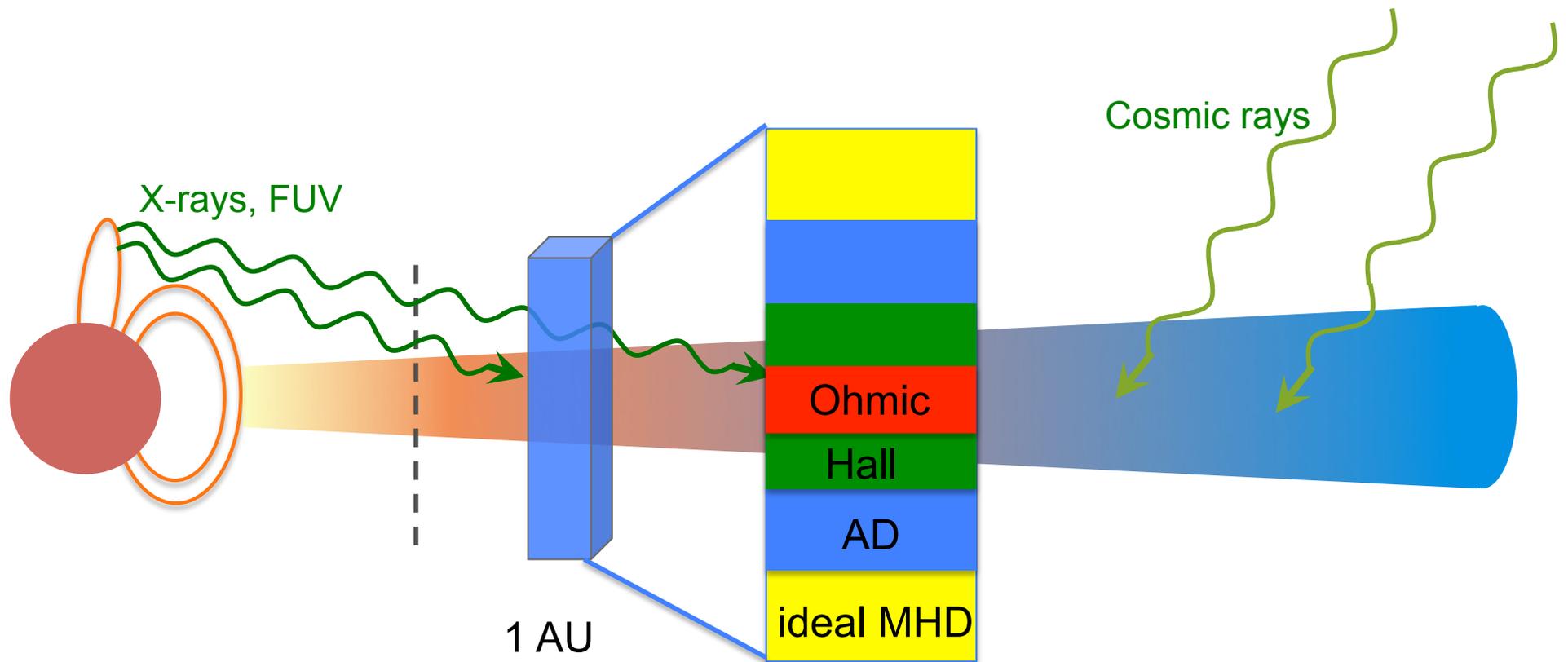
# Wind launching process (inner PPDs)



To achieve the right accretion rate, only need a weak  $B_z$ , but it  $\sim$ equi-partition at wind base.

(Bai & Stone, 2013b)

# Adding the Hall effect: inner Disk



All three non-ideal MHD effects are important.

Unstratified simulations: Kunz & Lesur 13 (local), O’Keeffe & Downes 14 (global)

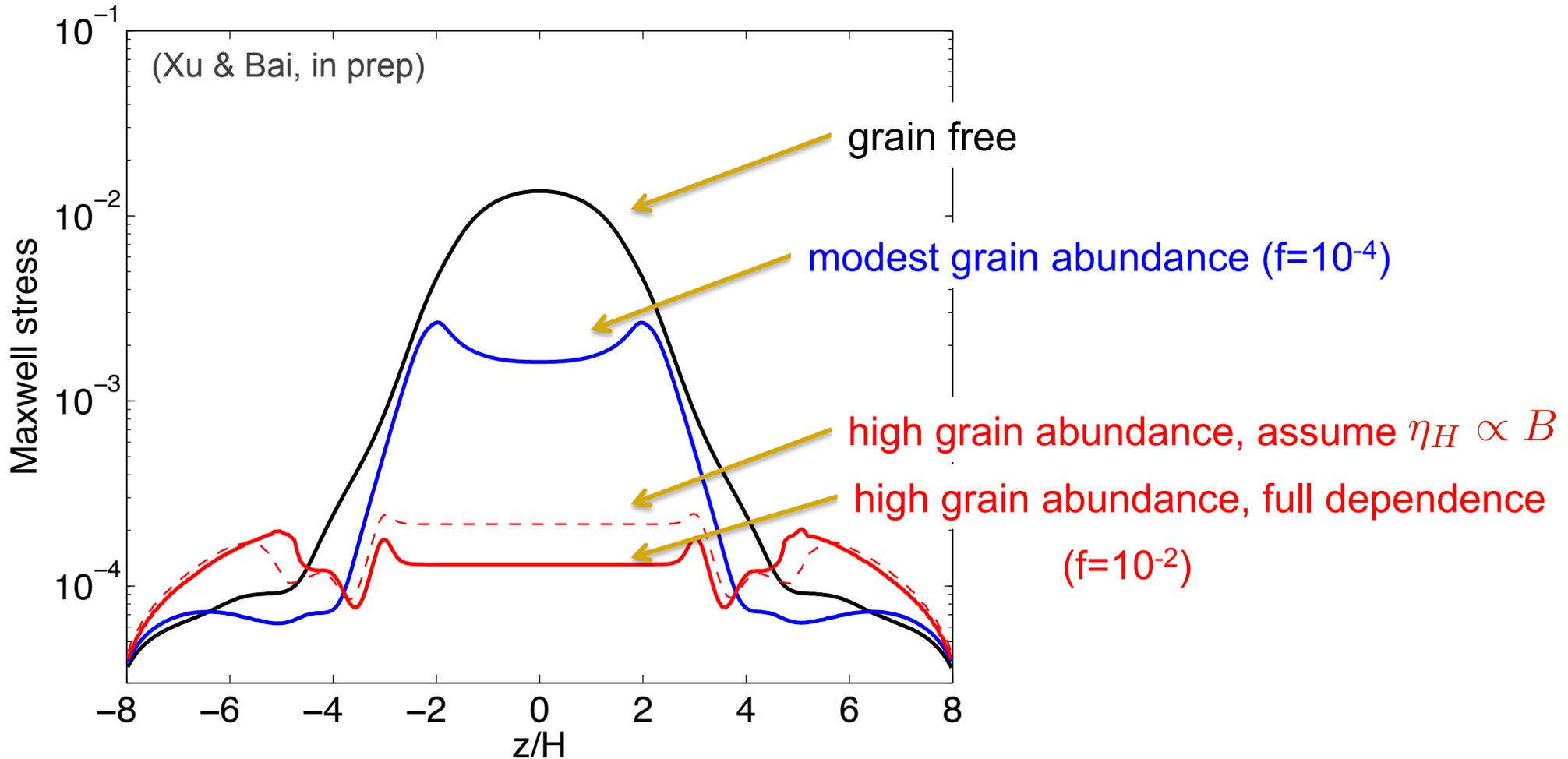
Stratified simulations in shearing-box: Lesur+14, Bai 14,15, Simon+15

