

# Super-Eddington Flows and Spectra

**Ken OHSUGA**

University of Tsukuba, JAPAN

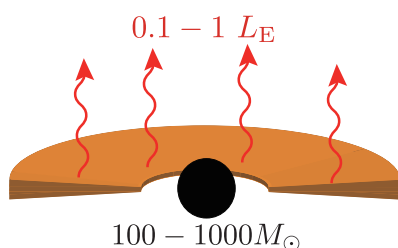
T. Kawashima, H. R. Takahashi, Y. Asahina, H. Kobayashi (NAOJ)  
T. Kitaki, S. Mineshige (Kyoto Univ.)

## Sub-Eddington or Super-Eddington

### IMBH + sub-Eddington disk

If the IMBHs exist, sub-Eddington disk can explain the huge luminosity of ULXs;

$0.1 L_{\text{Edd}}(10^3 M_{\text{sun}}) \sim 10^{40} \text{erg/s.}$

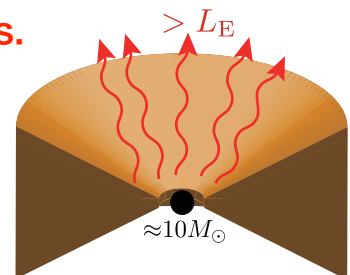


Makishima et al. 00, Miller et al. 04,  
Farrell et al. 09, Servillat et al. 11, etc.

### Super-Eddington Disk

Even if the BH mass is around  $10 M_{\text{sun}}$ , super-Eddington disks can reproduce the huge luminosity;

$10 L_{\text{Edd}}(10 M_{\text{sun}}) \sim 10^{40} \text{erg/s.}$

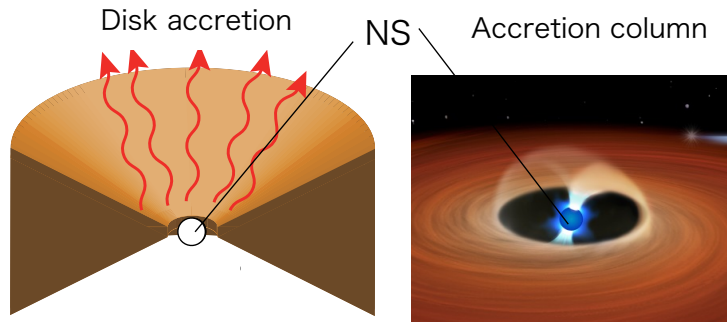


King 04, 08; Ohsuga+ 05, 09, 11,  
Poutanen+07; Gladstone+09; Middleton  
+11, Sadowski+13,15, Takahashi+16

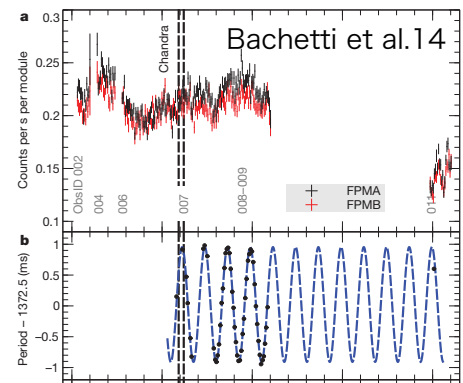
# Super-Edd. accretion onto NSs

## NS + Super-Eddington flow

If the central objects of ULXs are NSs, super-Eddington is necessary because the mass of NSs is a few  $M_{\text{sun}}$ .



Basko & Sunyaev 76; Ohsuga 07; Mushtukov+15, 18; King & Lasota 16; Kawashima et al. 16; Takahashi & Ohsuga 17, 18; Chashkina+17



