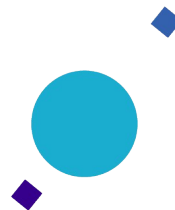


# Magnetospheric accretion and protostellar outflows in YSOs: a theoretical perspective

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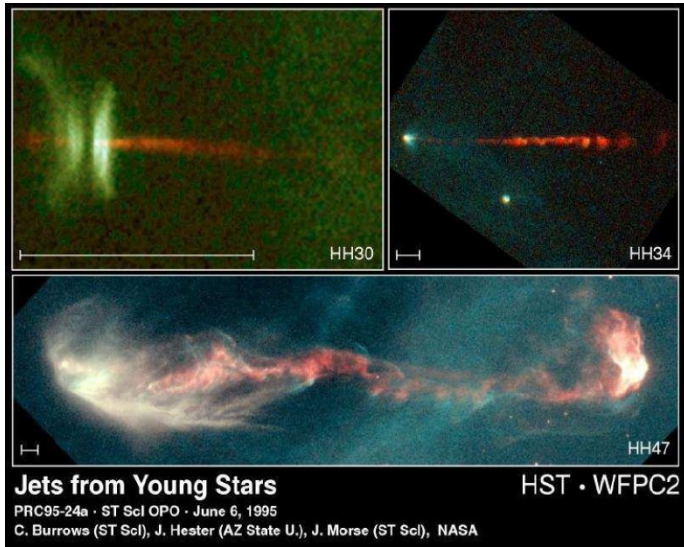
The accretion/outflow connection in young stellar objects – ESA/ESTEC

October 27<sup>th</sup> 2015

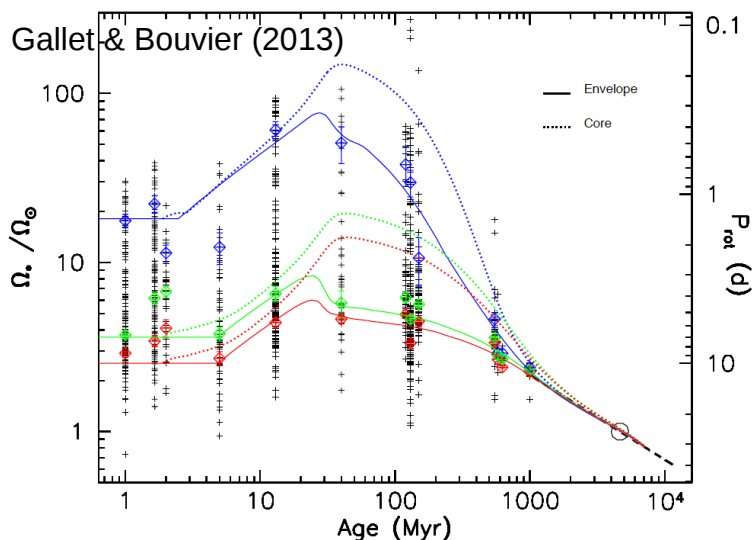
# Magnetospheric outflows

- Outflows are likely associated with magnetospheric accretion
- Magnetically driven outflows need two basic ingredients:
  - **rotation** (star, disk)
  - **large scale magnetic field** (measured, both intensity and topology)
- Combination of rotation and ordered magnetic field generates toroidal field  $B_\phi$
- $B_\phi$  provides:
  - **Poynting flux**  $E_\perp B_\phi / 4\pi$
  - **Magnetic torque**  $- B_\phi B_p / 4\pi$
  - **Poloidal force**  $-\nabla_{//} (rB_\phi)^2 / 8\pi r^2$
  - **Confinement tension**  $- B_\phi^2 / 4\pi r$

# Why bother?

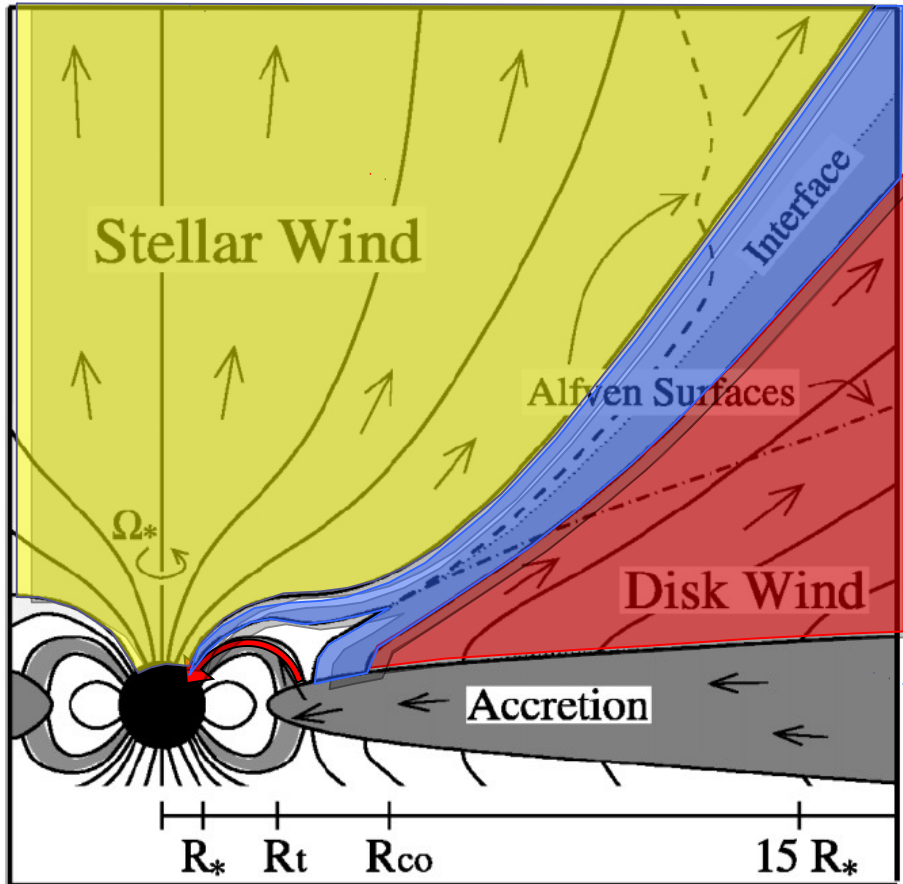


- Jets and outflows are clearly observed, showing strong correlation with accretion signatures, but not close enough to the source to give precise hints about their origin.



