Ion and neutral gas escape from the terrestrial planets

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	Н	0	02
Earth	0.6	9.7	19.4
Venus	0.5	8.2	16.4
Mars	0.13	1.8	3.6







What is the role of the intrinsic magnetic field?



The present paradigm: a magnetosphere created by an intrinsic magnetic fields protects the planetary atmosphere from the erosion resulted from the solar wind. The magnetospheric fields separate the moving interplanetary magnetic field "frozen-in" in the solar wind and the upper ionized part of the atmosphere, ionosphere. The ionospheric ions are not subjected to the induced and convective electric fields and cannot be accelerated to energies above the escape energy and leave the planet. Contrary, atmospheres of non-magnetized Mars and Venus are subject to the erosion because the interplanetary magnetic field can reach their ionospheres.











Escape rates from Venus





 Escape at the Earth through the polar region is 2 times higher than the total escape: 2 / 3 return to the ionosphere (Seki's et al., 2001).



Escape rate, 1/s	Reference	Comment
20.10^{24}	Nilsson et al., 2009	Cluster, solar maximum, limited coverage,
	Armenting of al 1000	Only outflow through cusp
45·10 ²⁴	Arvellus el al., 1989	sphere
5.10 ²⁴	Seki et al., 2001	Geotail, solar minimum, limited coverage,
5 10		gross balance (outflow – inflow)



Why are the ion escape rates about the same or even higher from the Earth?





The ion escape rate is determined by the total energy transferred from the solar wind into the ionosphere and is thus limited by the transfer efficiency and the size of the interaction region.

The intrinsic magnetic field does not protect planetary atmospheres



Transfer (coupling) function

 $k = \frac{P}{(A \rho V^{3\gamma})}$

P Total escape power (~Q·E)
A interaction area
V solar wind velocity
ρ solar wind density

$$Q = k A \left(\frac{\rho V^3}{E}\right)$$



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Size does matter!



Distribution of the planetary ions



Example 1. Escape dependence on the SW parameters



Ramstad et al., 2015

- The highest escape rate $\Phi = 5.6 \times 10^{24} \text{ s}^{-1}$ is observed for thin (0.1-0.5 cm⁻³) and moderate or high solar wind velocity
 - Largest interaction area
 - Strongest coupling



- Constrained solar wind and EUV flux
- Nominal conditions only to achieve the best statistics
- 8 years of observations 2007-2015





$$A_{\text{interaction}} \sim (\rho V^2 \cos^2(SZA))^{-1}$$

- Small SZA \rightarrow small A_{interaction} \rightarrow low escape rate
- Large SZA \rightarrow IMB above the anomaly region \rightarrow small/"no" A_{interaction} \rightarrow low escape





Induced magnetospheres protect Mars and Venus even better than the terrestrial magnetosphere.

But how about the total loss?







Total atmospheric loss (2). Extrapolation over 4 Ga

Time before the present [Ga]

 10^{10} s^{-1}] Solar wind flux From observations a 10^{9} $v_{
m sw}~[
m km/s]~F_{
m sw}~[
m cm^{-2}$ of Sun-like stars 10^{8} $\circ \pi^1$ UMa 10^{7} Solar wind velocity b) 1000 From theory 500 Solar wind density $n_{
m sw}~[{
m cm}^{-3}]$ c) 10^{2} $N_{sw} = F_{sw}/V_{sw}$ 10 10 Solar wind dynamic pressure 10^{2} d) $F_{\rm XUV} \ [{\rm m^{-2}\ s^{-1}}] \ p_{\rm dyn} \ [{\rm nPa}]$ $P_{dyn} = \frac{1}{2} N_{sw} V_{sw}^{2}$ 10^{1} 10^{0} 10^{-1} Solar photoionizing flux e) From observations 10^{15} of Sun-like stars 10^{14} Mars ion escape rate f) From MEX $Q \left[\mathrm{s}^{-1} \right]$ 10^{25} measurements 10^{24} $p_{loss} = \frac{m_{avr} g}{A} \int Q(t) dt$ $\Delta p_{
m surf} \, [
m mbar]$ 10 Mars equivalent surface pressure lost $\Delta P = 6 mbar$ Pre-Hesp. Noa. Amazonian Noa. 3 2 0

(Ramstad et al., 2018)



The total atmospheric loss for 4 Ga is < 10 mbar. Negligible!



Is it the end? No! How about the neutral escape?



 Atoms/molecules can gain energy above the escape energy in dissociative recombination reactions and elastic collisions building up extended coronae

$$\begin{split} &O_2{}^+ + e \Box {\rightarrow} O + O + (0.8...7.0 \text{ eV}) \\ &CO_2{}^+ e {\rightarrow} CO + O + 8.3 \text{ eV} \\ &H_{th} + O_{hot} {\rightarrow} H_{hot} + O_{hot} \end{split}$$

• Only a few eV can be gained. This escape channel is important for Mars only.



Wang et al., 2015 (models)



Sputtering

- Momentum transfer from precipitating (mostly) O⁺ ions
- Atmospheric sputtering results in the additional population of the extended hot corona and in the increase of the Jean escape, especially at high levels of solar activity.
- The only channel to remove carbon in the form of CO₂ and CO.



Comparison DR (Dissociative Recombination) and sputtering: 1D models at Solar Max Shematovich, 2006



• Jean escape from hot coronae and sputtering

$$\Phi_{J} = \frac{N_{ex}v_{o}}{2\sqrt{\pi}} (1 + \lambda_{esc})e^{-\lambda_{esc}} \qquad \qquad \lambda_{esc} = \frac{GMm}{kT(R+z)} = \frac{R+z}{H(z)} = \left(\frac{v_{e}}{v_{o}}\right)^{2}$$

$$v_{e} = \sqrt{\frac{2GM}{r}}$$

$$v_{o} = \sqrt{\frac{2KT/m}{r}}$$

- Measurements of the hot coronae profiles (N_{exo} , T)
 - Difficult due to very low densities. Extensive modeling involved
 - Remote sensing via UV scattered and stimulated emissions (*Dieghan et al., 2015; Chaffin et al., 2015*)
 - Indirect measurements via pick-up H⁺ and O⁺ flux tracing (*Rahmati et al., 2017*)

Hot H : $3 \cdot 10^{25} \text{ s}^{-1}$ (measured, Rahmati et al., 2017)Hot O : $7 \cdot 10^{25} \text{ s}^{-1}$ (measured, Rahmati et al., 2017)Sputtered CO: $(1-4) \cdot 10^{20} \text{ s}^{-1}$ (model, Wang et al., 2015)Sputtered CO2: $(9-20) \cdot 10^{21} \text{ s}^{-1}$ (model, Wang et al., 2015)





- Direct measurements are very difficult but possible in the future
- New generation of ENA instruments at Mars required



Upstream solar wind → neutral solar wind

$$F_{neutral} = N_{sw} V_{sw} \sigma \int_{r}^{\infty} n(l) dl = (0.1 - 1\%) F_{sw}$$

S/C CX neutral solar wind (0.1 - 1%)
Solar wind



- Paradigm change: The intrinsic magnetic field does not protect planetary atmospheres
- The atmospheric loss associated with ion escape is low (< 10 mbar) for Mars (and likely for Venus)
- The neutral escape is by far more important channel
- New measurements techniques based on ENA measurements can be applied to better characterized the neutral escape