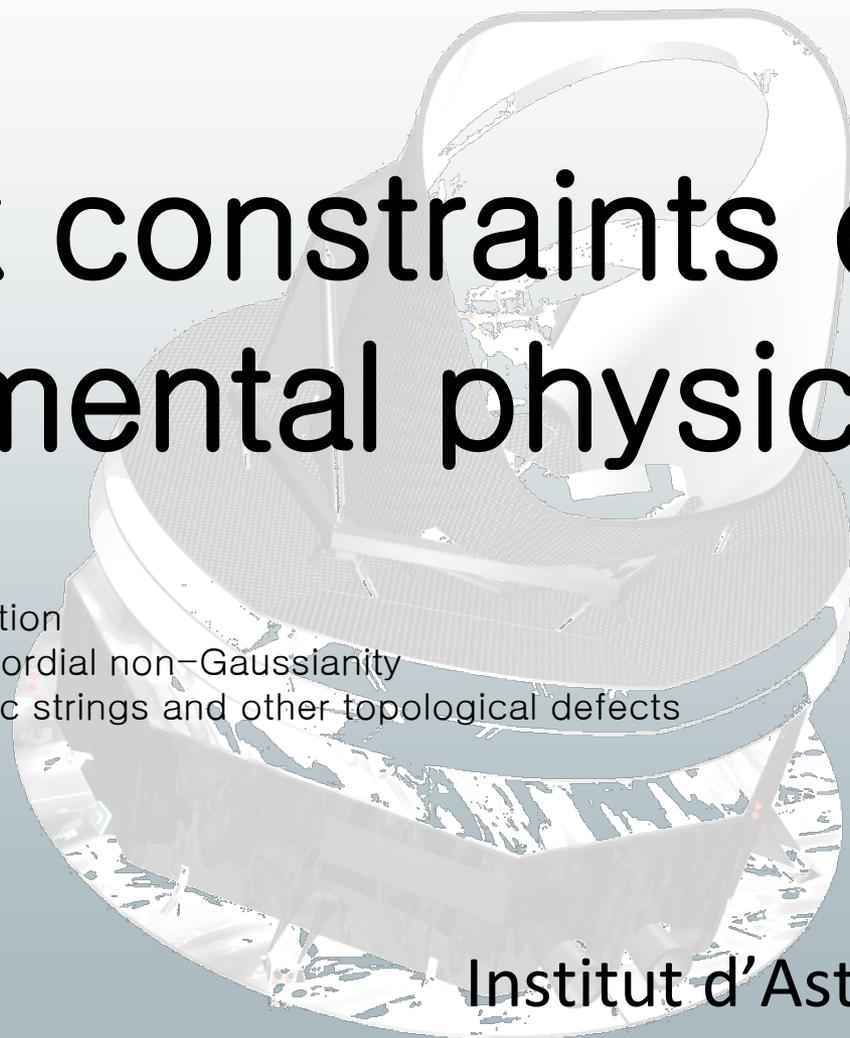


Planck constraints on fundamental physics



- XXII. Constraints on inflation
- XXIV Constraints on primordial non-Gaussianity
- XXV. Searches for cosmic strings and other topological defects