# Dependence on grain abundance

- **Grain-free:** (Lesur+14, Simon+15)
  - $\Omega \cdot B > 0$  : Very strong B field amplification to near equipartition, considerable  $\alpha$  (up to 0.1) to transport AM.
  - $\Omega \cdot B < 0$  : Bursty behavior with enhanced  $\alpha$  at ~5-10 AU.
- **Modest grain abundance (0.1 $\mu$ m, 0.01 solar):** (Bai, 14,15)
  - $\Omega \cdot B > 0$  : Modest B field amplification well below equipartition, enhanced but small  $\alpha$  ( $\ll 0.01$ ).
  - $\Omega \cdot B < 0$  : Horizontal B field reduced to almost 0, no bursty behavior.
- **High grain abundance (0.1 $\mu$ m, solar):** (Xu & Bai, in prep)
  - $\Omega \cdot B > 0$  : Hall diffusivity changes sign at normal field strength, B field amplification suppressed.

# Dependence on grain abundance

1 AU,  $\Omega \cdot B > 0$  :



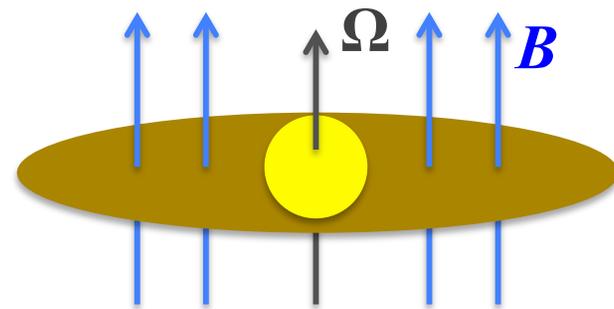
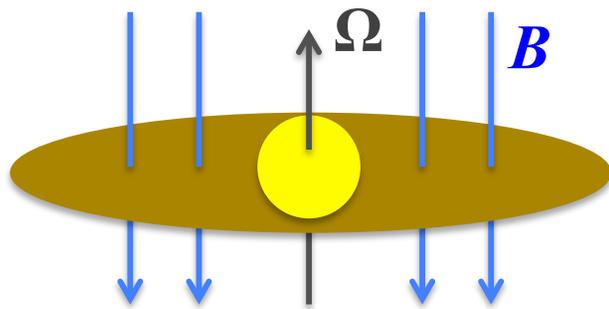
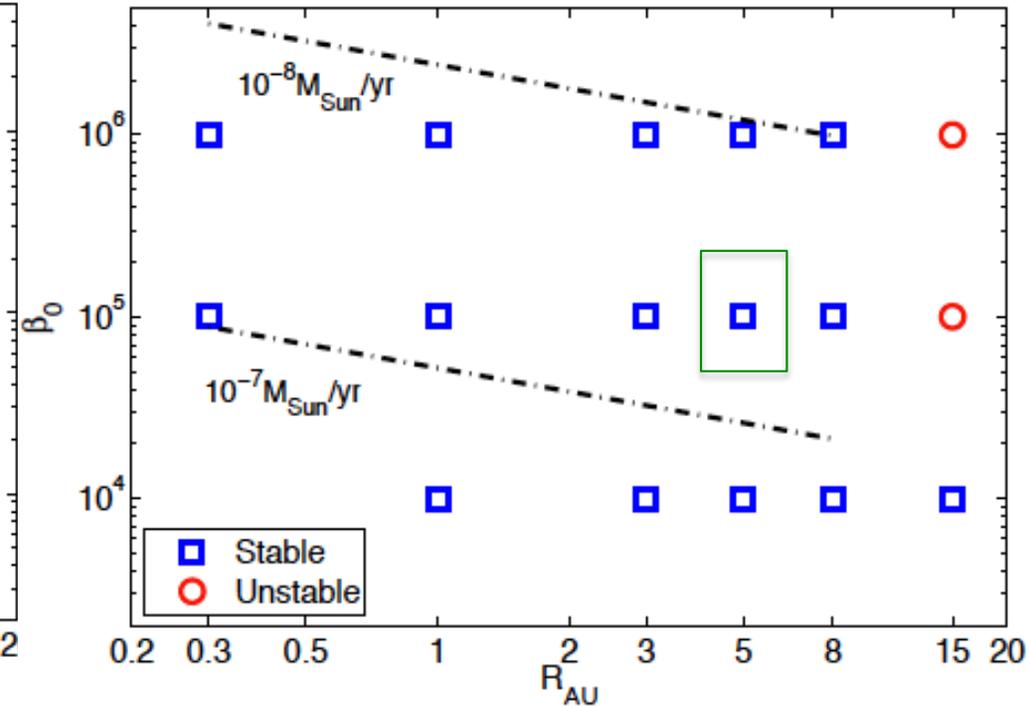
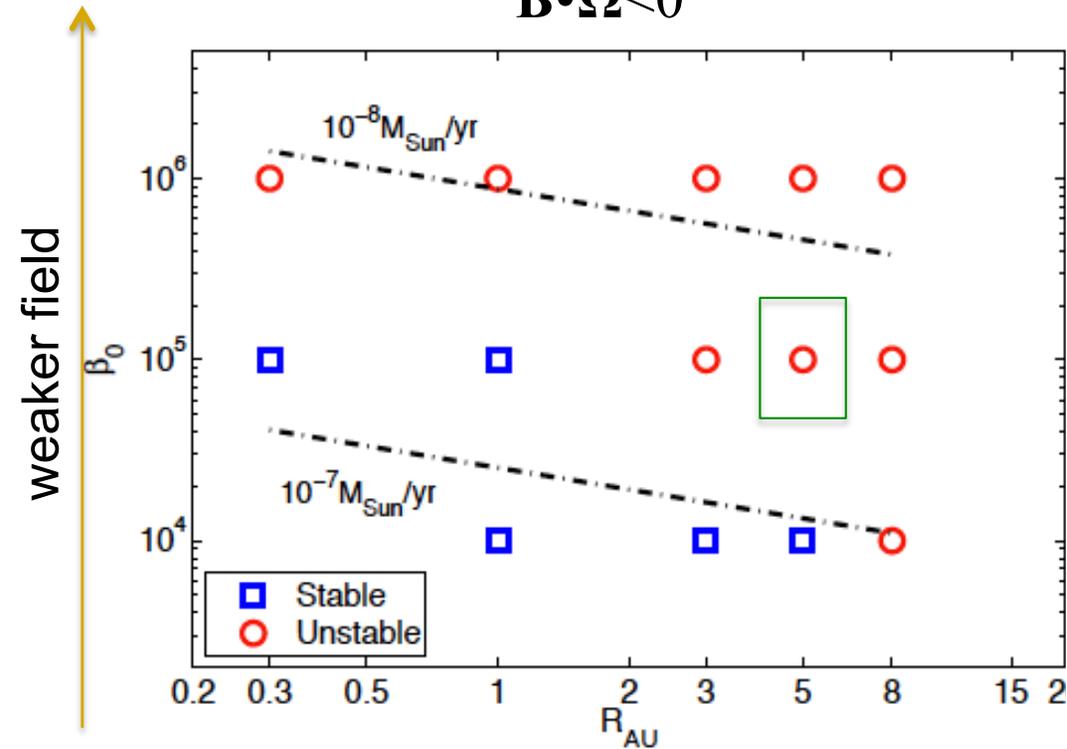
Assume grain size  $a=0.1\mu\text{m}$

