

# X-Ray Spectrum Analysis I. Low-resolution Spectra

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## Goal of this presentation

Providing you with a reference for topics relevant to spectroscopy of low-resolution (*i.e.* CCD) spectra:

- How do we fit spectra?
  - [and, by the way, what does it mean “*fitting a spectrum*”?]
- What files do we need? what are they?
- How do we turn the fitting wheel?
- How do we deal with calibration uncertainties?

If I make things too messy, *no panic!* Look at (*e.g.*):

<http://heasarc.gsfc.nasa.gov/docs/xanadu/xspec/manual/XspecSpectralFitting.html>



## Outline

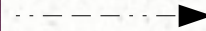
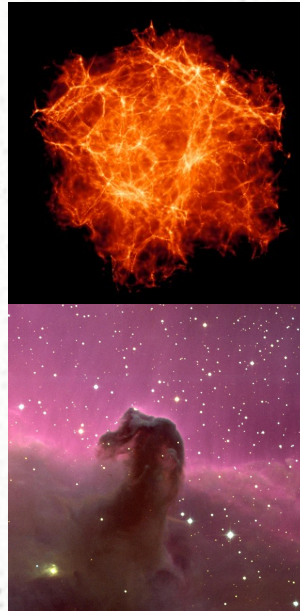
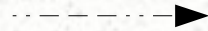
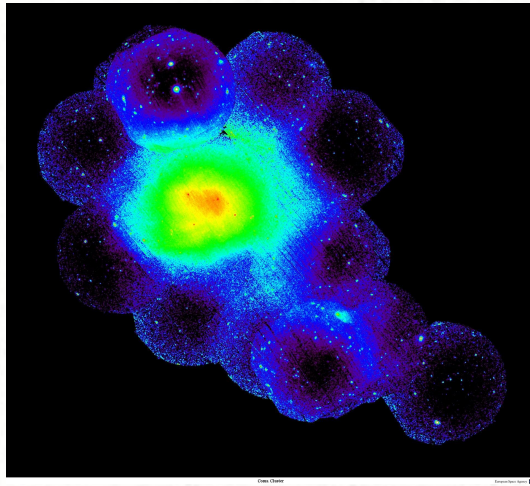
- This talk is primarily intended for users of CCD spectra:
  - ASCA/SIS
  - *Chandra*/ACIS
  - Swift/XRT
  - *Suzaku*/XIS
  - XMM-Newton/EPIC (-MOS and -pn)
- However, some basic principles can be applied to instruments with even lower resolution:
  - ROSAT/PSPC, ASCA/GIS, BeppoSAX, RXTE, *Suzaku*/HXD, NuSTAR ...
- [Bias: *Many examples refer to the EPIC cameras, but only because I am paid to ensure that they work ...*]

# Our ultimate goal is ...

Intrinsic source spectrum  $s(E)$  ...

... seen through IGM/ISM absorption  $a(E)$  ...

... detected as observed counts  $C(\text{PHA})$

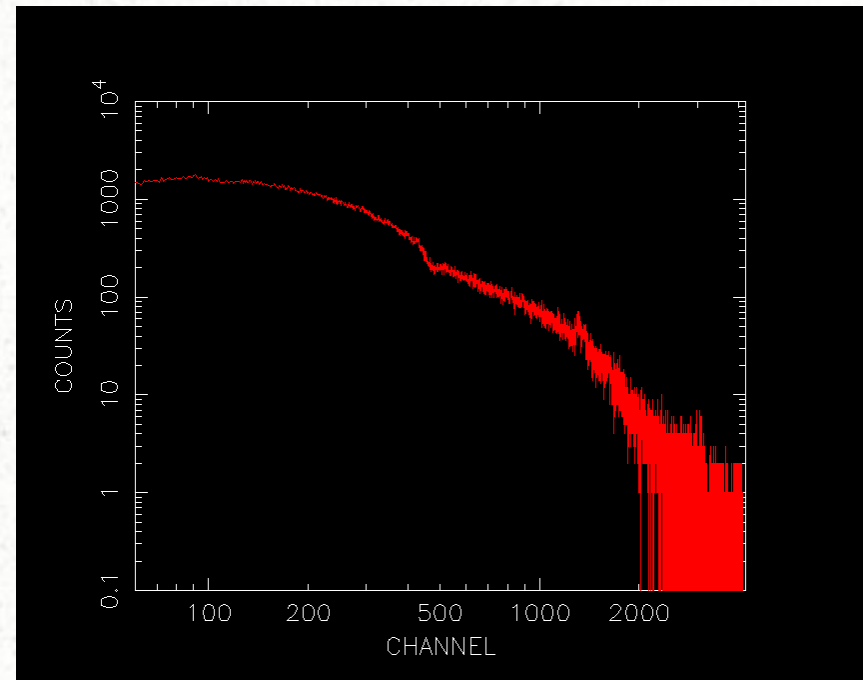
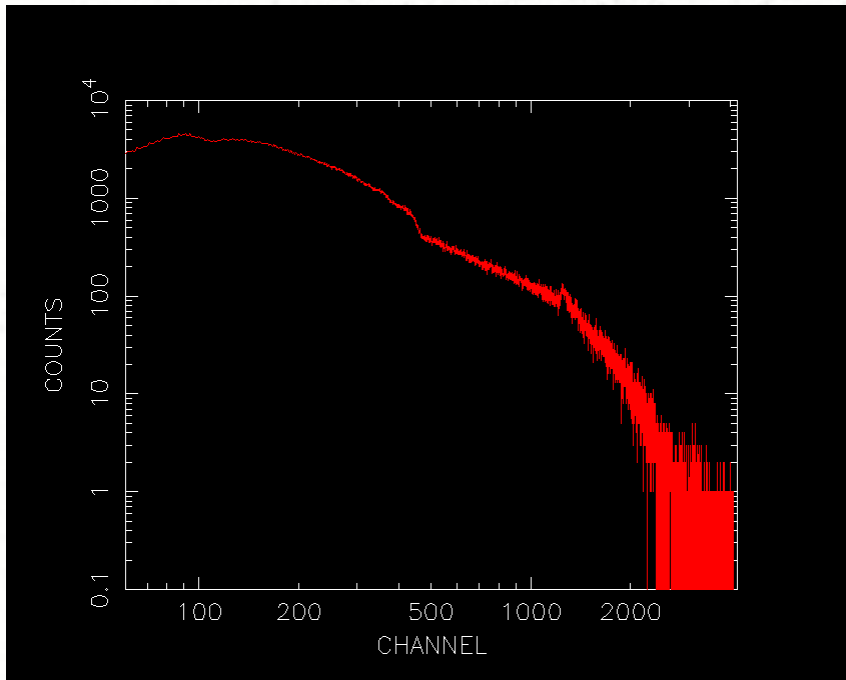


We measure  $C(\text{PHA})$ . We want to determine  $S(E)$  - occasionally  $A(E)$ . Easy, isn't it?