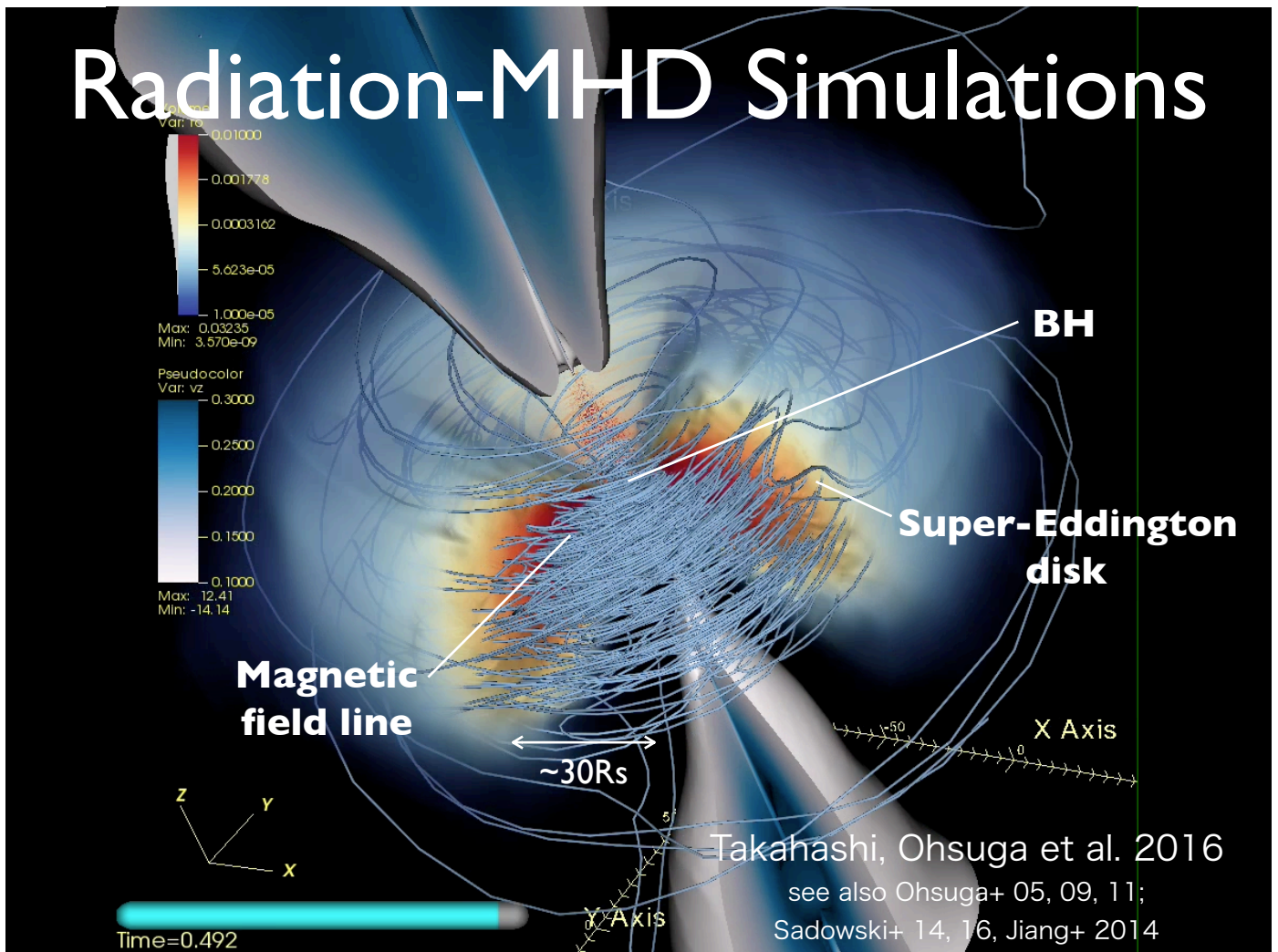
ULX pulsar;  
Engine is super-Eddington accretion onto NS.

also Fuerst+16, Israel+17, Carpano+18, etc.

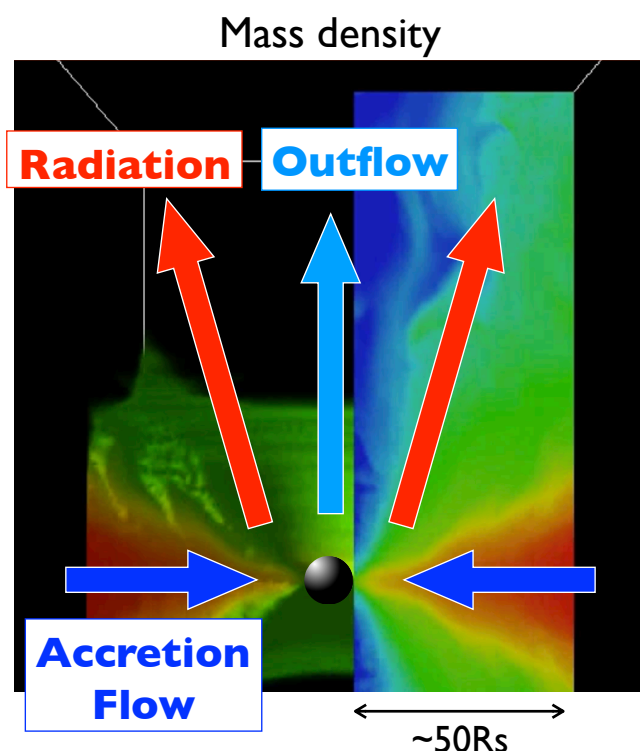
## Today's plan

- We introduce our **Radiation-HD/MHD simulations** of super-Eddington flows around BHs and explain the basic features of the super-Edd. flows.
- We show that the super-Edd. flows can explain observations of ULXs (Luminosity, spectra, clumpy outflow, ULX bubbles).
- We show the simulation results of the super-Edd. flows around NSs (Powerful outflows and accretion column appear).

# Radiation-MHD Simulations



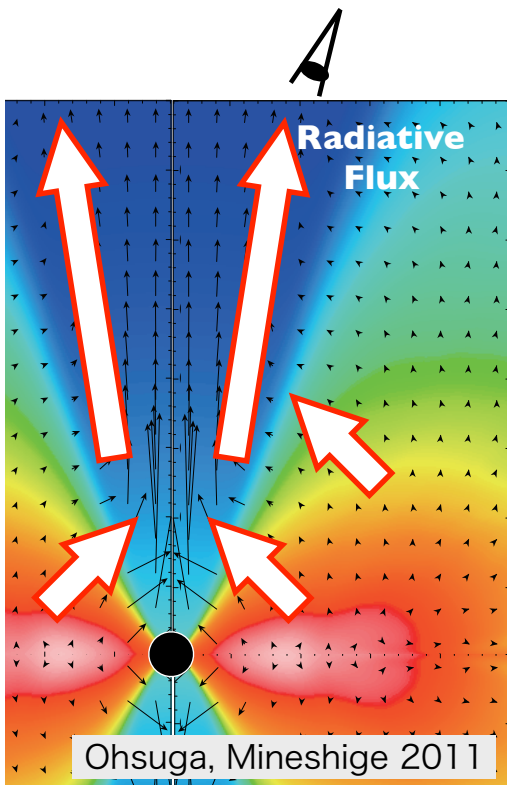
## Super-Eddington disk & outflows



Radiation-pressure  
supported disk  
&  
radiatively-driven jet  
( $\sim 0.3-0.5c$ )

Ohsuga et al. 2009  
Ohsuga & Mineshige 2011  
see also Ohsuga et al. 2005

# Luminosity



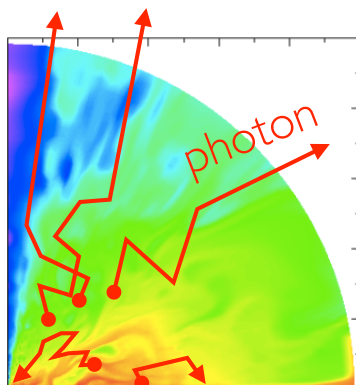
The radiative flux is mildly collimated since the disk is optically and geometrically thick.