Wind transport is still more efficient.

# Wind solutions modified, but some gets unstable

$\mathbf{B} \cdot \boldsymbol{\Omega} < 0$

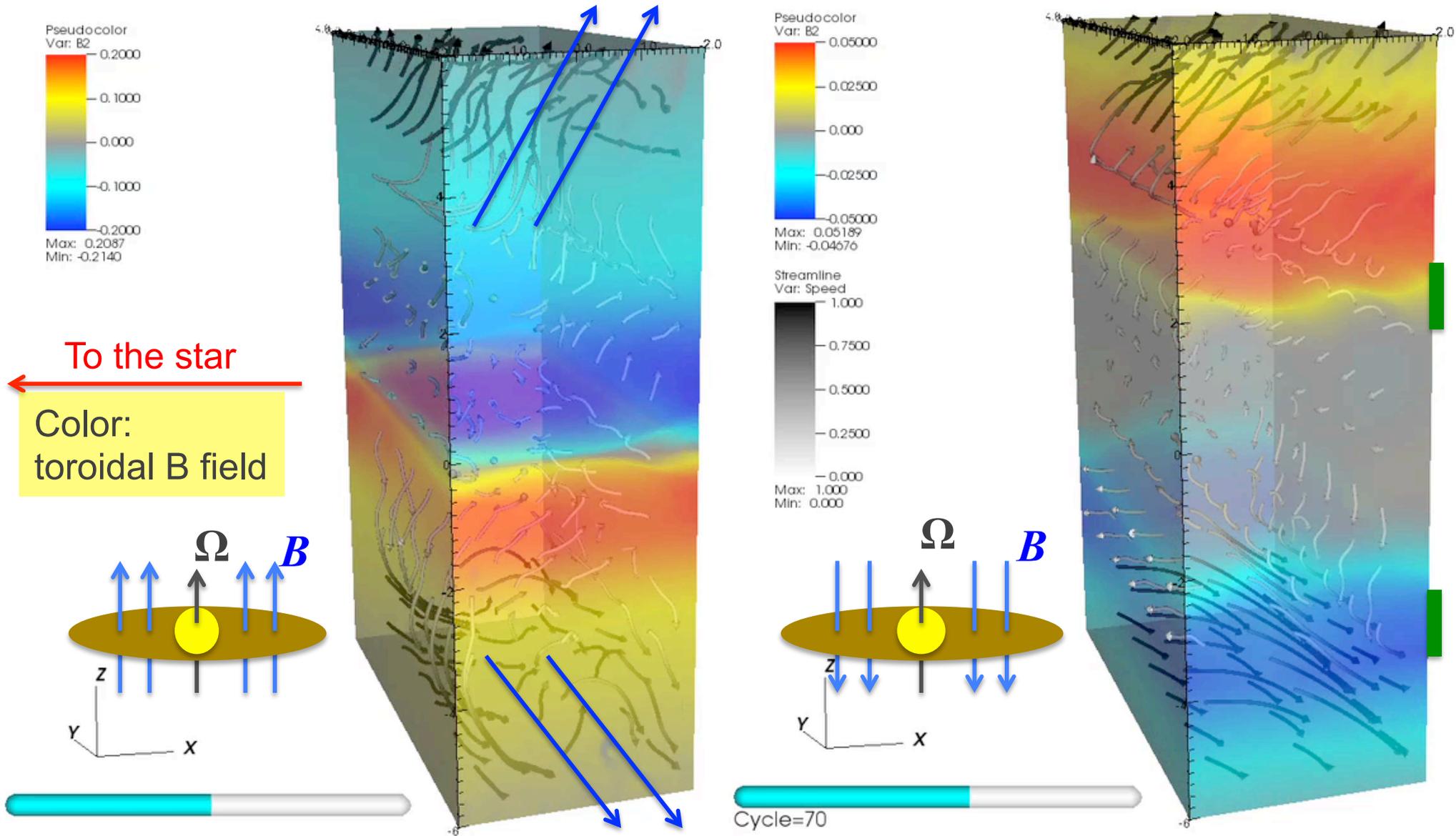
$\mathbf{B} \cdot \boldsymbol{\Omega} > 0$