*When all candles be out, all **cats** are grey*

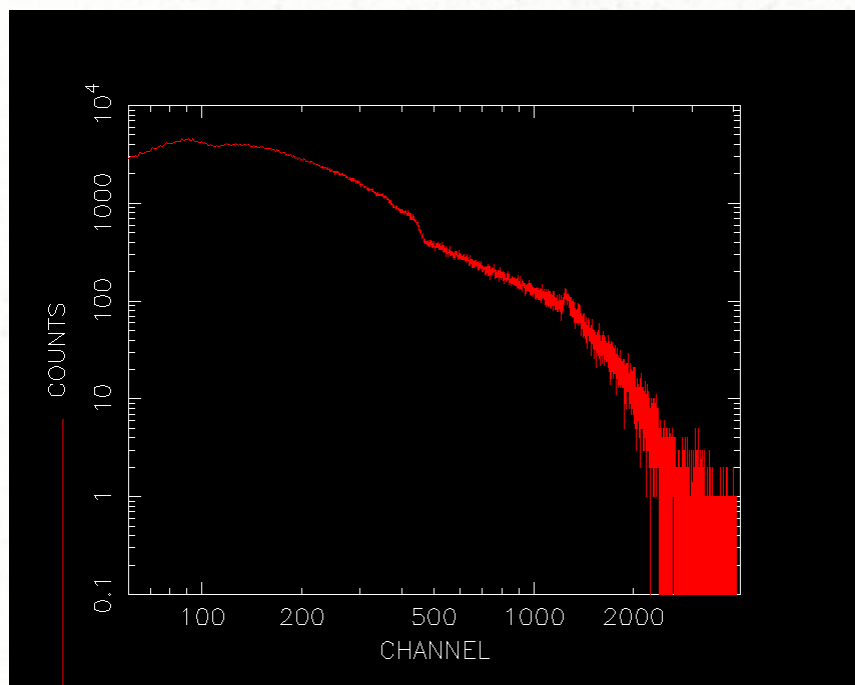
CCD spectra extracted by **dmextract**, **xmm/evselect**, or **xselect** look like this:



## *When all candles be out, all **cats** are grey*

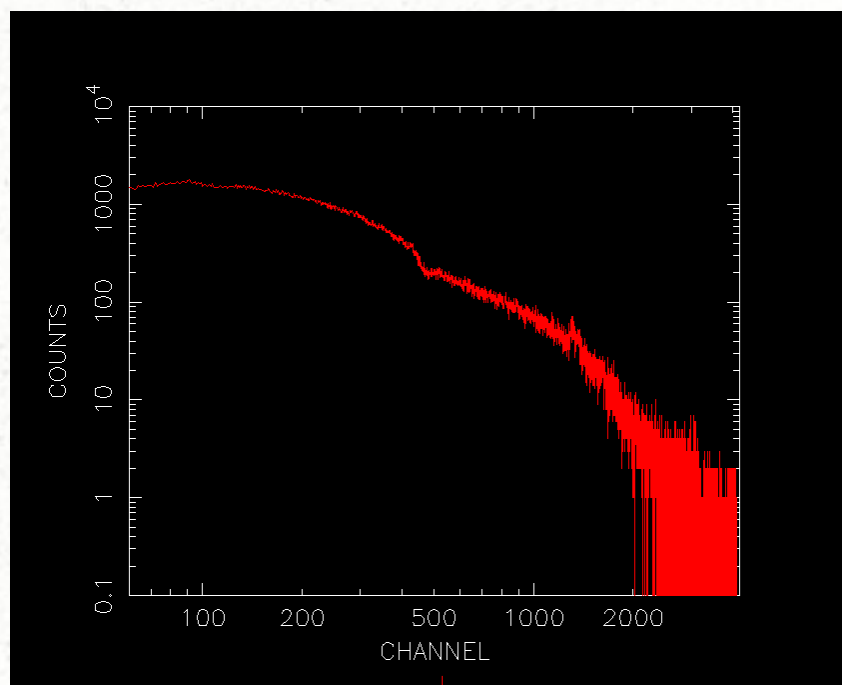
CCD spectra extracted by **dmextract**, **xmm/evselect**, or **xselect** look like this:

Ark120 – EPIC-pn (AGN)



These are “**COUNTS per bin**”, not flux!

Coma – EPIC-pn (Galaxy Cluster)



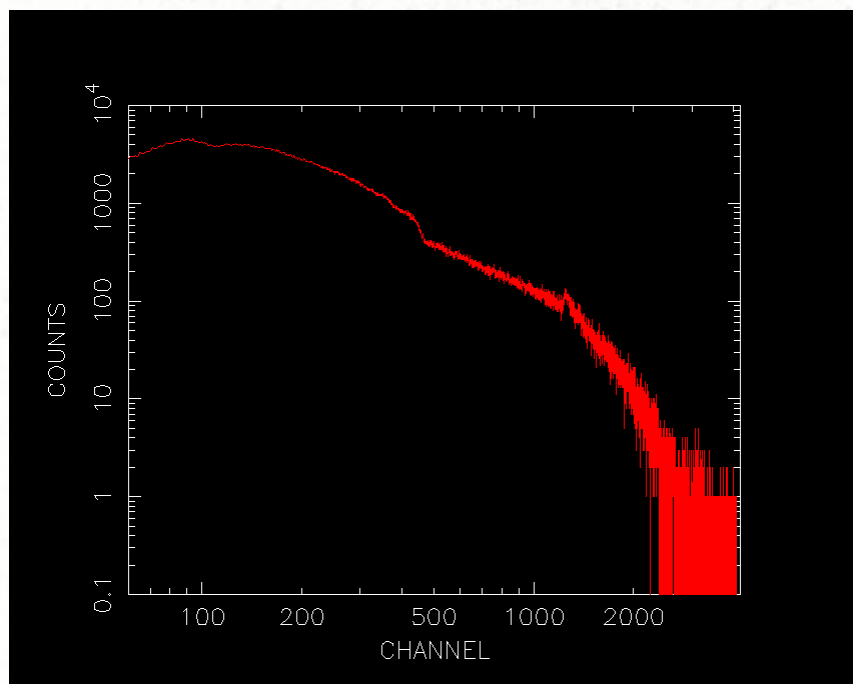
These are “**CHANNELS**”, not energy!