Thus, observed luminosity is much larger than the Eddington luminosity except for the edge-on view (e.g.,  $22L_{\text{Edd}}$  for  $\lesssim 20^\circ$  in the case of  $\dot{M} \sim 100L_{\text{Edd}}/c^2$ ,  $L_{\text{disk}} \sim 3L_{\text{Edd}}$ ).

**Super-Eddington flows can explain the large X-ray luminosity of ULXs.**

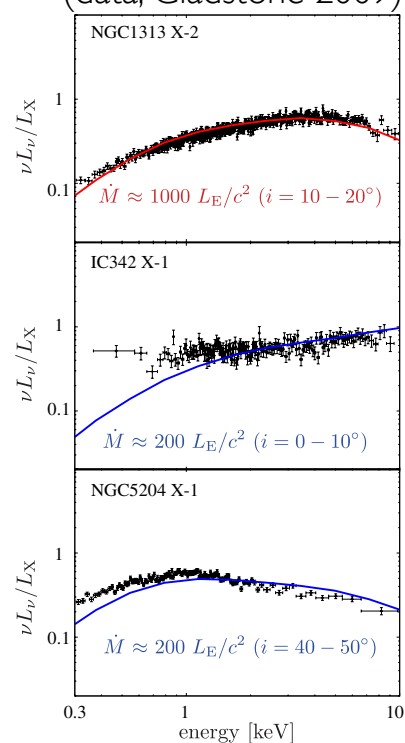
## X-RAY SPECTRA

Monte Carlo Radiation transfer  
(free-free, thermal & bulk compton)  
*Kawashima, Ohsuga et al. 2012*



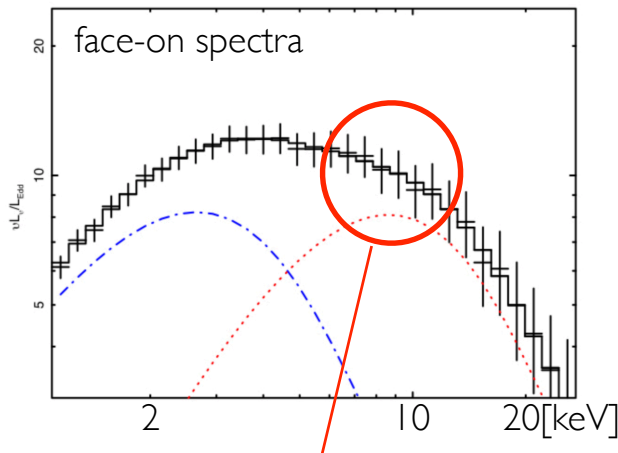
**Simulated spectra nicely fit the observations.**

*Kawashima et al. 2012*  
(data; Gladstone 2009)



# X-RAY SPECTRA

Kitaki, Mineshige, Kawashima, Ohsuga+ 2017



We find roll over  
at around several keV.

## fitting parameters

Model	Parameter (unit)	numerical values
DISKBB	$T_{\text{in}}$ (keV)	1.0
DISKPBB	$T_{\text{in}}$ (keV)	3.1
	$p$	0.86
SIMPL	$\Gamma$	2.9

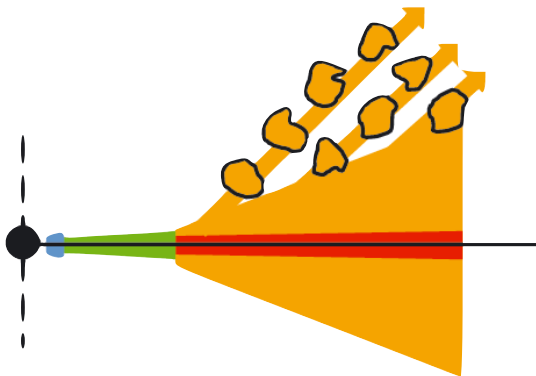
Resulting  $\Gamma \sim 2.9$  is consistent  
with the observations of  
 $\Gamma \sim 3.1$  (Walton+15).

\* $T_{\text{in}}$  is too high. But it would  
decrease when we employ larger  
computational box.