(Bai, 2014)

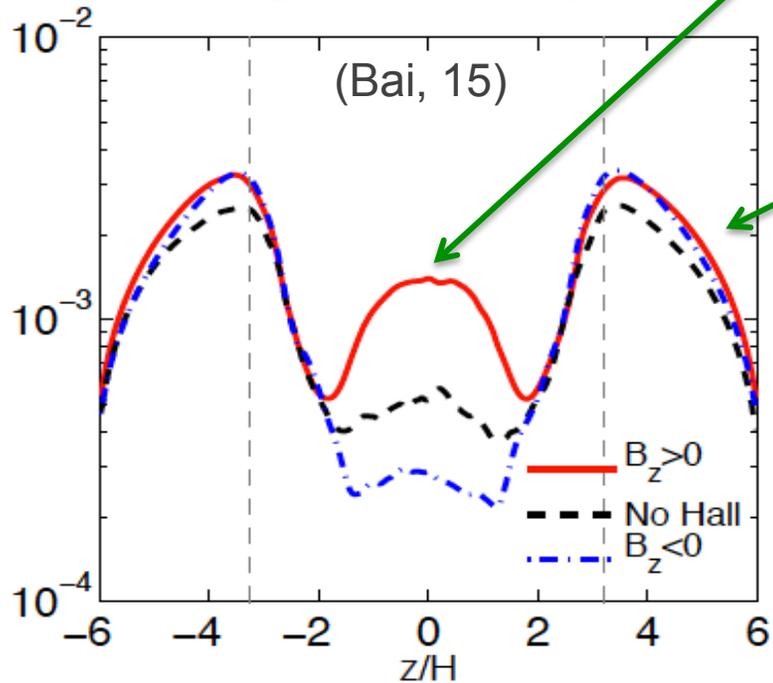
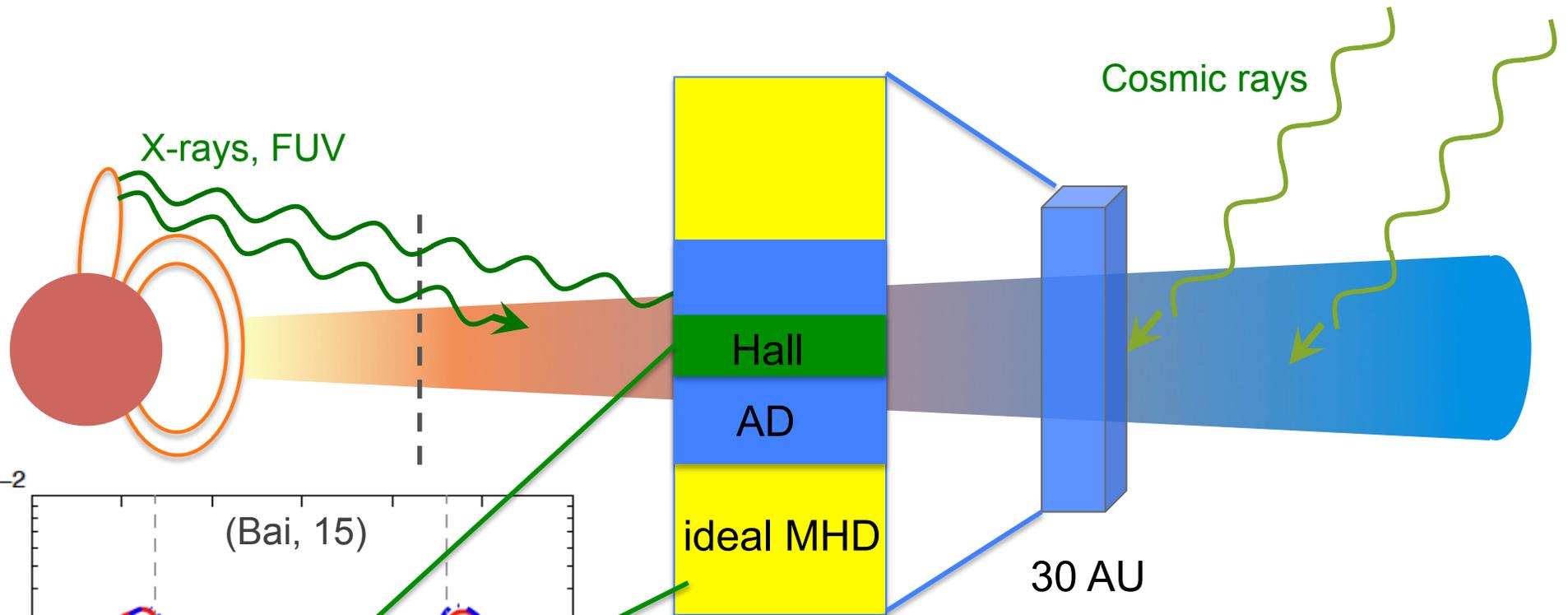
# Representative results at 5 AU

$$\beta_{z0} \sim 10^5$$



# Adding the Hall effect: outer disk

(Bai 15, Simon+15)



The Hall effect is only modestly important in the outer disk.