First problem: spectral extractors produce spectra in instrumental quantities

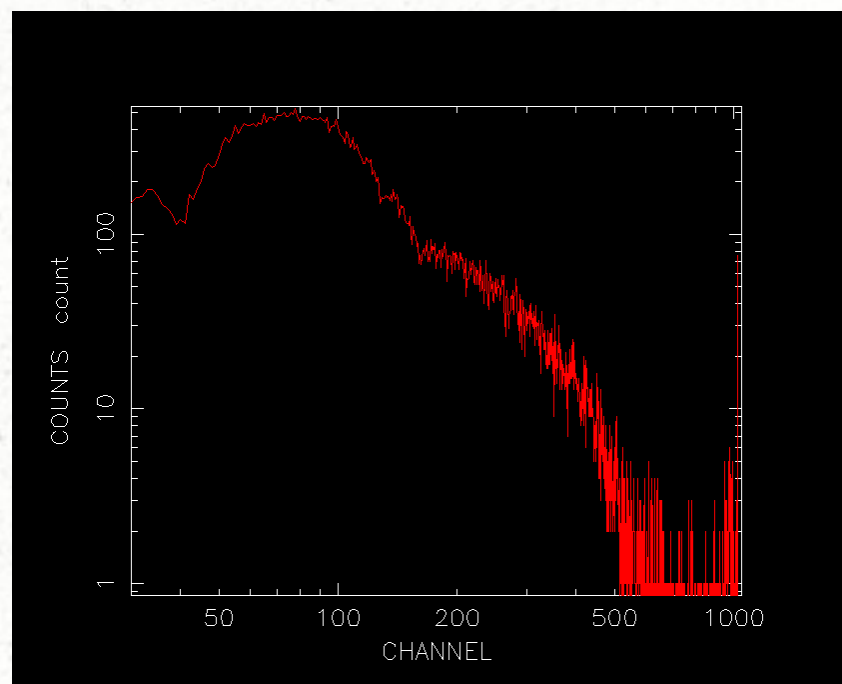
*When all candles be out, all **cats** are grey*

*“And now, for something completely different: the larch ...” (Monty Python, 1968)*

Ark120 – EPIC-pn (AGN)



Ark120 – SIS (AGN)



Second problem: the shape of the count spectra is dominated by the transfer function of the telescope+detector: we must “decode” it



## The spectral equation

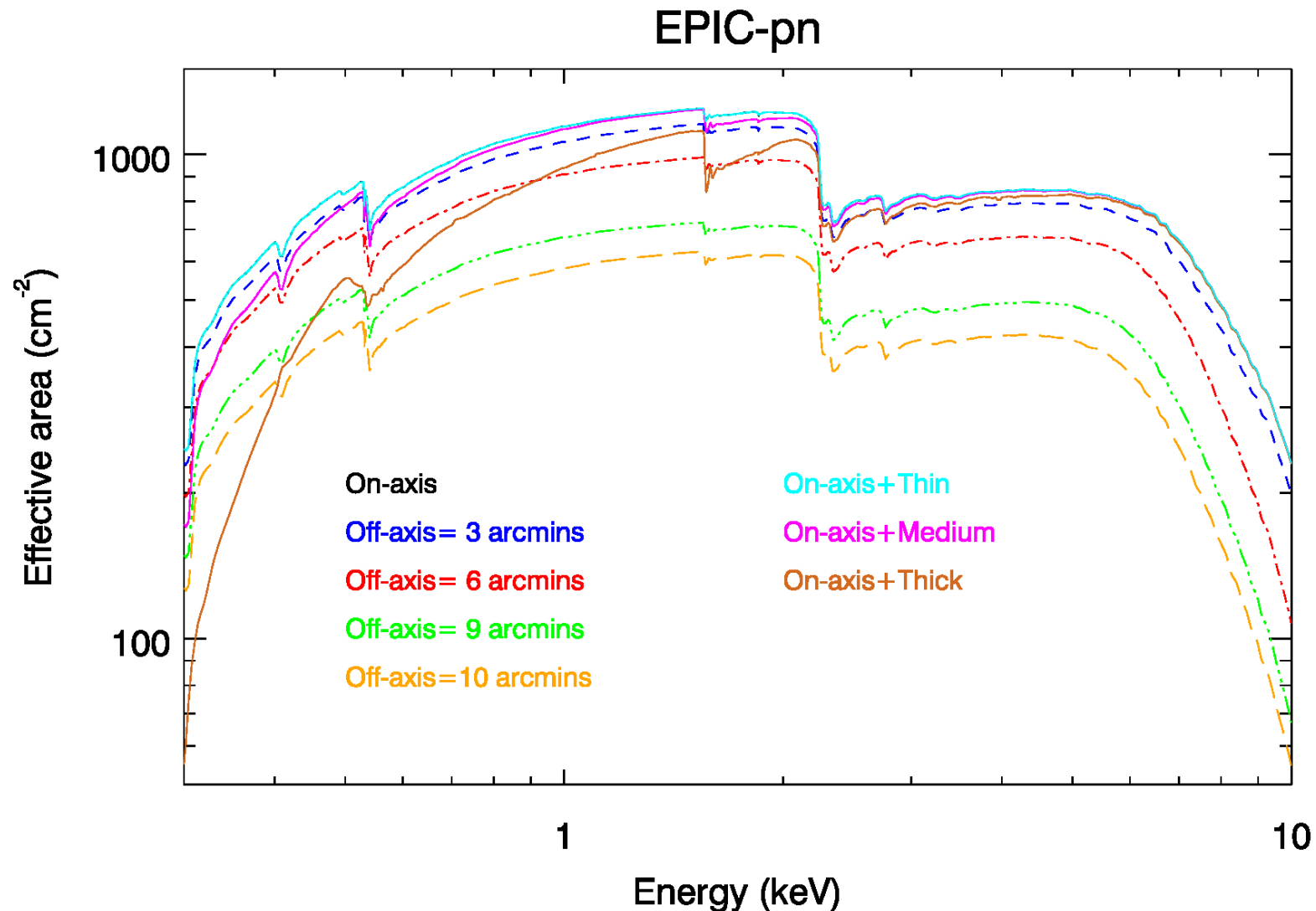
$$C(h) = (N\tau) \int dE R(h, E) A(E) s(E)$$

- $(N\tau)$  = exposure time
- $C(h)$  = observed spectrum, in units of *counts per spectral bin*
- $R(h,E)$  = redistribution matrix (a.k.a. “RMF file”), typically normalised to 1
- $A(E)$  = effective area (a.k.a. “ARF” or “ancillary file”) in units of *area*
- $s(E)$  = intrinsic spectrum (to be determined)
- $h$  = spectral channels, in units of *Pulse Height Analysis* (PHA) or *Pulse Invariant* (PI): digital instrumental quantities only loosely related to energy

We would need to invert this equation to get  $s(E)$ .  
However, in general this is not possible. Why?