François R. Bouchet

Institut d'Astrophysique de Paris

On behalf of the Planck collaboration

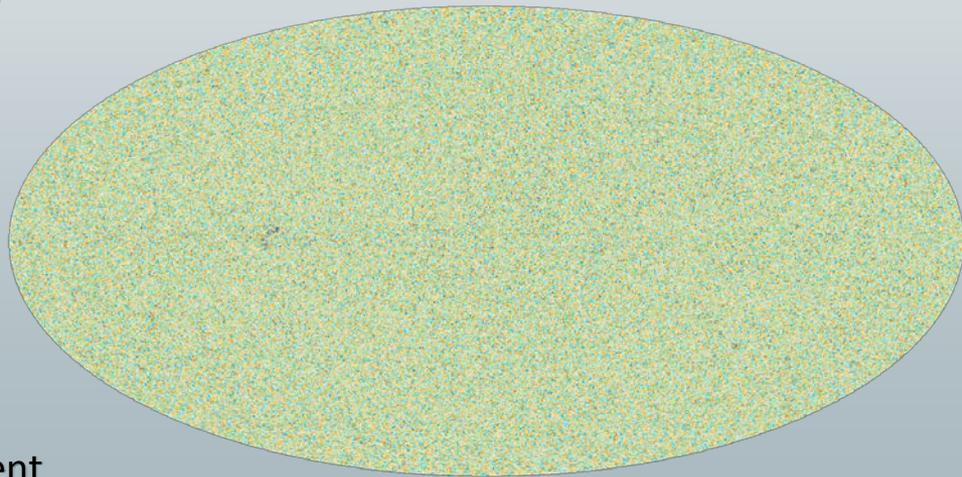
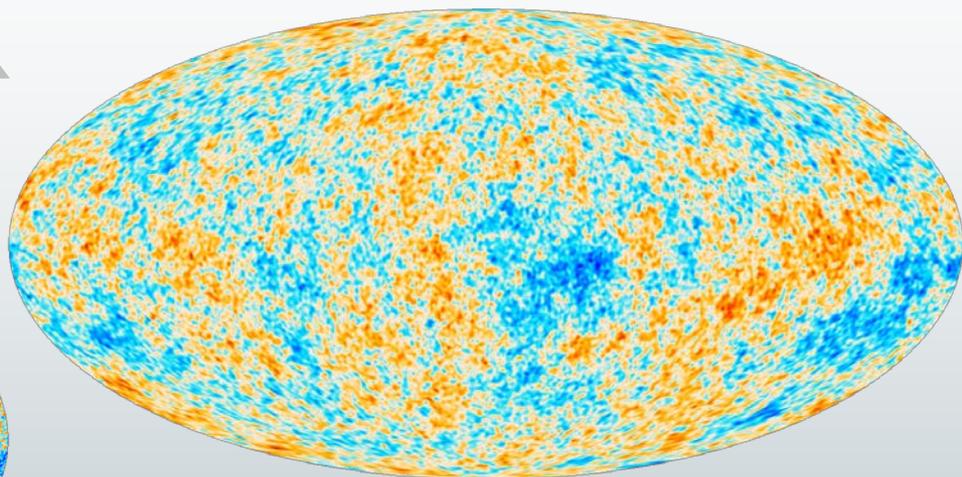


Our window



Smoothed map (suppressing scales $\theta < 1$ deg) :