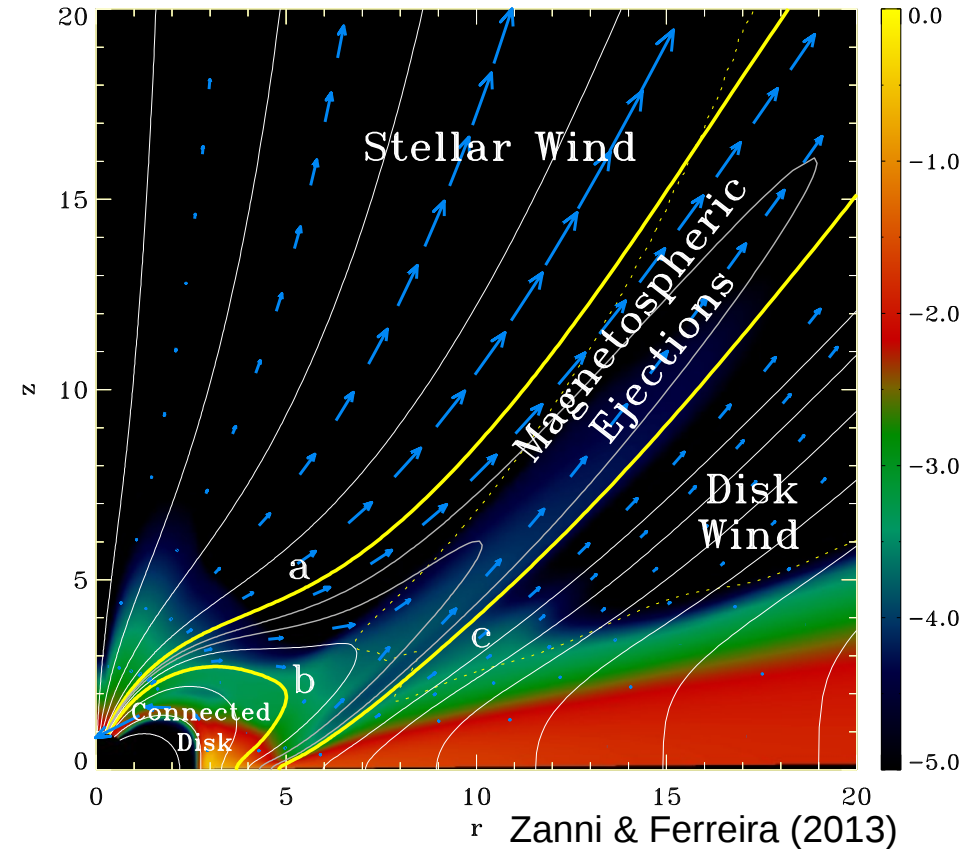
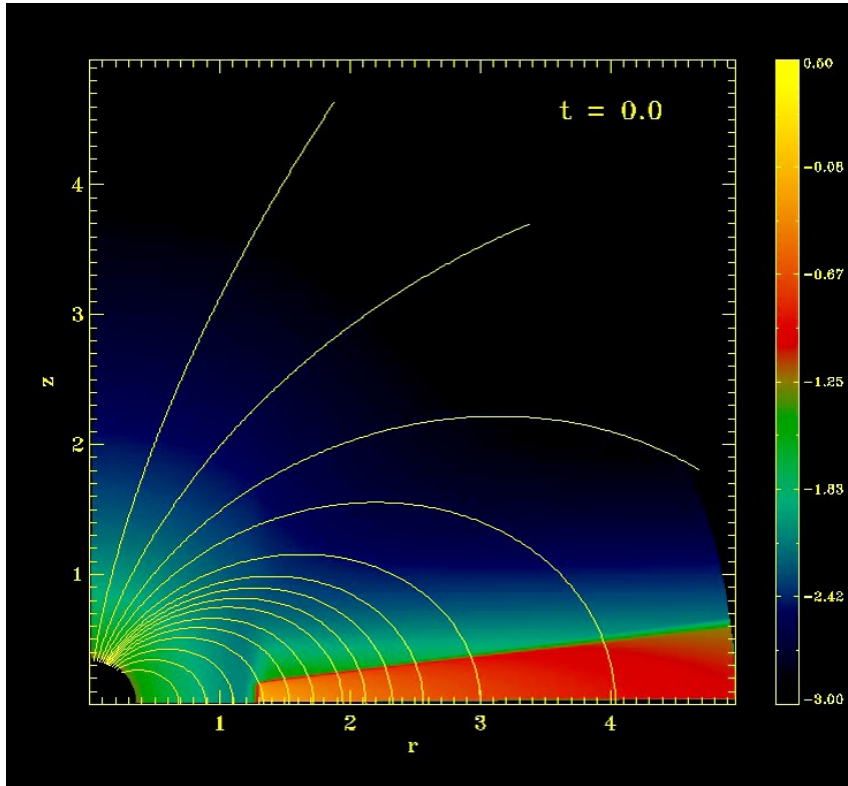
- CTTs show slow rotation periods, well below breakup (3-10 d Bouvier et al. 1993), that appear to be rather constant during the TTauri phase (Irwin et al. 2007). Can magnetospheric outflows extract the excess angular momentum?