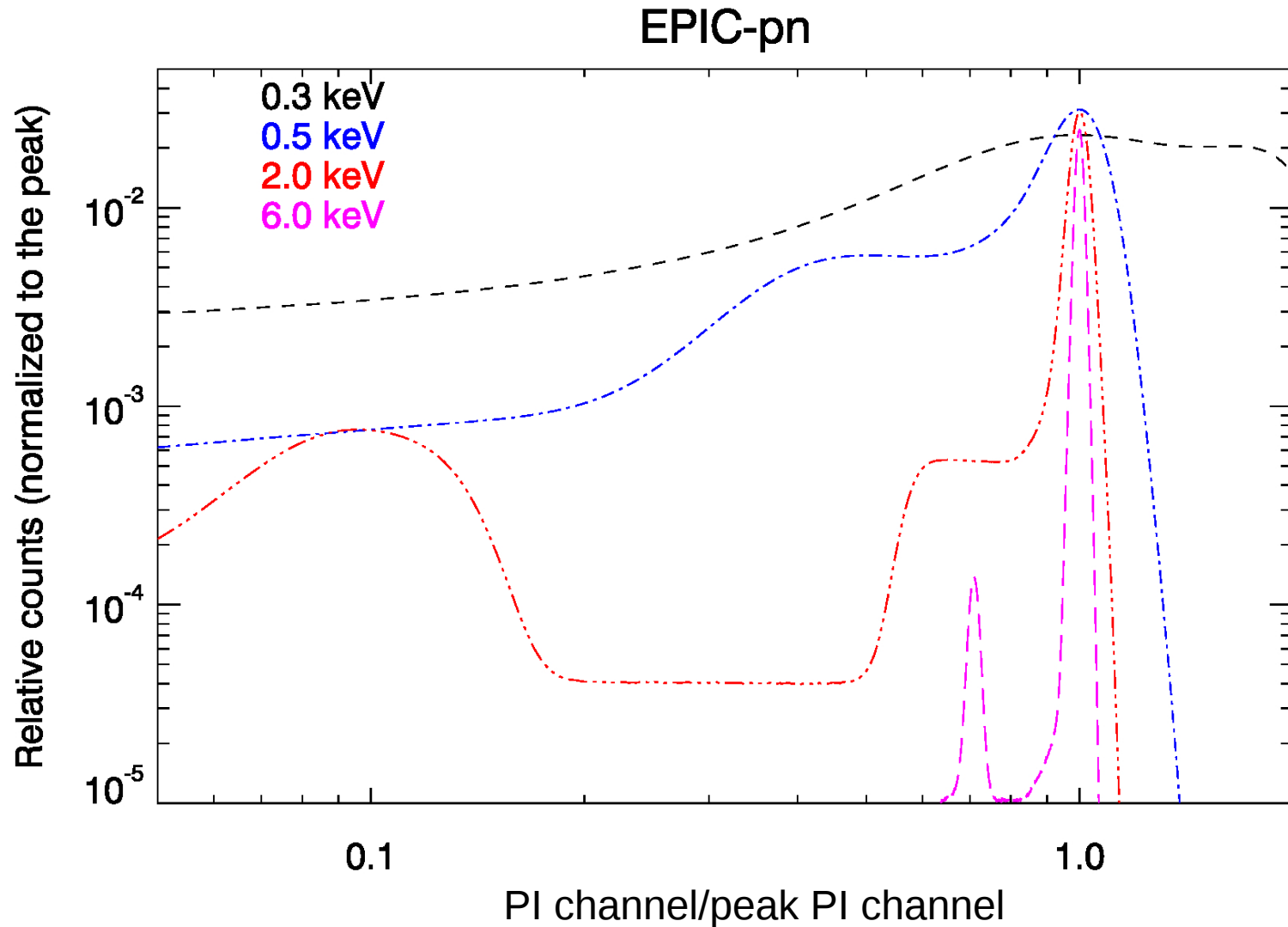
## The effective area $A(E)$



Measure (conventionally expressed in units of “area”) of the collecting power of telescope+filter+detector. It depends on energy and position (“off-axis”)

[Beware: not all observatories carry “optical photon blocking filters”]

## Redistribution matrix $R(E)$



Response of the detector to a monochromatic line. Highly dependent on the energy. The width of the core defines the instrument resolution:

$$\sigma_{PHA} = [n^2 + fE]^{0.5} \quad (n^2 \rightarrow \text{noise term})$$

## *Inverting the spectral equation?*

The redistribution is sampled at discrete spectral channels:

$$R_{hE}^i = \frac{\int_{E_{j-1}}^{E_j} R(i, E') dE'}{(E_j - E_{j-1})}$$

The whole spectra matrix is actually a discrete matrix equation:

$$C_h = T \sum_i \sum_E R_{hE}^i A_E^i S_E^i dE$$

The cross-talk among different energies prevents the  $R_{hE}^i$  matrix from being inverted.

Alternative: **Forward-folding approach**



## *Forward-folding approach*

- 1) Assume a model with its defining parameters
- 2) Define a set of parameter values
- 3) Convolve the model with the instrument response
- 4) Compare the (dis)agreement between the observed spectrum and the folded model through a *goodness-of-fit* statistical test
- 5) Change the parameter values to minimize the goodness-of-fitness test  $\equiv$  **fit**
- 6) Once the best-fit is found, calculate the confidence intervals on the best-fit parameters

Spectral packages are looping machines through the steps above (+ a few other cosmetic features)

# Background spectra

The *inevitable background* is due to various component:

- Space environment
- Instrument
- Astrophysical sources

## Synopsis of background components in XMM-Newton EPIC

	SOFT PROTONS	INTERNAL (cosmic-ray induced)	ELECTRONIC NOISE	HARD X-RAYS	SOFT X-RAYS
Source	Few x 100 keV solar protons, accelerated by magnetospheric reconnection events. Dominate times of high-BG.	Interaction of High Energy particles (cosmic rays) with detector - associated instrumental fluorescence. <a href="#">Main MOS ref.</a>	(1) Bright pixels & (parts of) columns. (2) CAMEX readout noise (pn). (3) (4) (5) (6) Artificial Low-E enhancements in outer MOS CCDs (Also dark current - thought negligible).	X-ray background (AGN etc). <a href="#">Single Reflections from outside FOV</a> . <a href="#">Out-of-time (OOT) events (pn)</a>	Local Bubble, Galactic Disk, Galactic Halo, <a href="#">Solar Wind Charge Exchange (SWCX)</a> <a href="#">SWCX</a> , <a href="#">Single Reflections from outside FOV</a> , <a href="#">Out-of-time (OOT) events (pn)</a>
Variable? (per Observation)	Flares (up to >1000%). Unpredictable. Significant quiescent component (long flares) - survive GTI screening. ( <a href="#">Also additional possible 'irreducible' component</a> ).	+/-10%. <a href="#">MOS</a> , <a href="#">MOS</a> : >2keV continuum unchanged, small changes in fluorescence lines. <1.5keV continuum varies - may be due to Al redistribution. <a href="#">pn</a> : Difference between continuum and lines (some correlation).	(1) +/-10%. (2) Very constant. (3) (4) Believed constant.	Constant.	Constant. Long obs. may see effect of <a href="#">SWCX</a> <a href="#">SWCX</a> (e.g. variations at 0.5-1.2 keV [OvIII/MgXI], but not at 2-4 keV).
Variable? (Obs. to Obs.)	Unpredictable. Affect 30%-40% of time. Flaring SP increasing? Quiescent SP not evolving. More SPs far from apogee. More SPs in winter than in summer. Low-E flares turn on before high-E.	<a href="#">Majority @ +/-15%. Can be x10 higher in high radiation periods</a> . No increase after solar flares. Plus above 'per Observation' variations.	(1) >1000% (pixels come and go, also [micro-meteorite damage]). (2) Mode-dependent (lowest eFF, then FF, LW, highest SW) (3) effects 5-20+% of obs. (4) effects 20-50% of obs. (factor increases with high-BG rate). (5) (6) >50% of obs for later Revs (Rev1300+)	Constant. <a href="#">OOT</a> events (pn) mode-dependent (LW:0.16%, FF:6.3%, eFF:2.3%)	Variation with RA/Dec (+/-35%). <a href="#">SWCX</a> <a href="#">SWCX</a> may affect observations differently. <a href="#">OOT</a> events (pn) mode-dependent (LW:0.16%, FF:6.3%, eFF:2.3%)
Spectral	Variable. Unpredictable. Continuum spectrum (no lines), fitted by unfolded xspec PL ( <a href="#">double-exponential or broken power law (break energy stable ~3.2 keV)</a> ) model for E>0.5keV (E<0.5keV, less flux is seen). <a href="#">Variable in intensity + shape (higher the intensity, flatter the slope)</a> .	Flat ( <a href="#">MOS index ~0.2</a> ) + fluorescence + detector noise. <a href="#">MOS</a> : 1.5keV Al-K, 1.7keV Si-K, 2.2keV Au. <a href="#">Det. noise &lt;0.5keV</a> . High-E lines (Cr 5.4, Mn 5.8, Fe-K 6.4, Au 9.1&11.4). (Here also) <a href="#">PN</a> : 1.5keV Al-K, No Si (self-absorbed), Cu-Ni-Zn-K (~8keV). MIP noise <0.3keV.	(1) low-E (<300eV), tail may reach higher-E. (2) low-E (<300eV). (3) (4) low-E (<500eV) (3) High-rate plus soft excess. (5) (6) Strong excess <1000eV.	1.4 power law. Below 5keV, dominates over internal component. Above 5keV, internal component dominates (in times of low-BG).	Thermal with ~<1keV emission lines. Extragalactic @>0.8keV, index=1.4. Galactic - emission/absorption varies. <a href="#">SWCX</a> <a href="#">SWCX</a> very soft, with unusual OvIII/OvII line ratios (plus others) - Strong OvIII & MgXI
Spatial - Vignetted?	Yes (scattered) - <a href="#">Vignetting is flatter than for photons - low-E SPs extremely flat, higher-E SPs steeper (MOS1 - pn shows more constant vignetting with energy)</a>	No - flat (see below).	(1,2) Bright pixels and CAMEX - No. MOS noise - (3) No/unclear (out-FOV) (see below) (4) Yes - evident in vignetting maps (in-FOV). (similar, smaller-magnitude vignetting asymmetries seen in pn). (5) (6)	Yes.	Yes.
Spatial - Structure?	Perhaps, in MOS due to the RGA. No structure seen in pn. <a href="#">SP feature seen in MOS1-CCD2 at low-E</a> . SPs observed only inside FOV.	Yes. Detector + construction. <a href="#">MOS</a> : outer CCDs more Al, less Si. <a href="#">CCD edges more Si</a> . Less Si out-FOV. Continuum diff. between out-FOV and in-FOV below Al line (redistribution?). More Au out-FOV. Changes in high-E lines. <a href="#">CCD-to-CCD: line intensity variations, energies/widths stable</a> . (Here also) <a href="#">PN</a> : Line intensities show large spatial variations from electronic board. Central 'hole' in high-E lines (~8keV). Residual MIP contribution near CAMEX readout (low-E, non-singles, parallel to readout).	Yes. (1) Individual pixels & columns. (Also [pn] sections of columns away from CAMEX, near to FOV centre) (2) Near pn readout (CAMEX), perpendicular to readout. (3) MOS1 CCDs 4 & 5, MOS2 CCDs 2 & 5 - unusual in- & out-FOV differences (esp. MOS1 CCD4) and spatial inhomogeneities. (4) MOS1 CCDs 2 & 5. (5) (6). Lower-level ~persistent low-E enhancement in MOS1 CCD2	No. <a href="#">Single reflections</a> : Diffuse flux from 0.4-1.4 deg (out-FOV) is ~7% of in-FOV signal. <a href="#">Effective area of 1 telescope ~3 sq.cm at 20-80 arcminutes off-axis</a> . <a href="#">OOT</a> events (pn) smeared along readout from bright sources of X-rays. (extra BG in pn LW mode due to frame store area).	No, apart from real astronomical objects. Exgal.>0.8keV spatially uniform. <a href="#">SWCX</a> <a href="#">SWCX</a> over whole FOV. <a href="#">Single reflections</a> : Diffuse flux from 0.4-1.4 deg (out-FOV) is ~7% of in-FOV signal. <a href="#">Effective area of 1 telescope ~3 sq.cm at 20-80 arcminutes off-axis</a> . <a href="#">OOT</a> events (pn) smeared along readout from bright sources of X-rays. (extra BG in pn LW mode due to frame store area).
Patterns	Distribution similar to genuine X-rays.	Distribution different from genuine X-rays.	Distribution different from genuine X-rays. (5) MOS E1/E2 connection	Genuine X-ray distribution.	Genuine X-ray distribution.