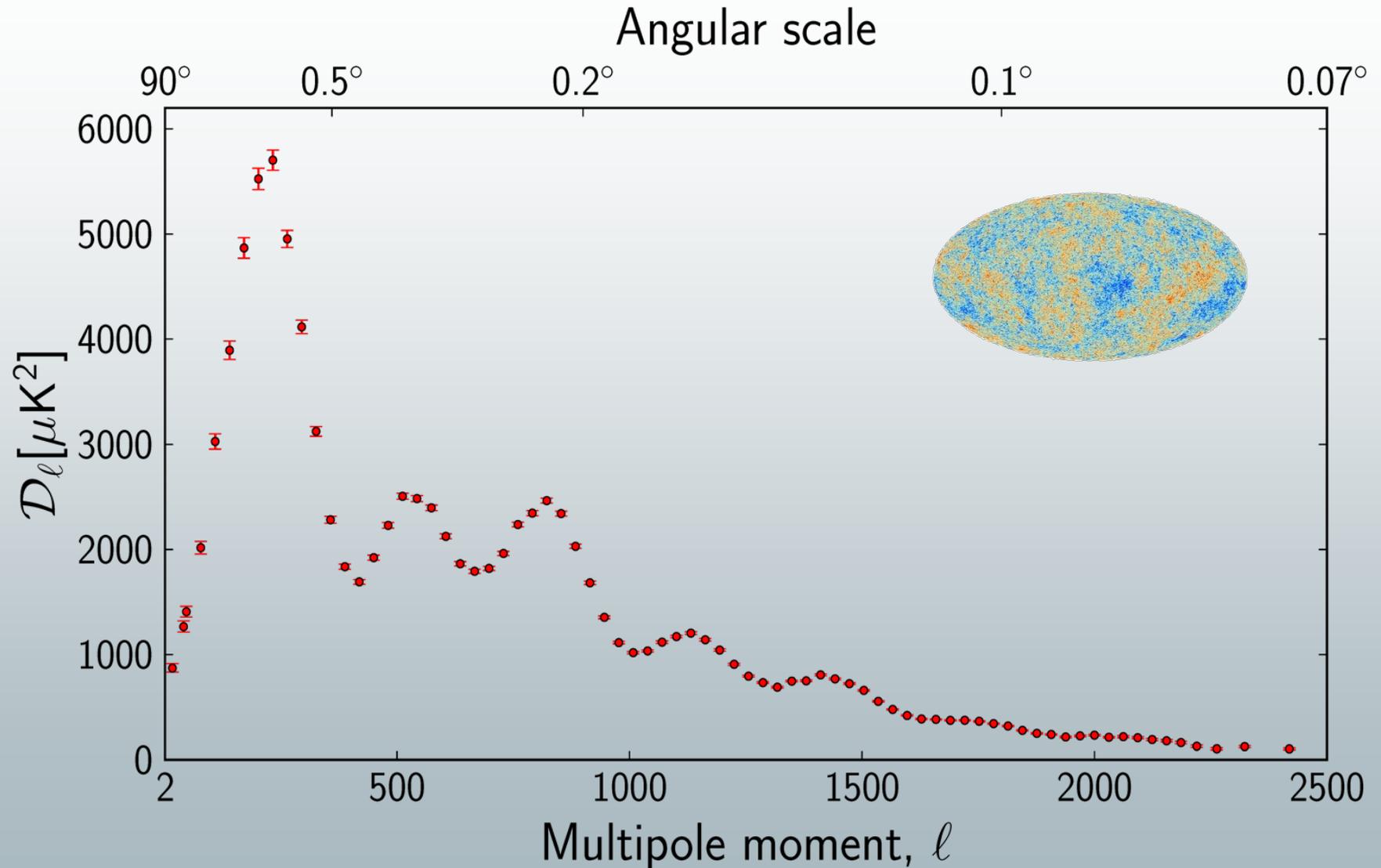
Quantum Fluctuations imprinted
When the age of the Universe was in the
interval [10-30, 10-12] seconds