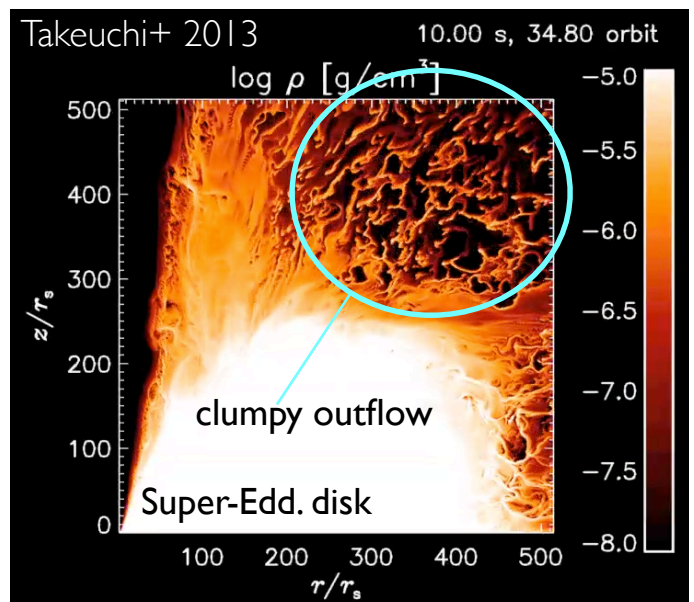
**Our results are consistent with the recent observations.**

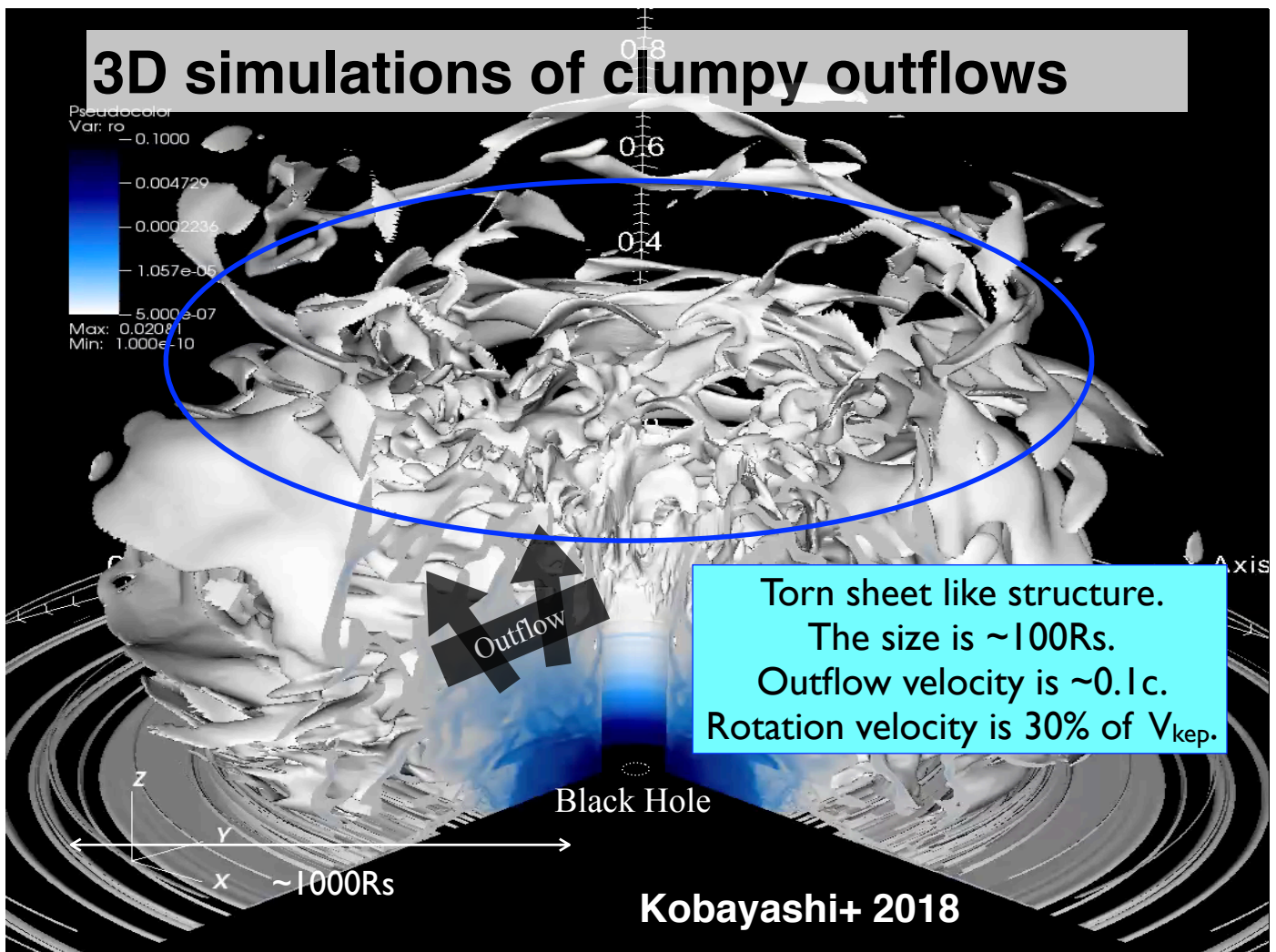
# CLUMPY OUTFLOWS

Some ULXs exhibit the time  
variations of X-ray luminosity,  
implying the launching of clumpy  
outflows (Middleton+2011).



2D simulations reveals that the  
outflows fragments into many  
clouds via RT instability.