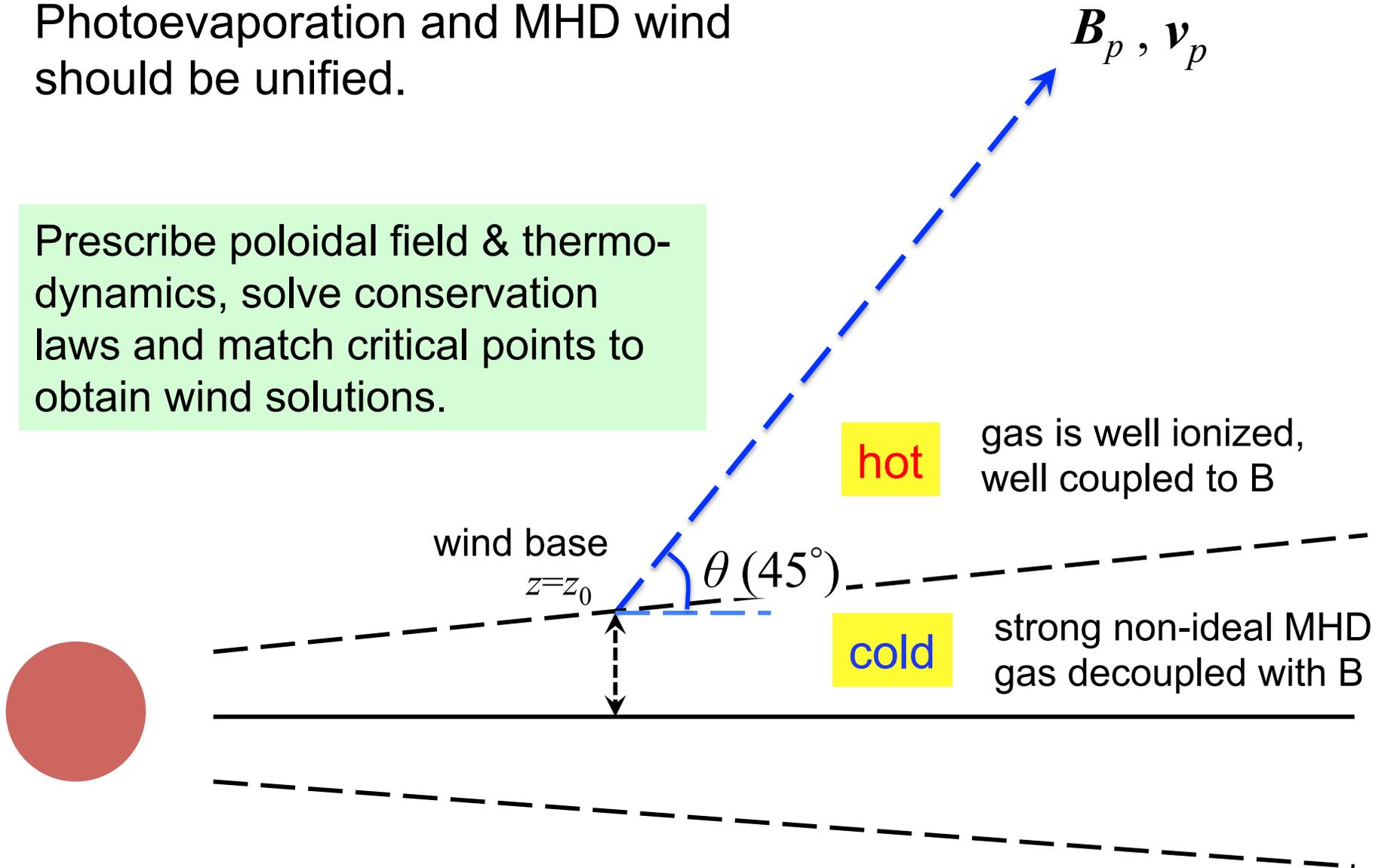
# Current status

- Net vertical magnetic flux is essential through entire disk.
- Disk microphysics is much better understood, but its parameter dependence is complex, especially on the ionization chemistry and grains (in the inner disk).
- Launching of disk wind is natural consequence of net  $B_z$ , with two major issues to be addressed:
  - Global wind structure and kinematics.
  - Local simulations have issues with wind symmetry.
- Level of turbulence/role of MRI less clear in outer disk (e.g., Simon+15, Flaherty+15): depend on FUV penetration (A. Gomez de Castro's talk). Wind likely dominates transport.

# Mass loss from PPDs: magnetized wind

Photoevaporation and MHD wind should be unified.