This implies that some components are focused by the telescope. Others aren't



## How to deal with background spectra

$$C_h = T[\sum_i \sum_E R_{hE}^i A_E^i (s_E^i + b_E^{i,f}) dE + b_E^{i,u}]$$

focused

not focused

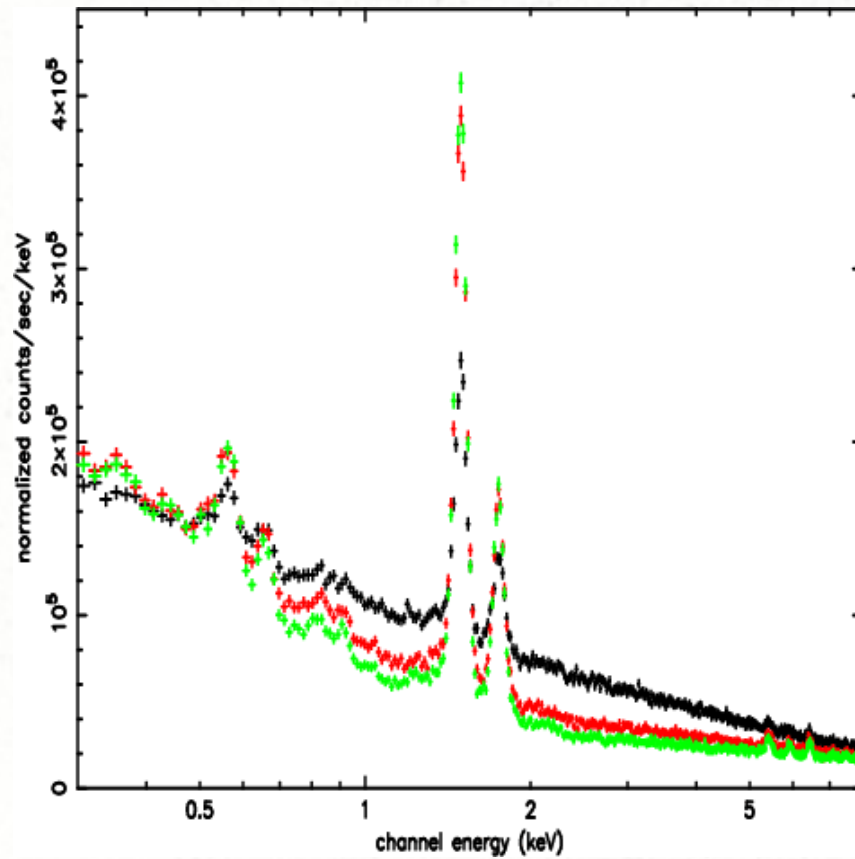
Three approaches are possible:

- Ignore the background. **Wrong**, even if in *Chandra* it is often very low
- Subtract the background. Easy, but:
  - “It reduces the amount of statistical information in the analysis [...]
  - The background subtracted data are not Poisson-distributed;
    - [For example, subtracting a background can give negative counts; this is definitely not Poissonian!]
  - Fluctuations, particularly in the vicinity of localized features, can adversely affect analysis”
- Model and fit simultaneously the source and the background. Appealing, but:
  - The background spectra is often awfully complex, time- and detector-position dependent, sometimes not known at all

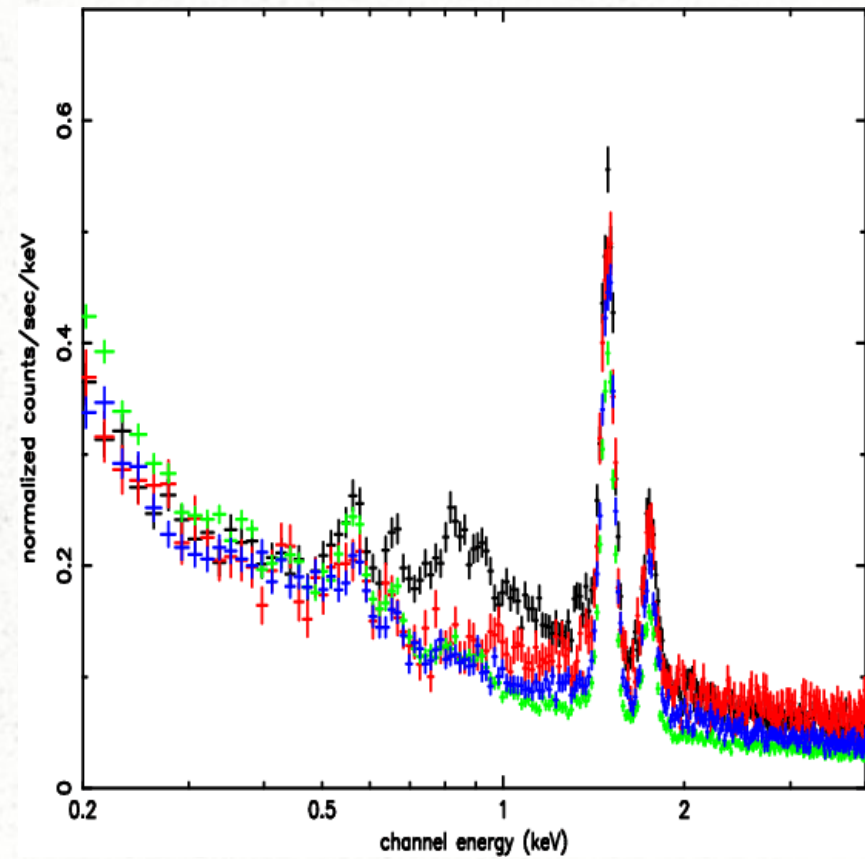


# Goodness-of-fit statistical tests

EPIC-MOS background spectra  
as a function of count rate



EPIC-MOS background spectra  
along different line-of-sights



# Models

Most software packages include the same suite of astrophysical models ( $\sim 10^2$ ):

- Additive:

- Phenomenological: po, bb, brems, gauss
- Astrophysical: comptt, diskbb, apec, diskline

blackbody  
power-law      bremsstrahlung      Gaussian profile

- Multiplicative:

- Absorption, cut-off ...

Comptonization      Thermal plasma  
Accretion disk blackbody      Relativistic line emission

- Convolution:

- Kernels, flux calculation ...

- Mixing

- Surface brightness, deprojection ...

- Colleagues in the community contribute their own (“external model”), either as functions or as FITS table

- You can create your own (it does not require a software guru)!