# Magnetospheric outflows



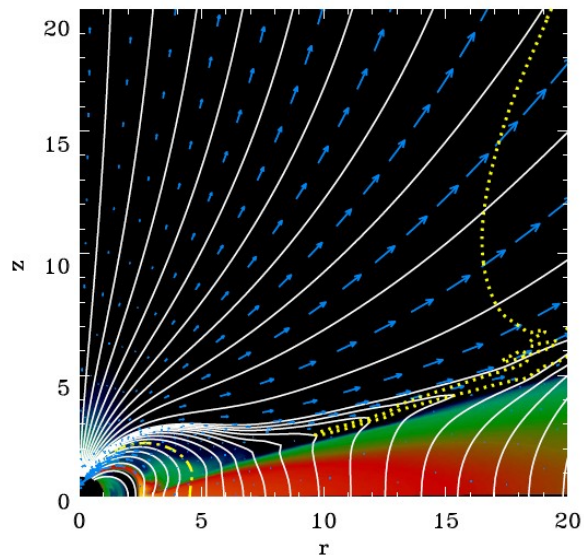
- **Stellar winds** extracting angular momentum from the star (Matt & Pudritz 2005)
- **Disk-winds** extracting disk angular momentum before it is accreted (e.g. X-wind, Shu et al. 1994)
- **Magnetospheric ejections** extracting angular momentum from the disk AND the star (Zanni & Ferreira 2013)

# A numerical model

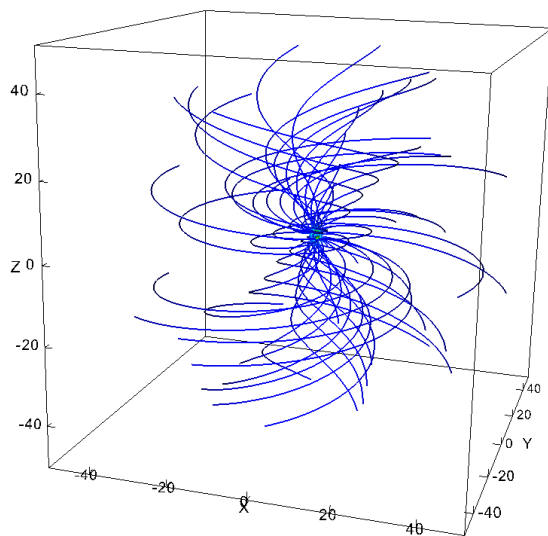


- Dipolar field aligned with the rotation axis of the star ( $B_* = 800 \text{ G}$ )
- Star ( $M_* = 0.5M_{\text{sun}}$ ,  $R_* = 2R_{\text{sun}}$ ) rotating with a 4.6 days period ( $\Omega_* = 0.1\Omega_k$ )
- Viscous and resistive “alpha” (Shakura & Sunyaev) disk with  $M_{\text{acc}} = 6 \times 10^{-9} M_{\text{sun}}/\text{yr}$

# Stellar winds

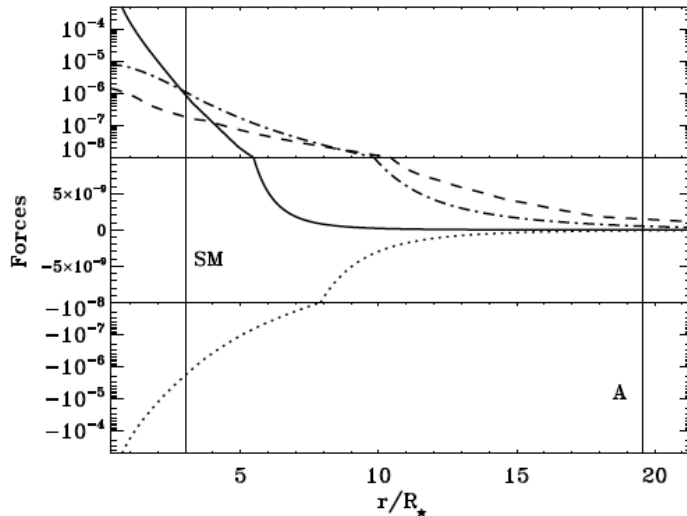


Zanni & Ferreira (2009)



- **Slowly rotating star:**
  - no centrifugal thrust
  - thermal driving ruled out (Matt & Pudritz 2007)
    - Virial temperature  $\sim 10^6$  K, no HeI or H I features
    - Too strong X-ray emission
  - Turbulent Alfvén waves (DeCampli 1981) driven by the accretion power? (Matt & Pudritz 2005)
- **Ejection rate likely low**
  - $M_{\text{wind}} < 1\% M_{\text{acc}}$  (Cranmer 2009)

# Stellar winds



- An example (Zanni & Ferreira 2009):

- $M_{\text{wind}} \approx 1\% M_{\text{acc}}$

- Lever arm  $R_A/R_* \approx 19$

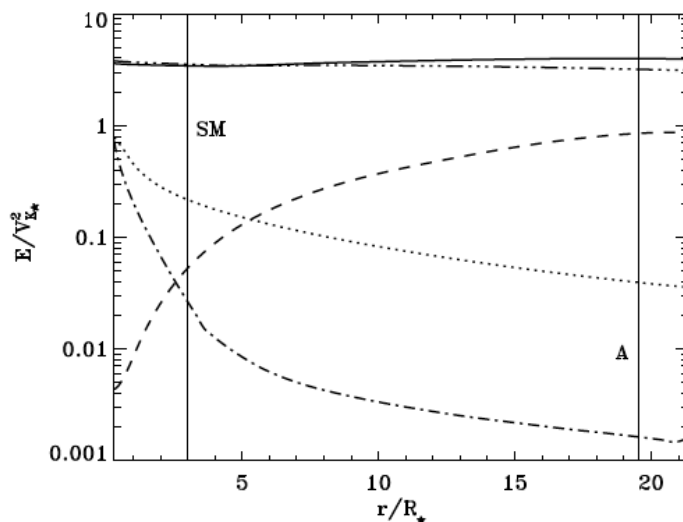
- Braking torque:  $M_{\text{wind}} R_A^2 \Omega_*$