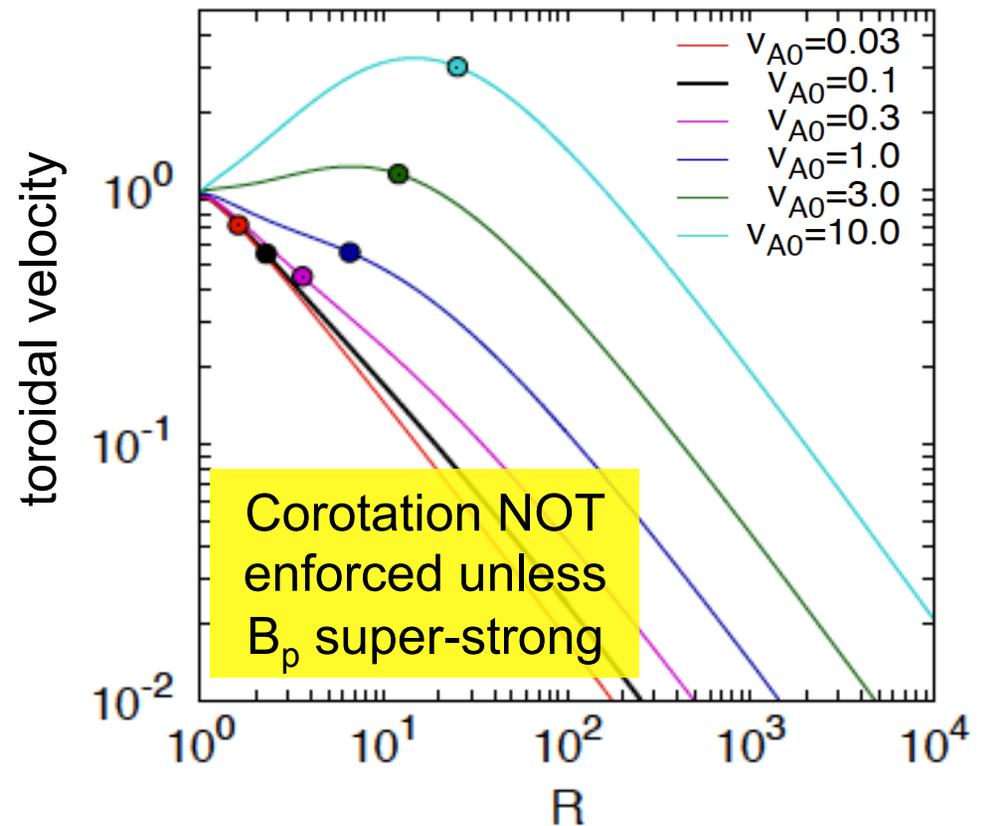
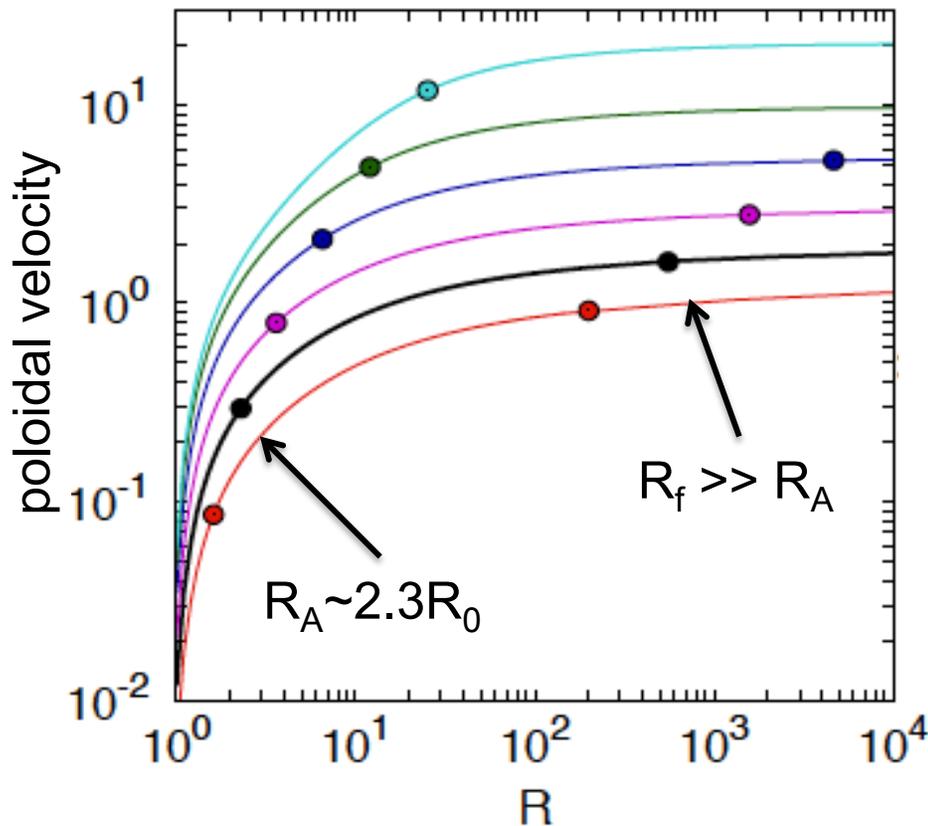
Prescribe poloidal field & thermodynamics, solve conservation laws and match critical points to obtain wind solutions.



(Bai, Ye, Goodman & Yuan, to be submitted)