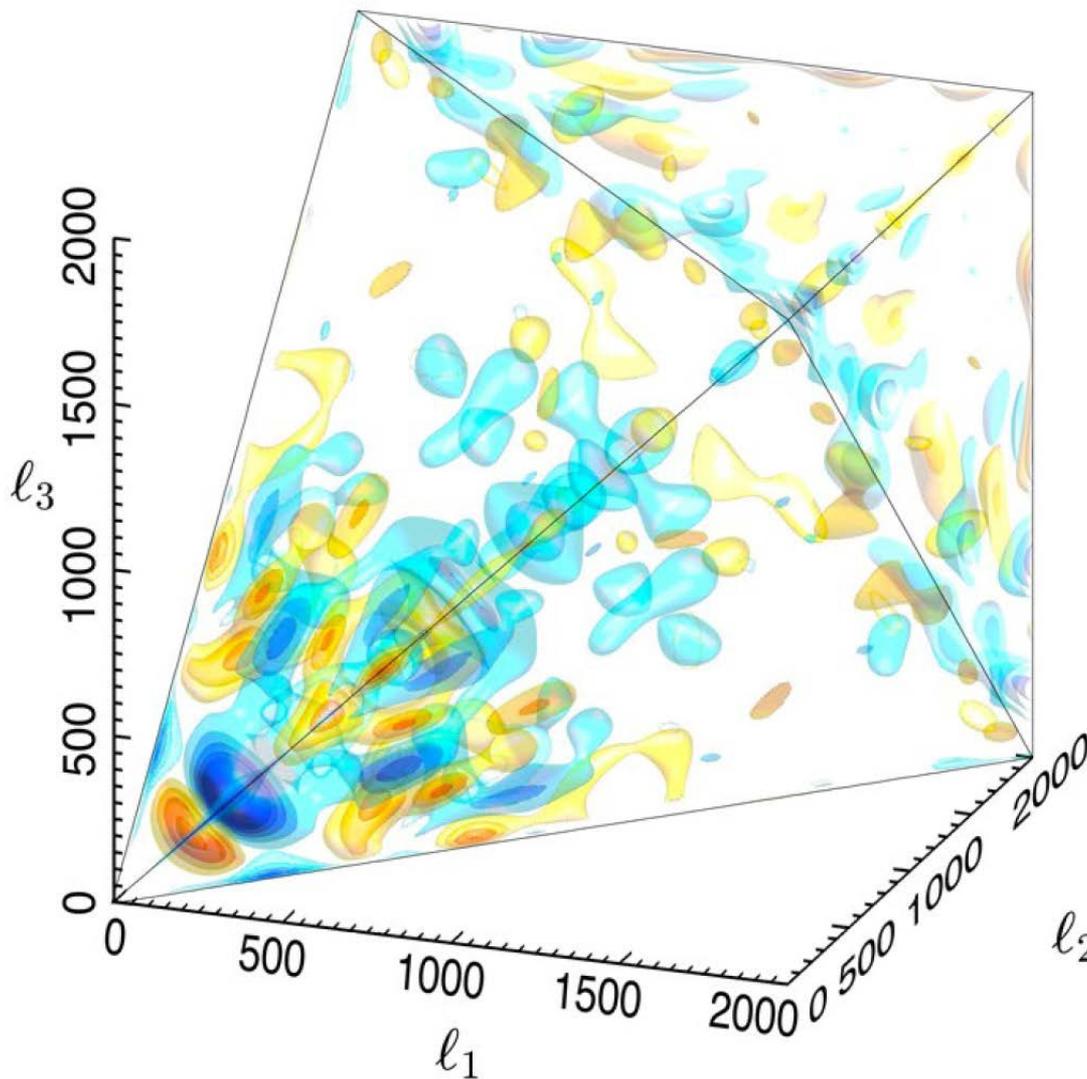
Difference map (scales $\theta < 1$ deg) :
Acoustic oscillations at small scales
< ct when $t=380\,000$ years (~ 150 Mpc today).
Which allows to take a census of the Universe content