- $P_{\text{wind}} \approx 2\% P_{\text{acc}}$

- Potentially very fast component ( $\sim 400 \text{ Km s}^{-1}$ ), but most of the energy still in Poynting flux

- Poor (self)collimation

- Despite large lever arm only 20% of the accretion torque is extracted



# Stellar winds and spin-down

- Is the stellar wind capable of balancing the accretion torque and brake the stellar rotation? (Matt & Pudritz 2005)

$$r_A = 2.11 \left( \frac{B_\star^2 R_\star^2}{\dot{M}_{\text{wind}} v_{\text{esc}}} \right)^{0.22} R_\star \quad \text{and} \quad \dot{M}_{\text{acc}} \sqrt{GM_\star R_t} = \dot{M}_{\text{wind}} r_A^2 \Omega_\star$$

Matt & Pudritz (2008, 2012)

- In order to balance a “Keplerian” accretion torque ( $\dot{M}_{\text{acc}} \sqrt{GM_\star R_t}$ ) of a star rotating at a small fraction of breakup (less than 10%) even assuming a kG dipolar component, a more **massive and powerful wind is required**

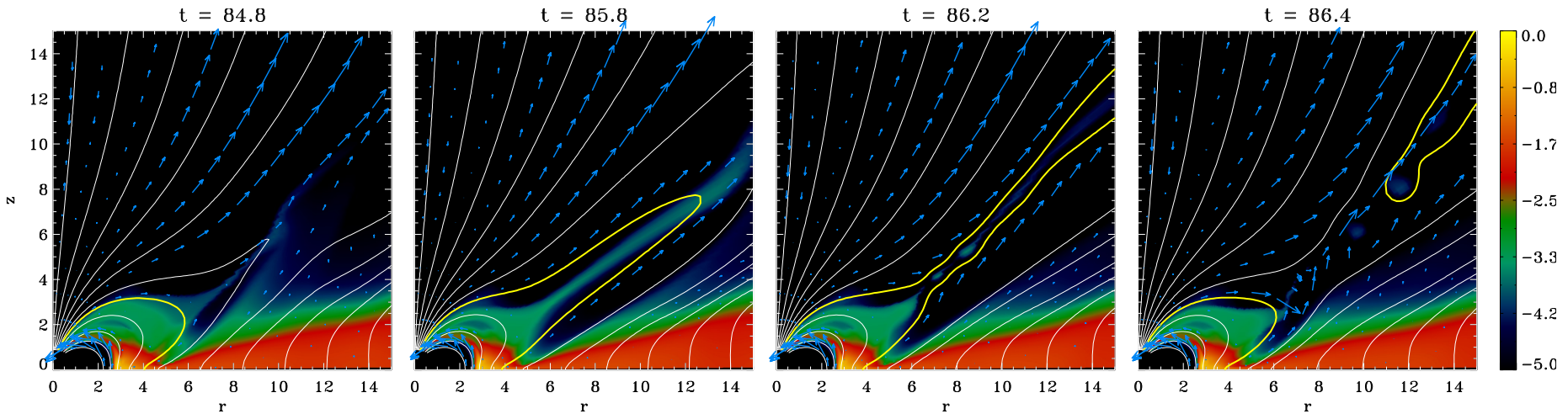
$$\dot{M}_{\text{wind}} > 10\% \dot{M}_{\text{acc}}$$

$$P_{\text{wind}} > 10\% P_{\text{acc}}$$

- For  $\dot{M}_{\text{acc}} \leq 10^{-8} \dot{M}_{\text{sun}}/\text{yr}$  the **contraction** timescale (Kelvin-Helmholtz  $\approx 10^6$  years) becomes comparable to the accretion spin-up timescale. **The mass and energy requirement becomes even more important**