# Wind kinematics

Expected wind launching conditions @ 1AU:  $v_{A0} \approx 0.1 v_K$

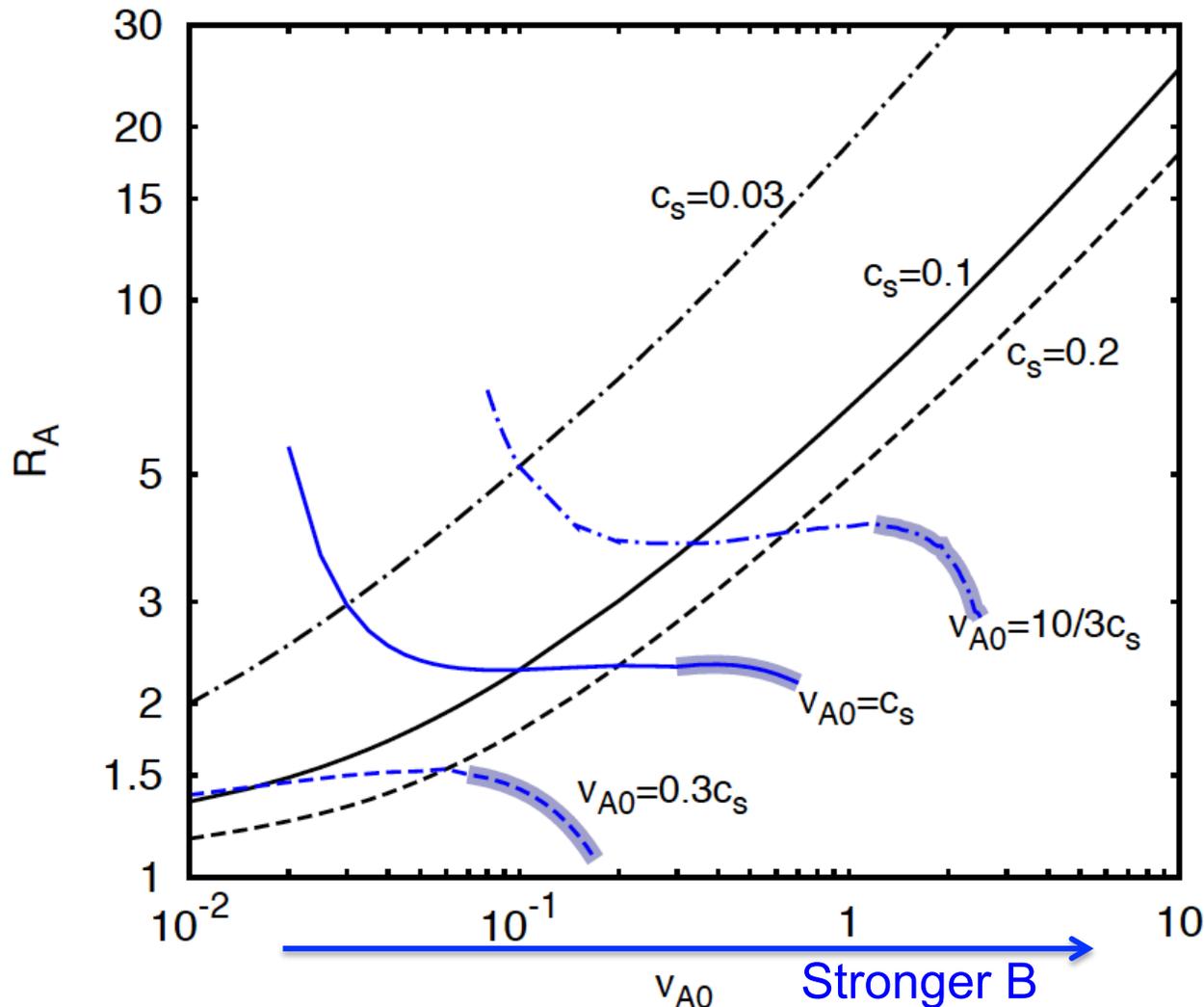


$$\frac{d\dot{M}_{\text{wind}}}{d \ln R} = \frac{1}{2} \frac{\dot{M}_{\text{acc}}}{(R_A/R_0)^2 - 1}$$

$\sim 0.12$  (can be significant)

In general, PPD wind is driven by magnetic pressure gradient.

# Thermal effect: magneto-photoevaporation?



The Alfvén radius strongly depends on

$$v_{A0}/c_s$$

B field strength

External heating

Both effects strongly affect the efficiency of wind-driven accretion!

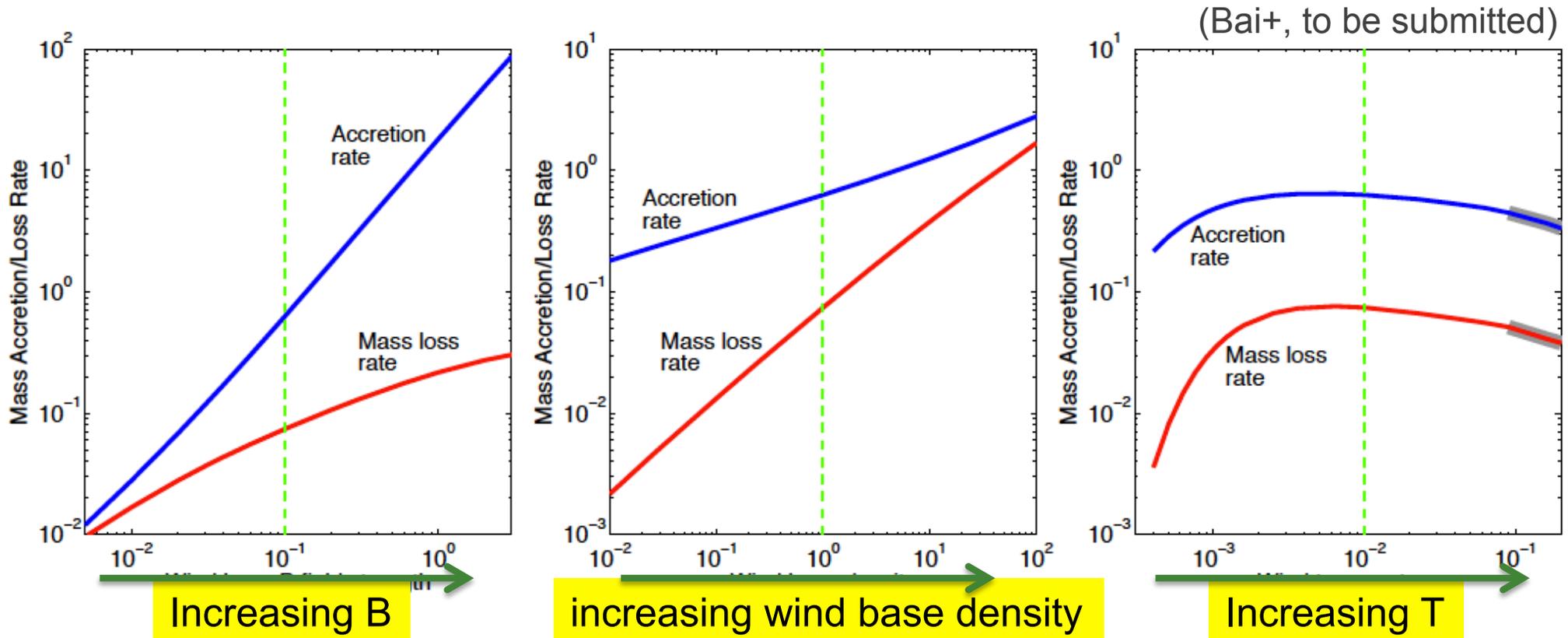
With B field, mass loss rate is always higher than pure thermally driven wind.

# Towards a global picture: wind-driven accretion

- Dependence on  $v_{A0}/c_s$  translates to:

$$\text{Note } v_{A0} = \frac{B_{p0}}{\sqrt{4\pi\rho_b}}$$

- Physical field strength  $B$
- Penetration depth of FUV  $\Rightarrow$  density at wind base
- Wind temperature