The Planck spectrum of Temperature anisotropies



The Planck bi-spectrum of Temperature anisotropies





Base Λ CDM model 6 parameters

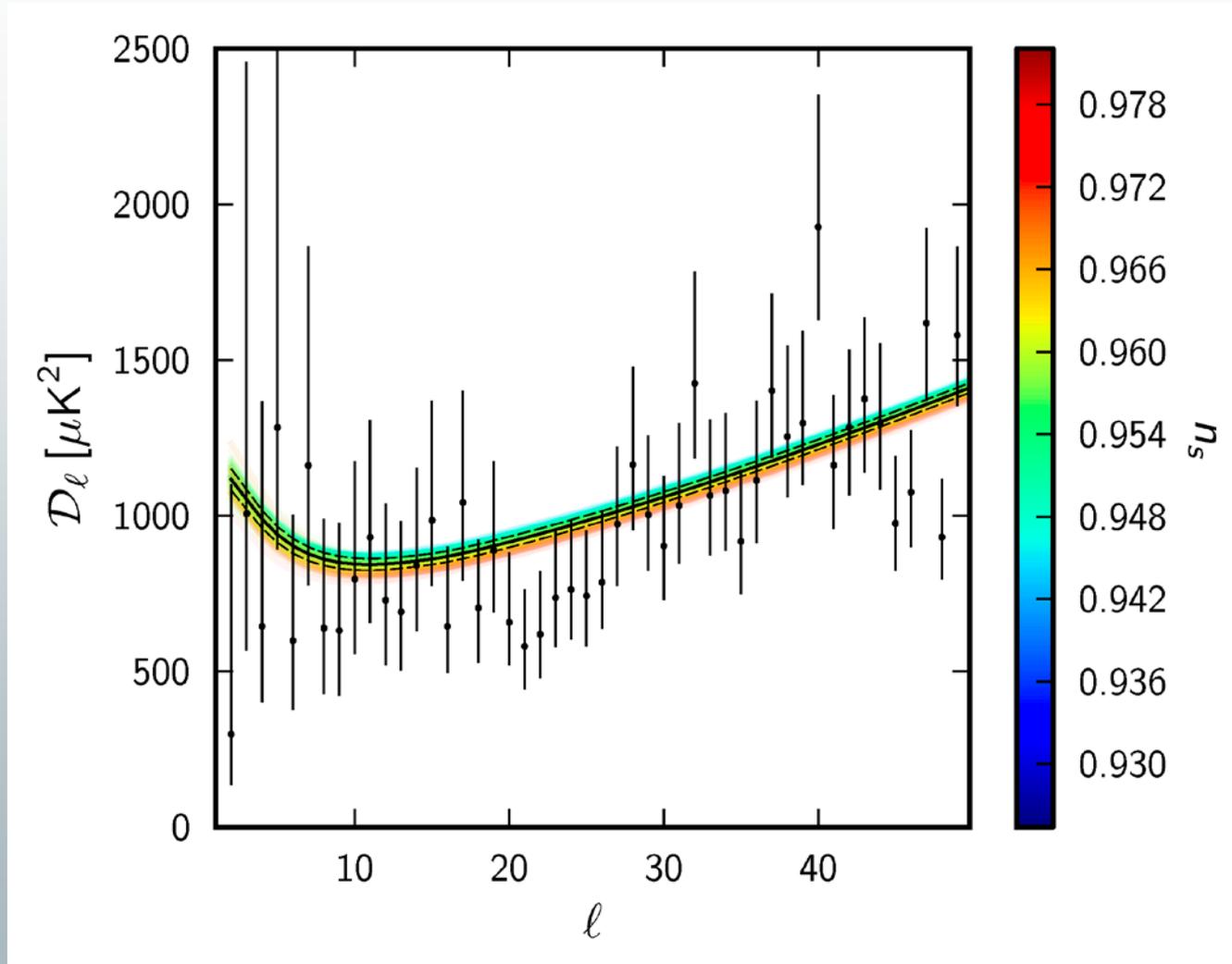


Parameter	<i>Planck</i> (CMB+lensing)		<i>Planck</i> +WP+highL+BAO	
	Best fit	68 % limits	Best fit	68 % limits
$\Omega_b h^2$	0.022242	0.02217 ± 0.00033	0.022161	0.02214 ± 0.00024
$\Omega_c h^2$	0.11805	0.1186 ± 0.0031	0.11889	0.1187 ± 0.0017
$100\theta_{MC}$	1.04150	1.04141 ± 0.00067	1.04148	1.04147 ± 0.00056
τ	0.0949	0.089 ± 0.032	0.0952	0.092 ± 0.013
n_s	0.9675	0.9635 ± 0.0094	0.9611	0.9608 ± 0.0054
$\ln(10^{10} A_s)$	3.098	3.085 ± 0.057	3.0973	3.091 ± 0.025

The sound horizon, θ , determined by the positions of the peaks, is now determined with 0.05% precision (links together $\Omega_b h^2$, $\Omega_c h^2$, H_0)

$$\theta_s = (1.04148 \pm 0.00066) \times 10^{-2} = 0.596724^\circ \pm 0.00038^\circ$$

Exact scale invariance of the primordial fluctuations is ruled out, at more than 7σ (as predicted by base inflation models)





Inflation has a few variants...