## Comparison with Observations

### Sheet like structure

Size (azimuthal direction)  $\sim 100R_s$

Outflow velocity  $\sim 0.1-0.2c$

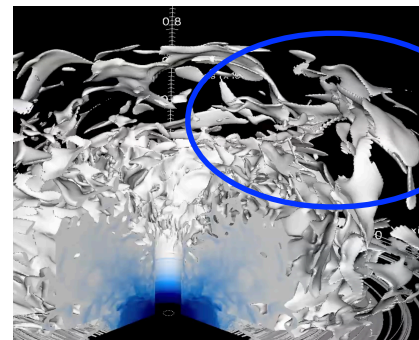
Rotation velocity  $\sim 30\%$  of  $V_{\text{kep}}$

### (1)Time variation

Timescale of the luminosity variation  
( $100R_s/0.3V_{\text{kep}}$ ) is

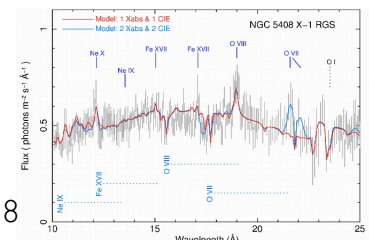
$$\sim 2.5 \left( \frac{M_{\text{BH}}}{10 M_{\odot}} \right) \left( \frac{\ell_{\text{cl}}^{\theta}}{10^2 r_s} \right) \left( \frac{r}{10^3 r_s} \right) s$$

consistent with the observations  
(Middleton+11) in the case of  
 $M_{\text{BH}} \sim 10-100 M_{\text{sun}}$ .



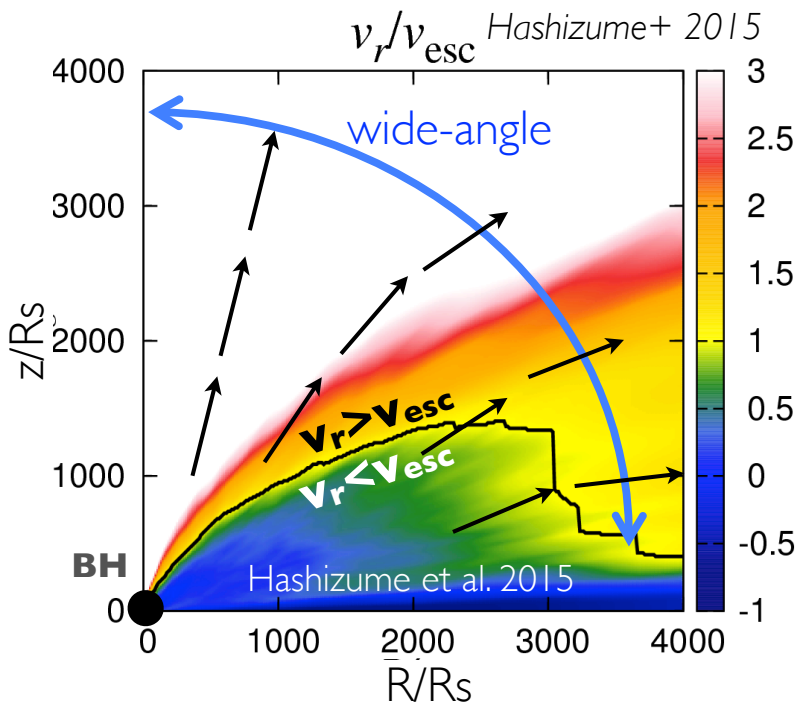
### (2)Absorption lines

Outflow velocity of  $\sim 0.1-0.2c$   
agrees with the observations  
of blueshifted absorption lines.



Pinto+16,  
see also Kosec+18

# ULX BUBBLE



At the direction of  $45^\circ$ , we find  $v_r > v_{esc}$  near the black hole.

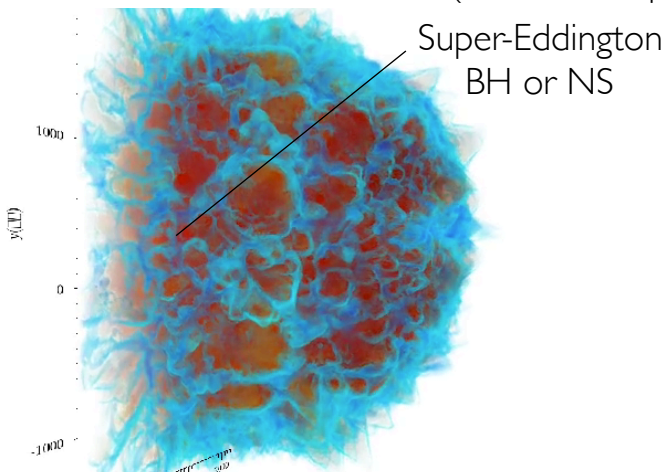
In  $\theta \sim 45^\circ - 80^\circ$ , outflow velocity gradually increases and exceeds  $v_{esc}$  at  $r \sim 1000 - 4000 R_s$ .

Such a wide angle outflow is accelerated by radiation force.

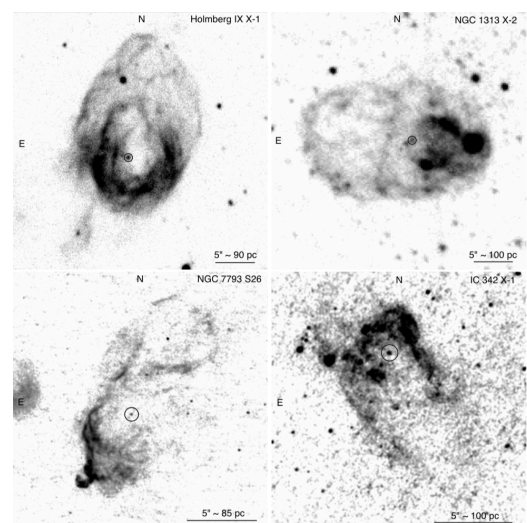
# ULX BUBBLE

Simulated kinetic power is comparable to the X-ray luminosity ( $L_{kin} \sim L_X \sim 10^{39-40} \text{ erg/s}$ ). Such feature agrees with the observed ULX bubbles.