## *A detour on presentation skills*

Never write a slide with more words than you can read in 30 seconds

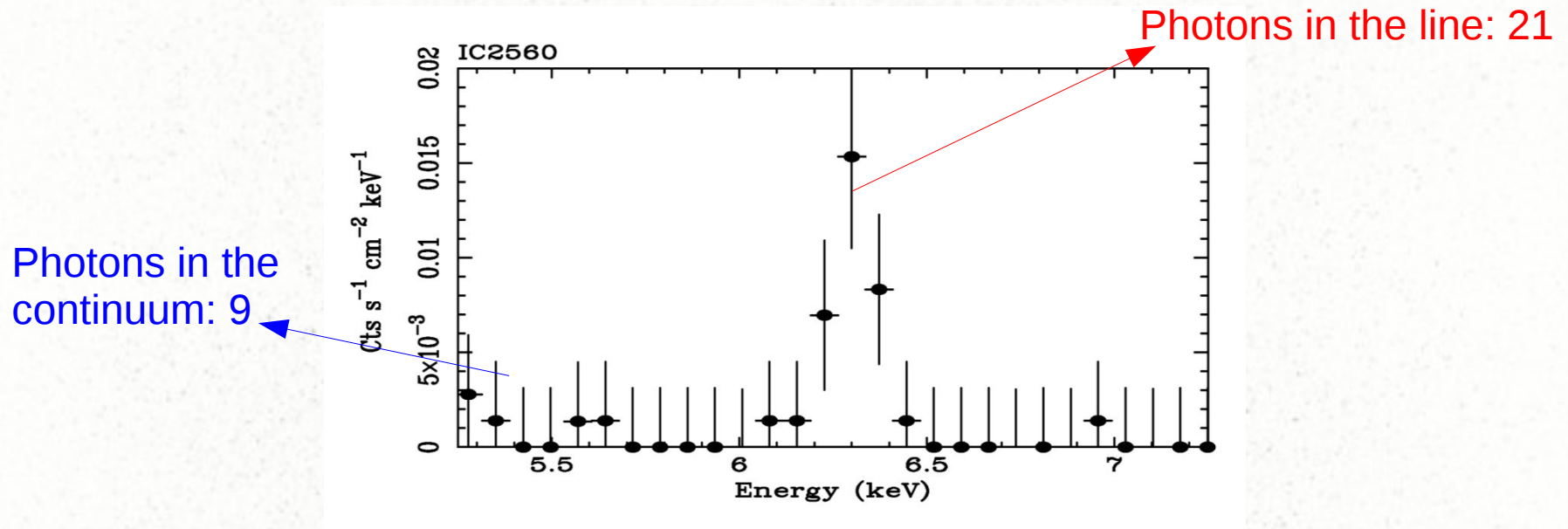


# Goodness-of-fit tests

- The most common goodness-of-fit statistics test is the “chi-squared” ( $\chi^2$ ):  
$$\chi^2 = \sum \frac{(\text{observed} - \text{expected})^2}{\text{expected}}$$
  - It requires that the distribution of background-subtracted counts in each spectral channel is well approximated by a Gaussian (5-10 counts)
  - Different alternatives for the denominator: the **XSPEC default is biased**.  
Use `weight churazov`, or `weight model` instead
- Alternatively, one can use the Cash (C-)statistics  $C = 2 \sum_{i=1}^n s_i - N_i + N_i \ln(N_i/s_i)$ .
  - Applicable to data following the Poissonian statistics only (i.e.: non-background subtracted spectra).
  - XSPEC implements a flavour (the “W-statistics”) which can be directly applied to background-subtracted spectra
  - *[issue with spectra with very low number of counts: K.Arnaud recommends to rebin the spectra to ensure that each channel has got at least one count – reason unknown]*
  - It does not yield a metrics of the absolute quality of a fit (one need to use Monte-Carlo simulations in this case)
- XSPEC version 12.8 allows you to use different statistics to calculate the best-fit parameters, and the absolute quality of the fit.  
Recommendations: C-statistics for the former,  $\chi^2$  for the latter

## To (re)bin or not to (re)bin?

- Rebin your spectra is pure evil, and may lead to loss of scientific information:



- However, a minimum level of spectral rebinning is required to avoid oversampling the intrinsic resolution of the instrument

# Shannon theorem

Let  $f(t)$  be a continuous signal. Let  $\tilde{g}(\omega)$  be its Fourier transform, given by

$$g(\omega) = \int_{-\infty}^{\infty} e^{i\omega t} f(t) dt. \quad (1.6)$$

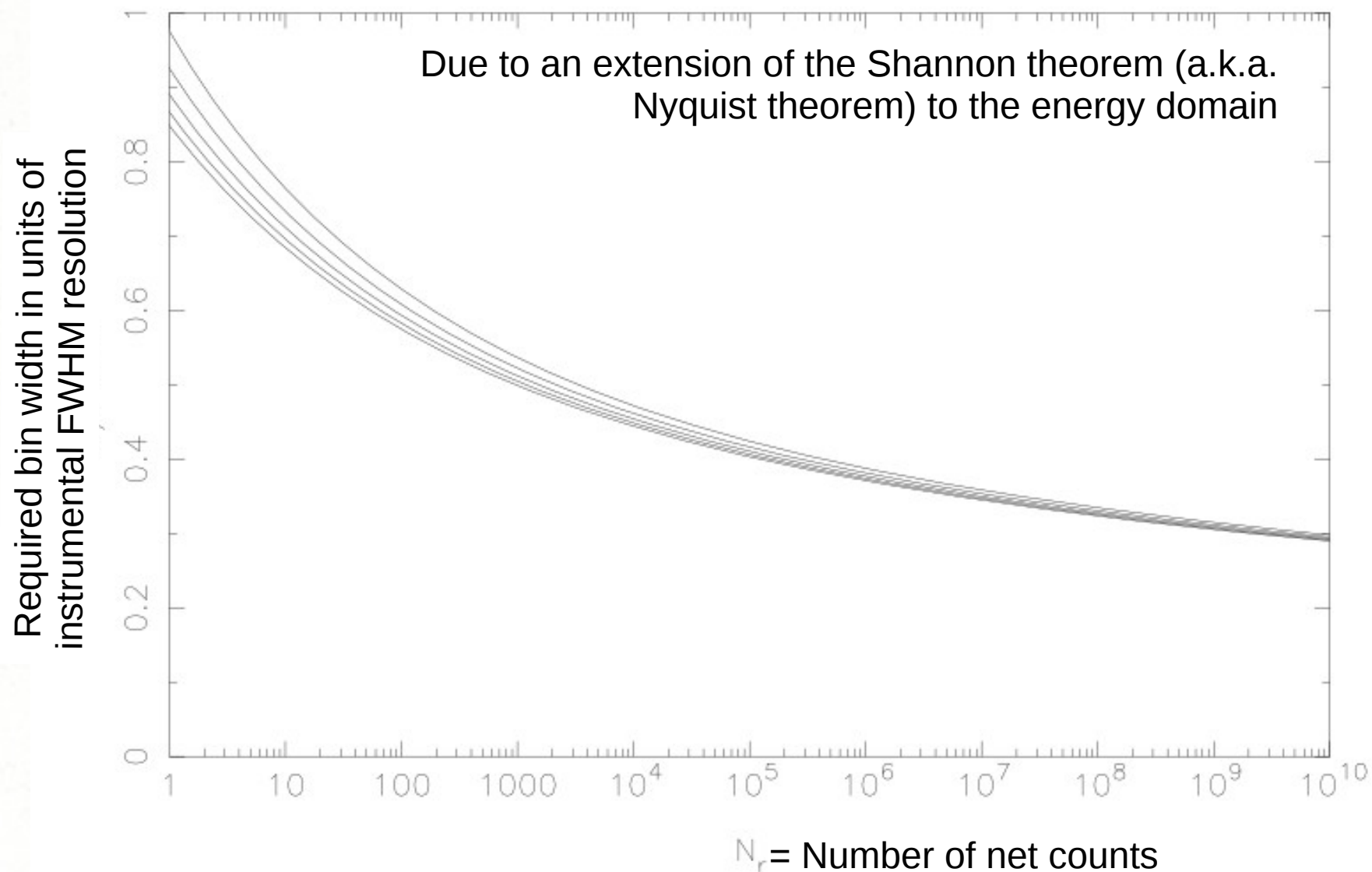
If  $g(\omega) = 0$  for all  $|\omega| > W$  for a given frequency  $W$ , then  $f(t)$  is band-limited, and in that case Shannon has shown that