- assisted brane inflation
- anomaly-induced inflation
- assisted inflation
- assisted chaotic inflation
- B-inflation
- boundary inflation
- brane inflation
- brane-assisted inflation
- brane gas inflation
- brane-antibrane inflation
- braneworld inflation
- Brans-Dicke chaotic inflation
- Brans-Dicke inflation
- bulky brane inflation
- chaotic inflation
- chaotic hybrid inflation
- chaotic new inflation
- Chromo-Natural Inflation
- D-brane inflation
- D-term inflation
- dilaton-driven inflation
- dilaton-driven brane inflation
- double inflation
- double D-term inflation
- dual inflation
- dynamical inflation
- dynamical SUSY inflation
- S-dimensional assisted inflation
- eternal inflation
- extended inflation
- extended open inflation
- extended warm inflation
- extra dimensional inflation
- ...



- F-term inflation
- F-term hybrid inflation
- false-vacuum inflation
- false-vacuum chaotic inflation
- fast-roll inflation
- first-order inflation
- gauged inflation
- Ghost inflation
- Hagedorn inflation
- higher-curvature inflation
- hybrid inflation
- Hyper-extended inflation
- induced gravity inflation
- intermediate inflation
- inverted hybrid inflation
- Power-law inflation
- K-inflation
- Super symmetric inflation

- Quintessential inflation
- Roulette inflation
- curvature inflation
- Natural inflation
- Warm natural inflation
- Super inflation
- Super natural inflation
- Thermal inflation
- Discrete inflation
- Polarcap inflation
- Open inflation
- Topological inflation
- Multiple inflation
- Warm inflation
- Stochastic inflation
- Generalised assisted inflation
- Self-sustained inflation
- Graduated inflation
- Local inflation
- Singular inflation
- Slinky inflation
- Locked inflation
- Elastic inflation
- Mixed inflation
- Phantom inflation
- Non-commutative inflation
- Tachyonic inflation
- Tsunami inflation
- Lambda inflation
- Steep inflation
- Oscillating inflation
- Mutated hybrid inflation
- Inhomogeneous inflation
- ...



But there is a simple reference model



- $\Omega_k = 0$, single inflaton field, obeying slow-roll conditions (small derivatives of the potential, ϵ, η, ξ), deriving from a standard Lagrangian, with fluctuations behaving as in flat space at asymptotically early times and short distances (Bunch-Davies Initial conditions)
- ➔ Adiabatic fluctuations, $P_R(k) = A_s (k/k_*)^{n_s - 1}$, $n_s < 1$, no running, $N_* \in [50, 60]$, with negligible non-Gaussianity.
- We shall probe Ω_k , n_s , running ($dn_s/d\ln k$), non-slow-roll, features in $P(k)$ (parametric and non-parametric), isocurvature fluctuations, existence of defects, LEO non-gaussianity and a few more...