Simulations of ULX Bubbles (Asahina+ in prep.)

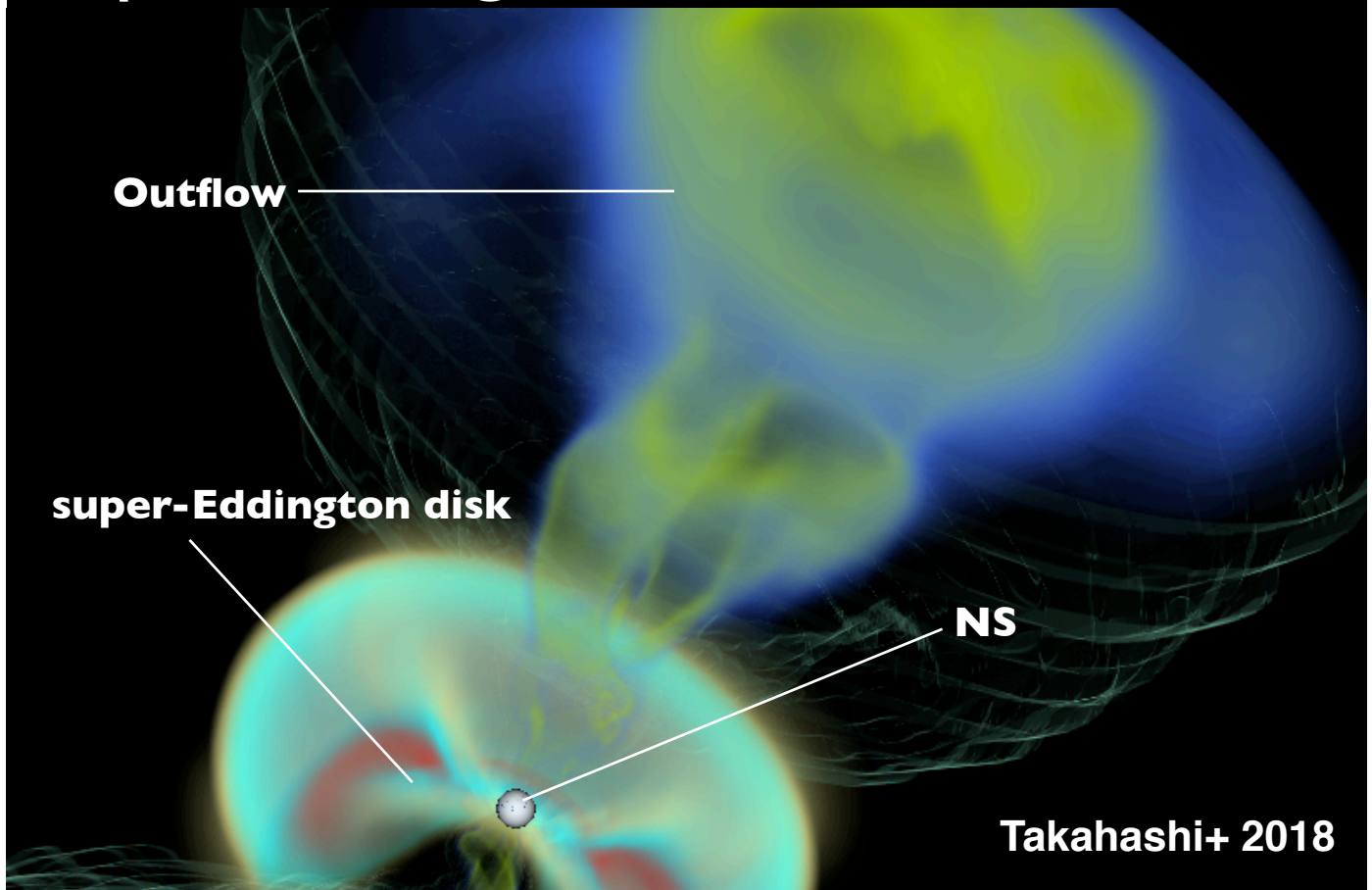


Feng & Soria 11

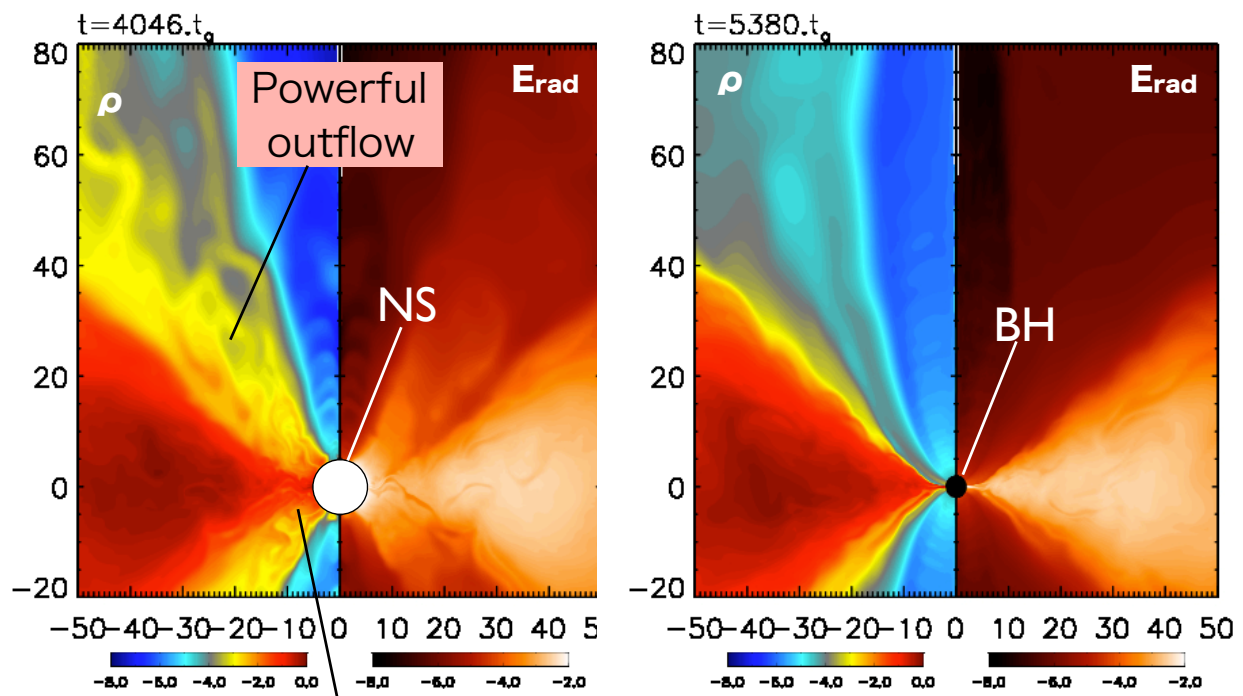


also Pakull & Mirioni (2003), Grise et al. (2006), Pakull et al. (2010), Soria et al. (2010) Cseh et al. (2015)