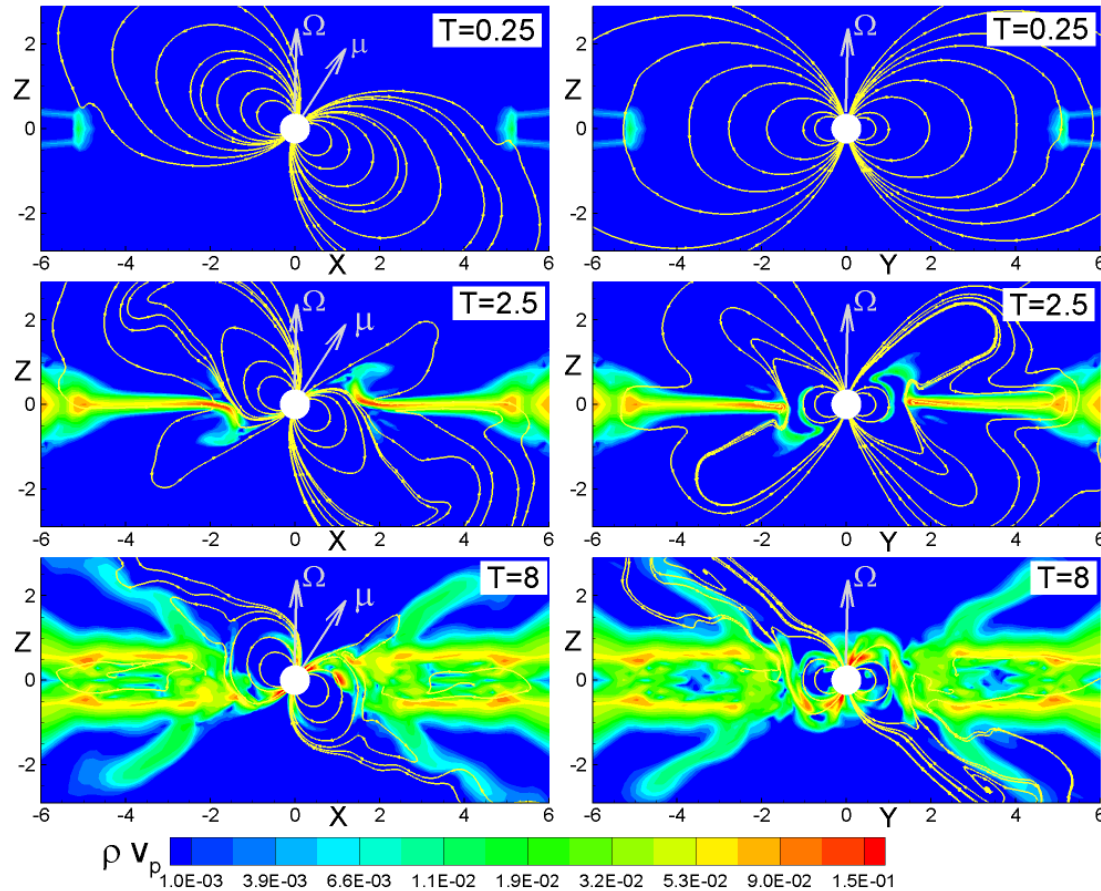


# Magnetospheric ejections



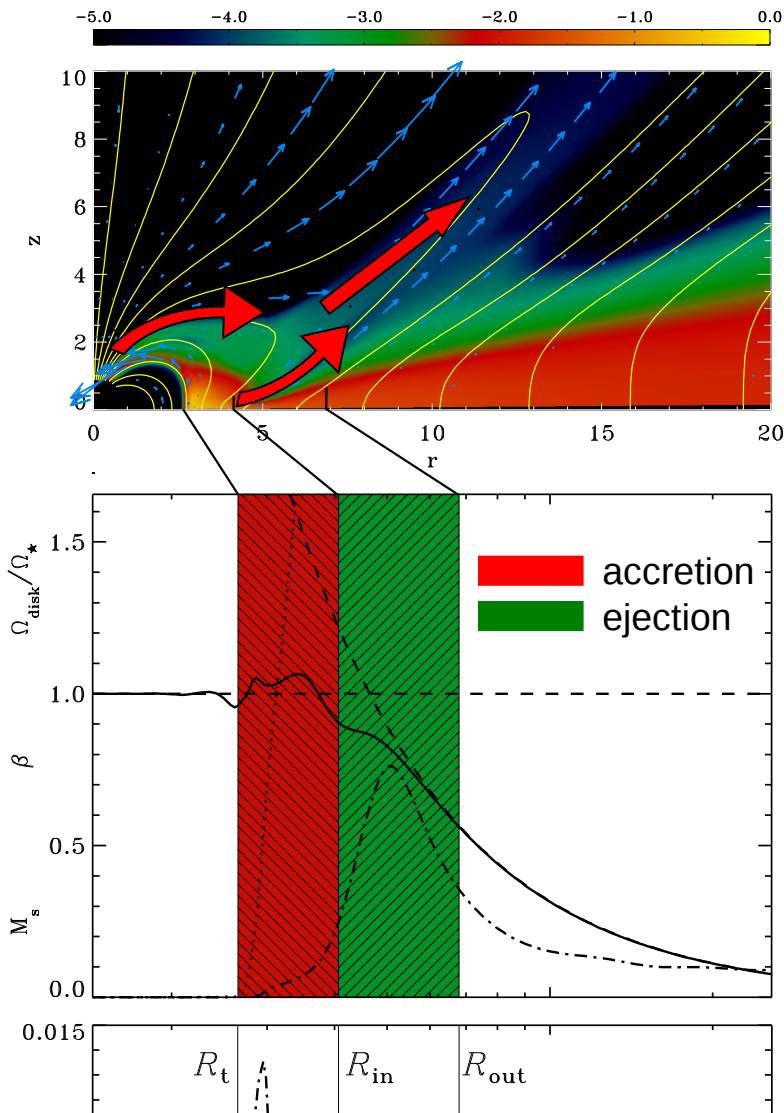
- Star-disk differential rotation builds up toroidal field component
- Due to the  $B_\phi$  pressure, initial potential structure (dipole) relaxes by inflating and opening
- Magnetospheric ejections **flow along the interface between the disk and the stellar open field lines**
- On a spatial scale around 10-20 stellar radii the ejection **disconnects in a reconnection event** (plasmoids)

# Conical winds (Romanova et al. 2009)



- Magnetospheric ejection phenomenon **closely related to “conical winds”** (Romanova et al. 2009)
- **Observable even in 3D** models with misaligned dipole!