$$f(t) = f_s(t) \equiv \sum_{n=-\infty}^{\infty} f(n\Delta) \frac{\sin \pi(t/\Delta - n)}{\pi(t/\Delta - n)}. \quad (1.7)$$

In (1.7), the bin size  $\Delta = 1/2W$ . Thus, a band-limited signal is completely determined by its values at an equally spaced grid with spacing  $\Delta$ .

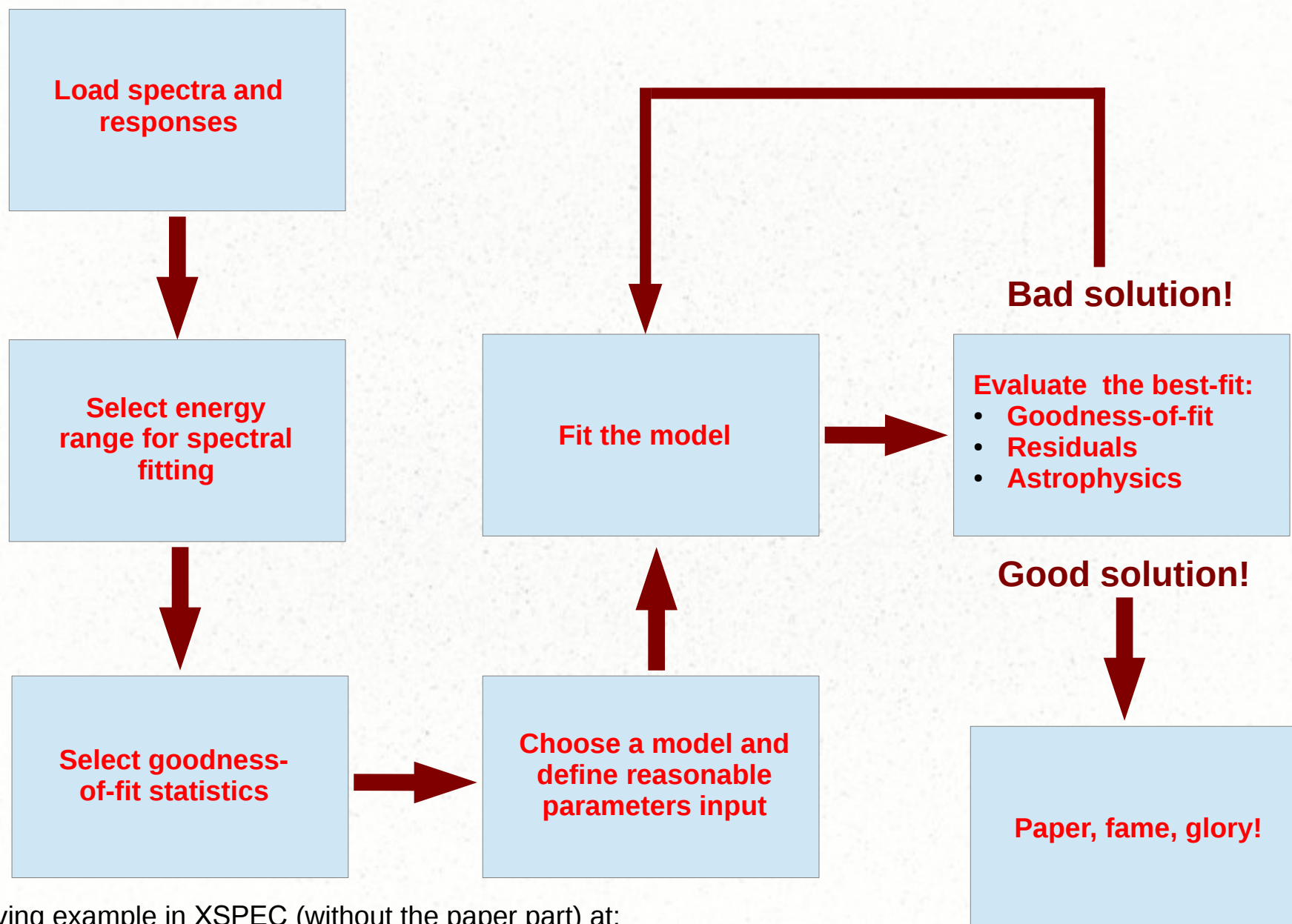


# An ideal rebinning strategy



**specgroup** in SAS implements this, and many other spectral rebinning schemes

# Forward-folding in action



Living example in XSPEC (without the paper part) at:

<http://heasarc.gsfc.nasa.gov/docs/xanadu/xspec/manual/XspecWalkthrough.html>