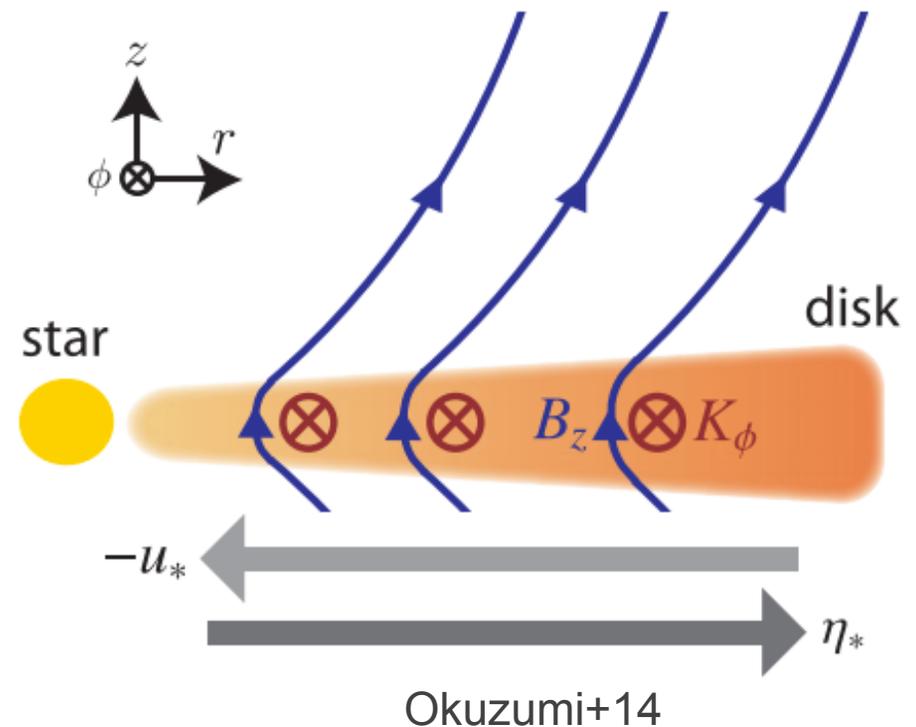


# Towards a global picture: wind-driven accretion

- Need to know the distribution of magnetic flux

## Magnetic flux transport:

Advection-diffusion framework  
(Lubow+94) with more recent  
development (Guilet & Ogilvie 12-14,  
Okuzumi, Takeuchi+14)



Need to incorporate wind and  
non-ideal MHD physics.

(e.g., Bai 14, Tsukamoto+15)

# Towards a global picture: wind-driven accretion

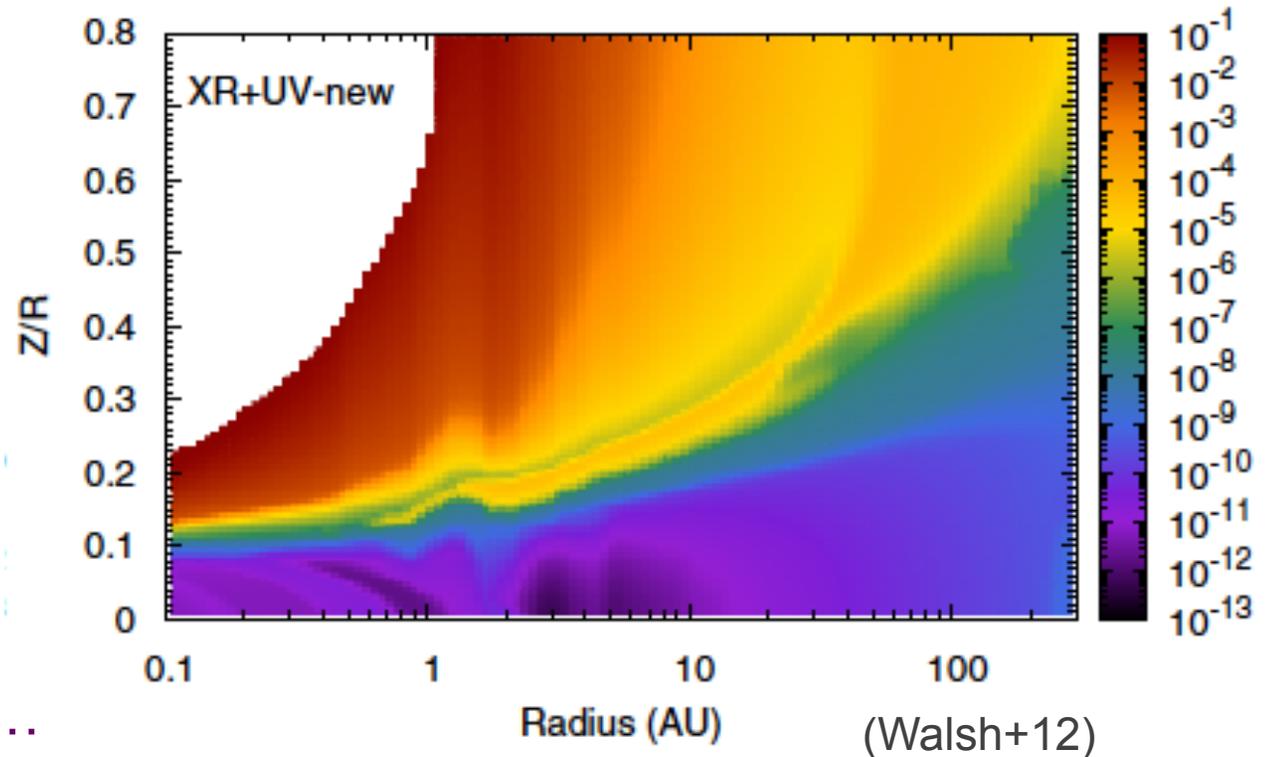
- Need to know how deep FUV penetrates

A complex task of coupling UV radiative transfer and photo-chemistry.

(e.g., Glassgold+04, Nomura & Millar 05, Woitke+09, Walsh+10,12, Bethell & Bergin 11, Akimkin+ 13)

and with wind dynamics...

(e.g., Panoglou+12)



Main source of UV opacity on tiny dust grains: how abundant are they?

# Summary

- **Initial condition: disk formation to be better understood**
  - Hall effect can induce a bimodality in disk size, as well as amount of B flux
- **MRI is disfavored, accretion is largely wind driven**
  - Inner disk: suppressed by Ohmic+AD; outer disk: strongly damped by AD
  - MND wind launching from far-UV ionization front.
- **Hall effect is polarity dependent, important at inner disk**
  - Disk dynamics depends on grain abundance: more grains suppress activities.
  - Despite unsettled issues, wind still likely dominates AM transport.
- **Disk wind-driven accretion as the key to global disk evolution**
  - Results strongly depend on both B flux distribution and thermal effects
  - Need to better understand magnetic flux transport