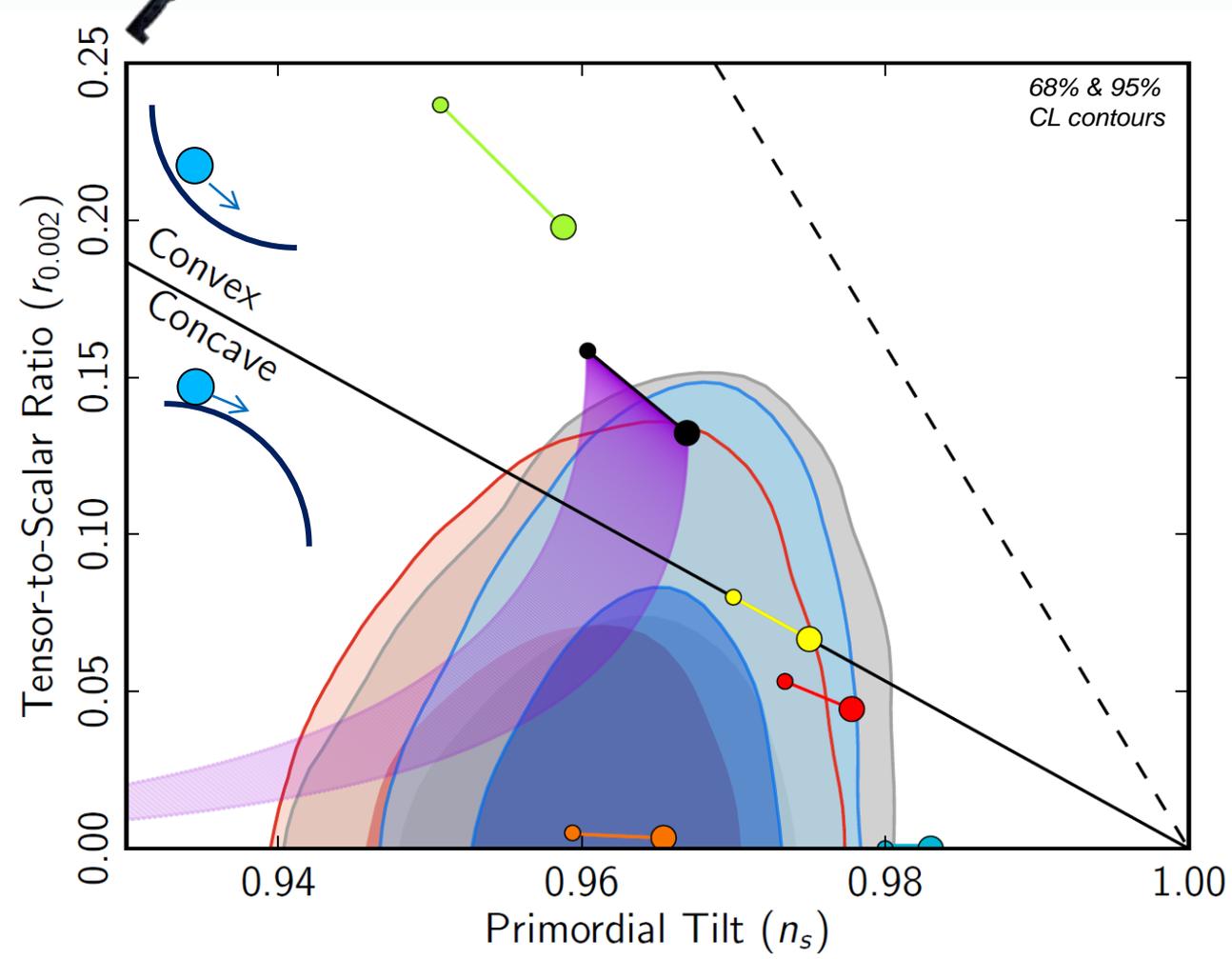


Few “specimen” in the slow roll approximation

- Power law potential and chaotic inflation $V(\phi) = \lambda M_{\text{pl}}^4 \left(\frac{\phi}{M_{\text{pl}}}\right)^n$
 - Simplest is $n=2$ (Linde 1983)
 - Axion monodromy $\rightarrow n=1, n=2/3$ (2010, 2008)
- Exponential potential & power law inf. $V(\phi) = \Lambda^4 \exp\left(-\lambda \frac{\phi}{M_{\text{pl}}}\right)$ (1985) $a(t) \propto t^{2/\lambda^2}$
- Inverse power law $V(\phi) = \Lambda^4 \left(\frac{\phi}{M_{\text{pl}}}\right)^{-\beta}$ (1990)
- Hill-top model $V(\phi) \approx \Lambda^4 \left(1 - \frac{\phi^r}{2\mu^p} + \dots\right)$ (Linde 1982, Albrecht & Steinhardt)
- SB potential $V(\phi) = \Lambda^4 \left(1 - \frac{\phi^2}{\mu^2}\right)^2$ (1990)
- Natural inflation $V(\phi) = \Lambda^4 \left[1 + \cos\left(\frac{\phi}{f}\right)\right]$ (1990) (1994)
- Hybrid inflation $V(\phi, \chi) = \Lambda^4 \left(1 - \frac{\chi^2}{u^2}\right)^2 + U(\phi) + \frac{g^2}{2} \phi^2 \chi^2$ $U(\phi) = \alpha_h \Lambda^4 \ln\left(\frac{\phi}{\mu}\right)$ SBSUSY
- R2 inflation $S = \int d^4x \sqrt{-g} \frac{M_{\text{pl}}^2}{2} \left(R + \frac{R^2}{6M^2}\right)$ (Starobinsky 1980, Mukhanov & Chibisov 1981)
- ... (Naming arbitrarily stopped at 1983)