# Magnetospheric ejections



- Magnetospheric ejections exchange angular momentum both with the disk and the star

- From the disk:

$$\dot{M}_{ej}/\dot{M}_{acc} \approx 0.1$$

$$\dot{J}_{ej}/\dot{J}_{acc} \approx 0.5 \quad (\text{"Keplerian" torque} \approx 0.34)$$

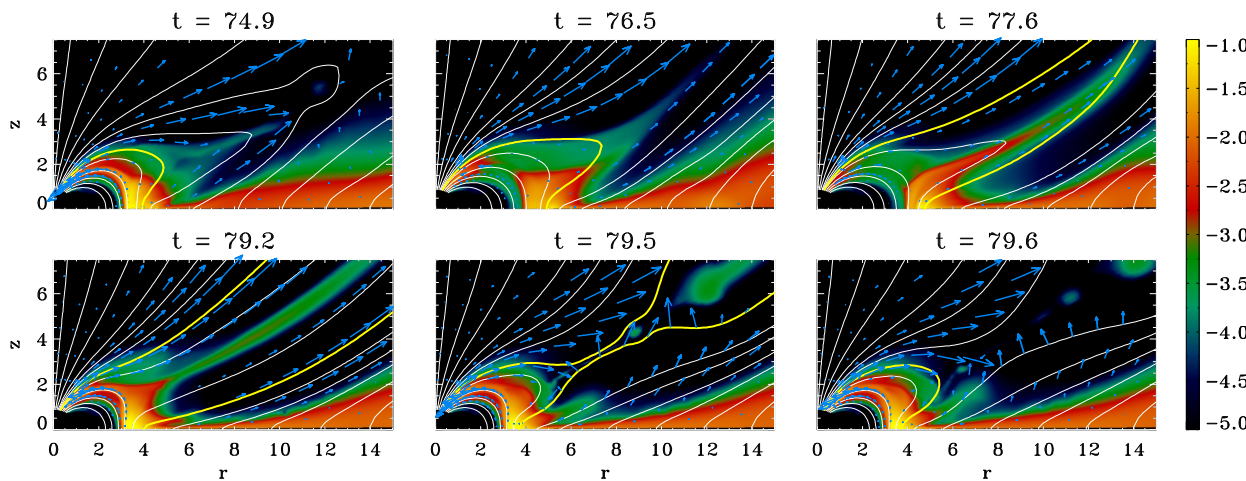
- The torque is strong enough so that the disk rotate slower than the star

- If also the MEs rotate slower than the star it is possible to extract energy and angular momentum directly from it .

- Energy extracted from the disk and the star can not push this rather massive ejection beyond  $100 \text{ Km s}^{-1}$

# MEs and spin-down

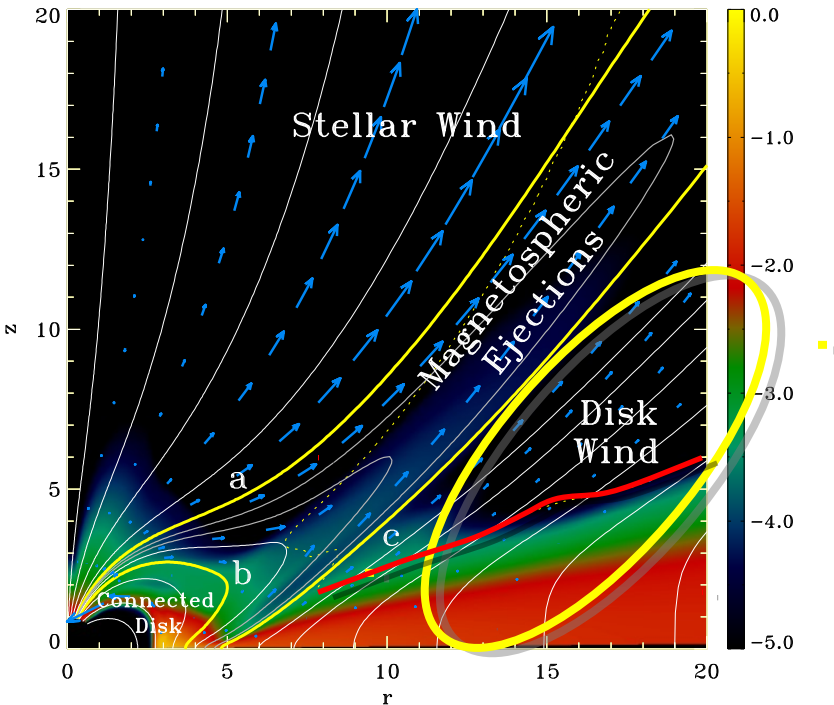
- Extracting stellar angular AM and a sizeable fraction of the disk AM before it is accreted, MEs (with some light SW) can balance the accretion torque for  $R_t \sim 0.6 R_{co}$  already.
- It has been shown (Zanni & Ferreira 2013) that in a “propeller” regime the MEs are able to balance the contraction spin-up too



- As the truncation radius gets closer to corotation ( $R_t \sim 0.75 R_{co}$  in the simulation shown), the stellar centrifugal barrier makes accretion difficult

- When the disk is pushed outwards fieldlines open up and strong outflow extracts angular momentum and pushes disk back in
- When the disk gets closer to the star accretion is possible but the magnetosphere closes again and disk is pushed out → accretion-ejection cycles

# ... and the disk-winds?

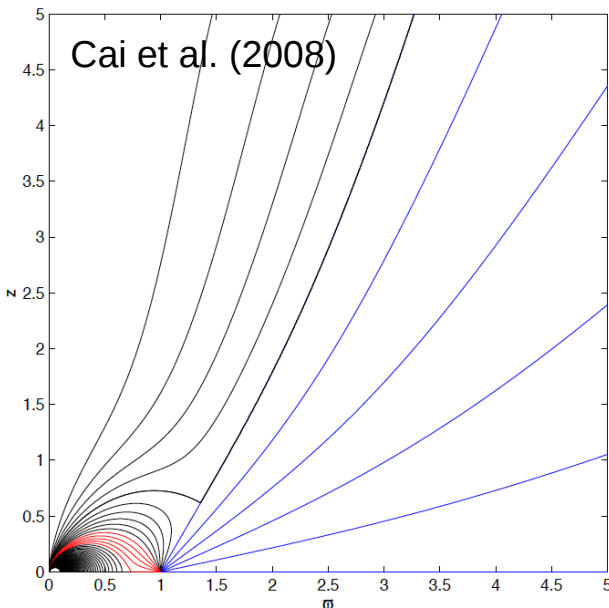


- Simulations display very weak disk winds, with a  $\approx 1\%$  mass ejection rate and Alfvén surface lying close to the disk surface
- A powerful extended disk wind seems to require a proper disk magnetic field

- Wide angle, fan-shaped powerful X-winds (Shu et al. 1994) are not observed in star-disk interaction numerical experiments.
  - wrong range of parameter space?
  - resolution limits in the launching region around corotation?