# Super-Eddington Flows around NSs



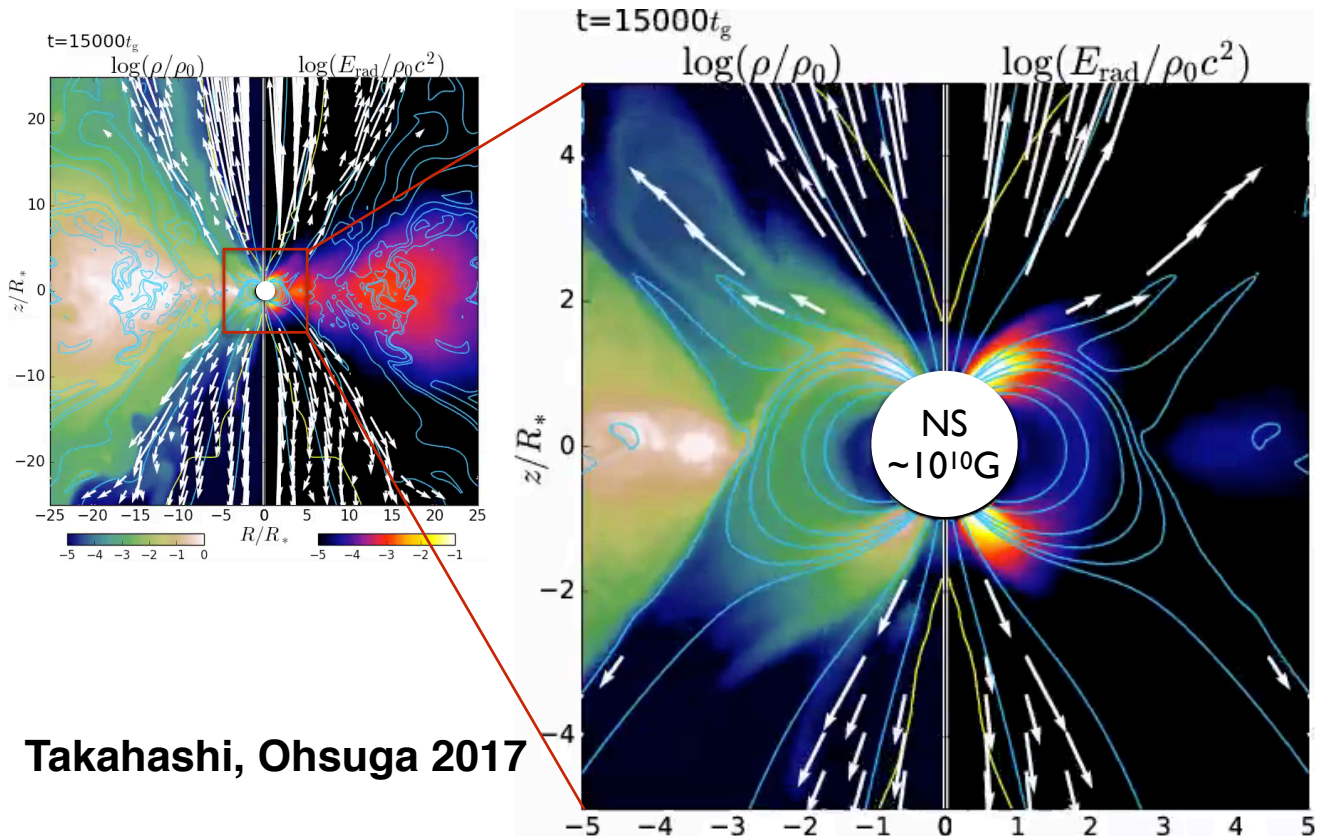
## Non-magnetized NS



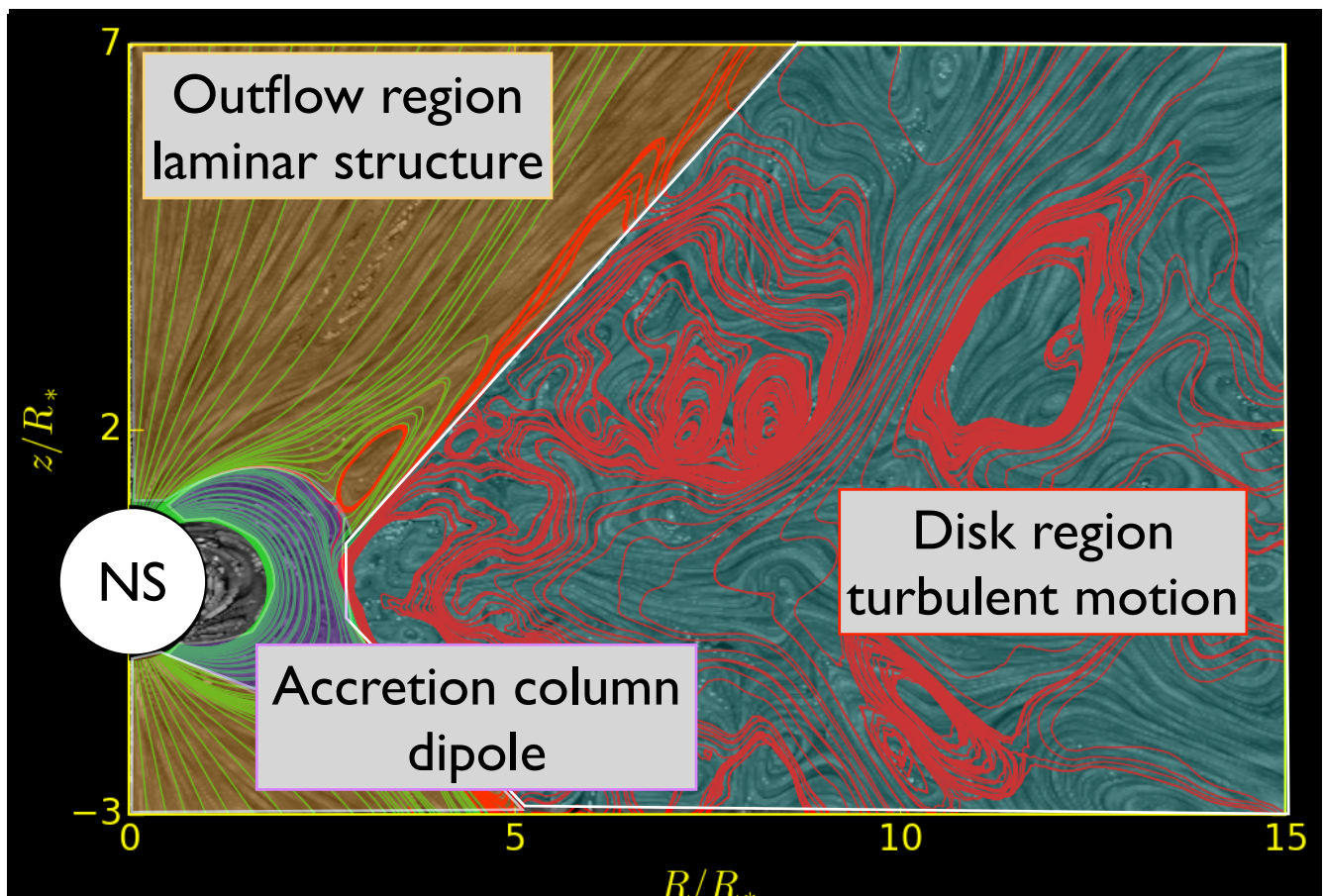
High-density region around the NS surface.

see also Ohsuga 2007

# Magnetized NS



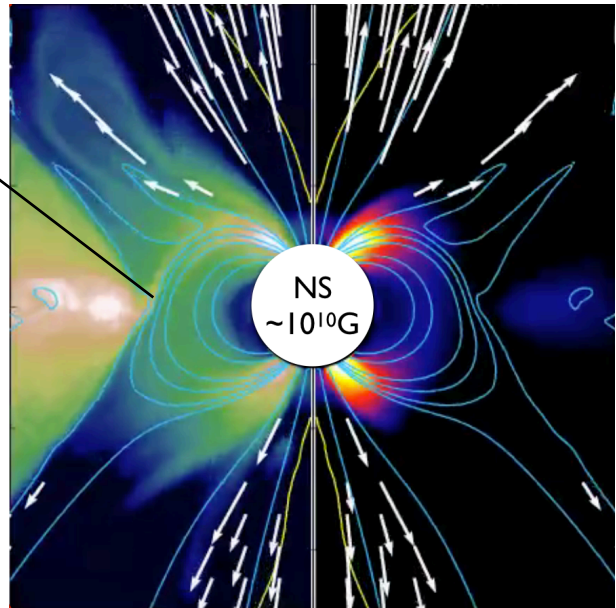
## Flow structure



# Spin-up rate

Magnetospheric radius ( $R_M$ ) is determined by  **$P_{\text{mag}} \sim P_{\text{rad}}$** , since radiation pressure-dominated disk is truncated.

Gas loses angular momentum at around  $R_M$ , leading to the spin-up of NS. Estimated spin-up rate is about  **$-3 \times 10^{-11} \text{ s}^{-1}$** . It roughly agrees with the observations of ULXPs.



## SUMMARY

- Super-Eddington flows can explain the basic features of the ULXs (X-ray luminosity, spectra, clumpy outflow, ULX bubbles).
- Super-Eddington accretion onto NS is more powerful engine, since the energy as well as the matter is not swallowed.
- ULXPs would be powered by the super-Eddington flows onto magnetized NS (super-Eddington accretion column).