# X-Winds

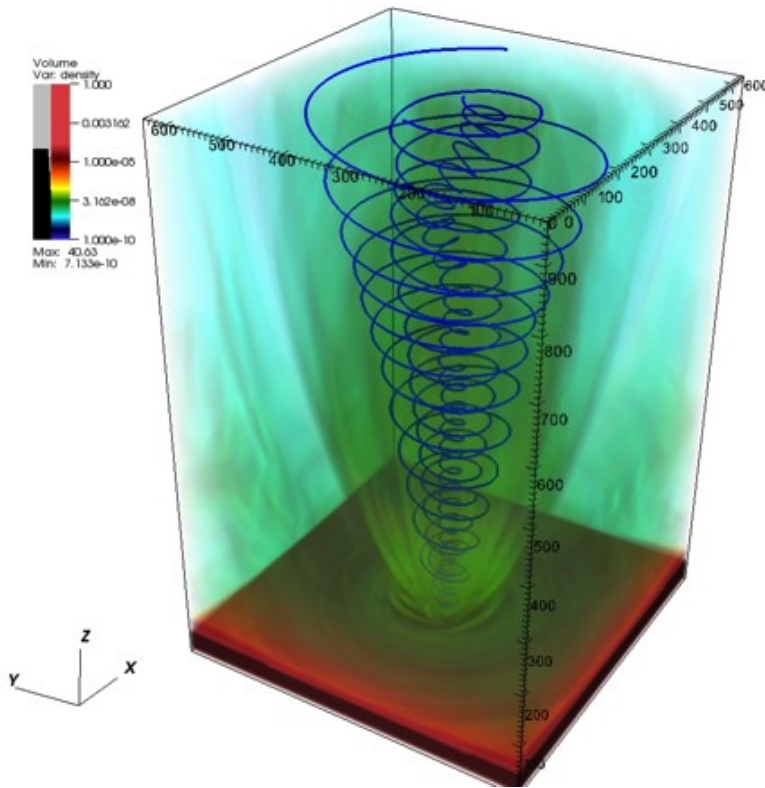
- An X-Wind (Shu et al. 1994) is a wide angle, fan shaped disk-wind launched from a small disk region around corotation where the open magnetospheric flux is trapped. It extracts a sizeable fraction of disk mass ( $M_{\text{xwind}} > 10\% M_{\text{acc}}$ ) and **ALL the disk angular momentum before it is accreted onto the star.**



- It is actually possible to calculate an outflow solution (Cai et al. 2008, Anderson et al. 2005)
- **Coupling with the disk is critical.**
  - Can the disk support this magnetic topology?
  - Is it possible to extract a high amount of mass and energy from such a tiny region? (Ferreira & Casse 2013)

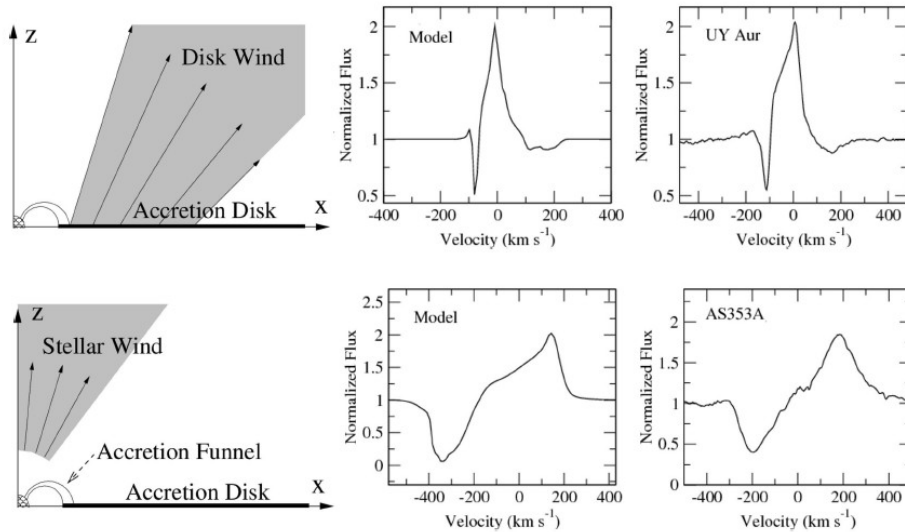


# Extended disk-winds



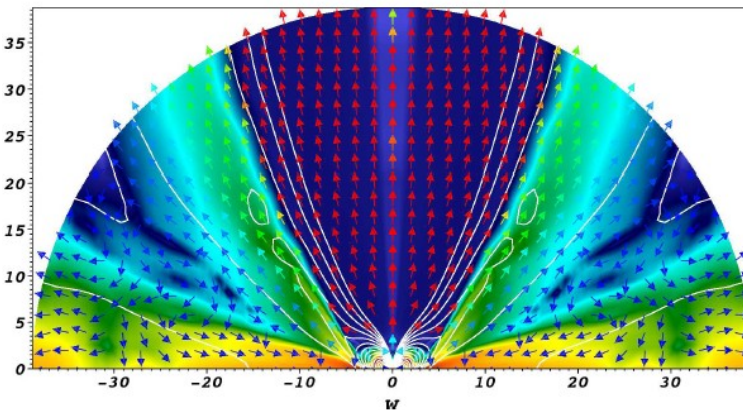
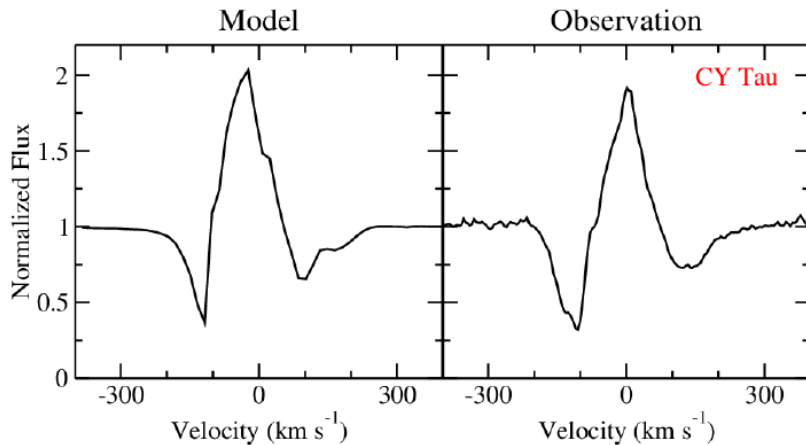
- **Powerful disk-winds** (i.e. extracting all the accretion power vs. local turbulent dissipation) **require large scale field close to equipartition** ( $\beta < 10$ )
- It is possible to compute solutions:
  - Very **fast** ( $V > 300 \text{ Km s}^{-1}$ ), less collimated, light ( $M_{\text{wind}} \approx 1\% M_{\text{acc}}$ )
  - **Slower** ( $V < 300 \text{ Km s}^{-1}$ ), more collimated with  $M_{\text{wind}} \approx 10\% M_{\text{acc}}$
- Launched from a distance from the central star disk-winds are unlikely to contribute to the stellar spin evolution problem

# Magnetospheric outflows: synthetic observations



- He I  $\lambda 10830$  line profiles
- **DW** (edge-on) shows slow and narrow blue-shifted absorption
- **SW** (pole-on) shows wide and fast blue-shifted absorption
- **Conical winds** (MEs?) similar to DW

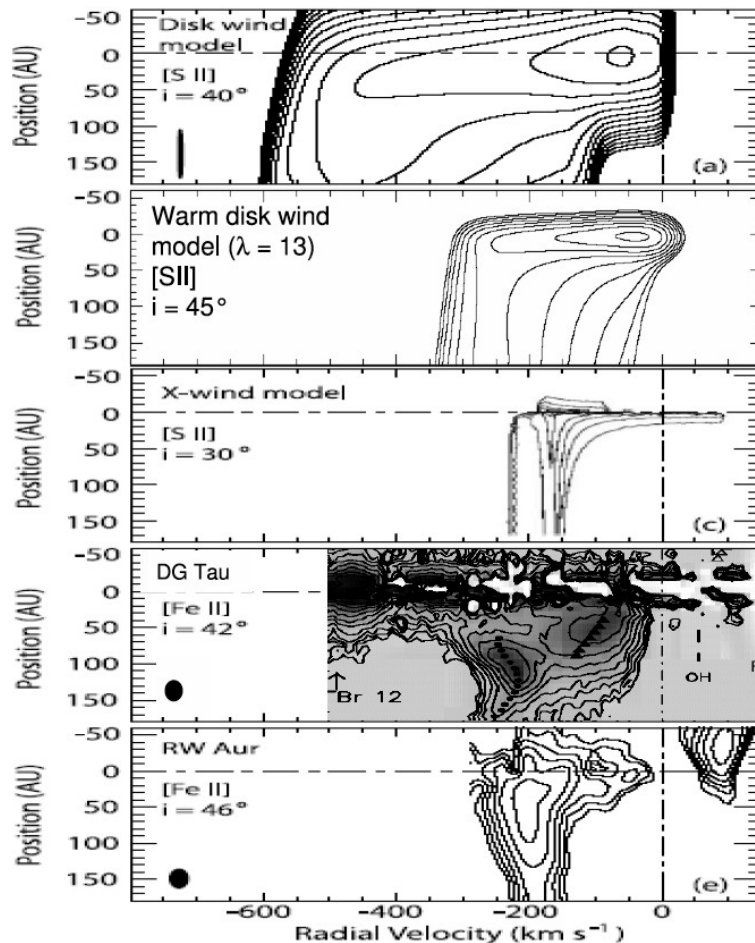
Kurosawa et al. (2011)



Kurosawa &  
Romanova (2012)



# PV diagrams



- Cold disk wind model,  $\lambda = 50$   
 $r_0 = 0.07 - 1 \text{ AU}$  (Garcia et al. 2000)
- Warm disk wind model,  $\lambda = 13$   
 $r_0 = 0.07 - 3 \text{ AU}$
- X-wind model,  $\lambda = 3.5$   
 $r_0 = 0.06 \text{ AU}$  (Shang et al. 1998)
- Pyo et al. 2003
- Pyo et al. 2006

# Summary

- **Stellar winds:**
  - Loosely self-collimated, likely low ejection rates and high terminal speeds
  - Can significantly contribute to stellar spin evolution problem, specially for higher ejection rates → driving power problem.
- **Magnetospheric ejections:**
  - Ballistic propagation (“conical winds”), rather high ejection rates and lower ( $< 100 \text{ Km s}^{-1}$ ) propagation speed
  - Significantly contribute to the stellar spin evolution, but can balance stellar contraction during propeller phases → extreme variability not observed.
- **Disk-winds:**
  - X-winds: support from numerical simulations controversial. Driving power issues.
  - Powerful extended disk winds require proper disk magnetic field
- **Line profile fitting extremely useful but ambiguous.** Spectro-interferometry?
- **Combine with larger scale diagnostics**
- **Be brave and produce more synthetic observations from simulations**