Constraint on representative Inflation models



$V_* = (1.9 \times 10^{16} \text{ GeV})^4 \frac{r_*}{0.12}$

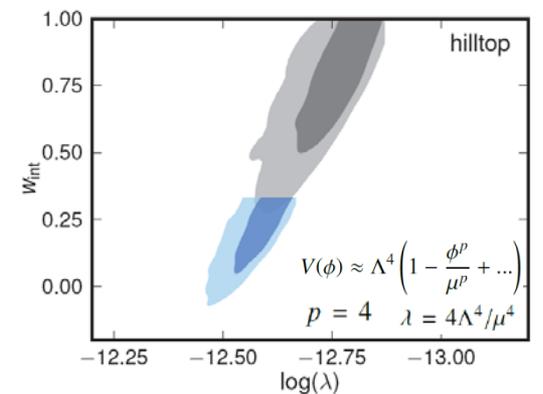
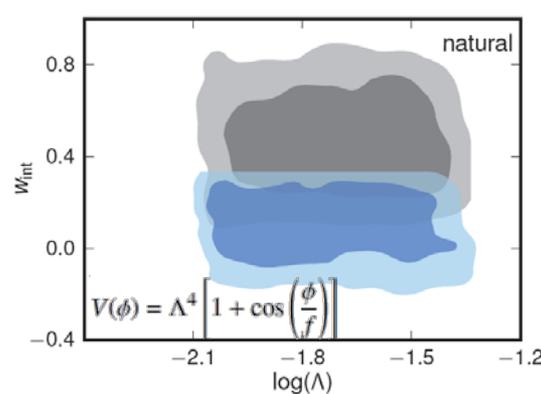
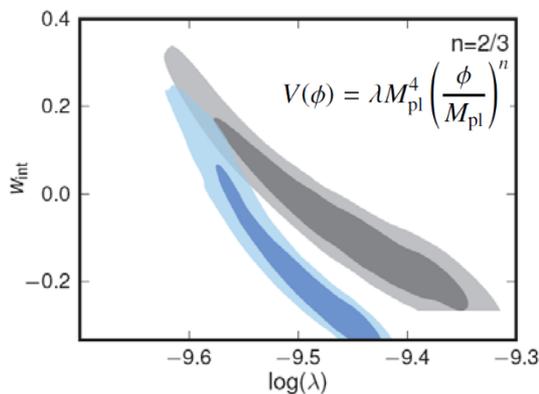
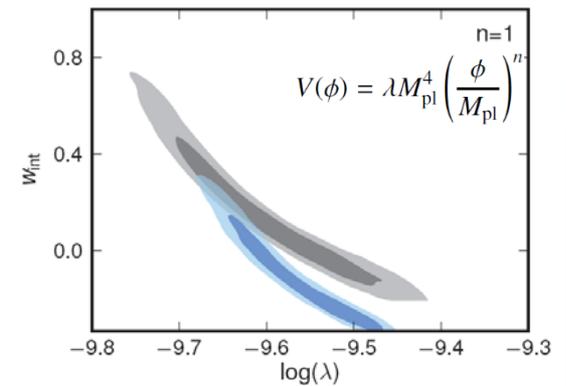
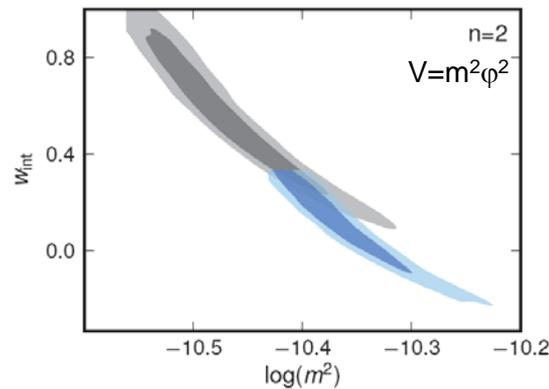
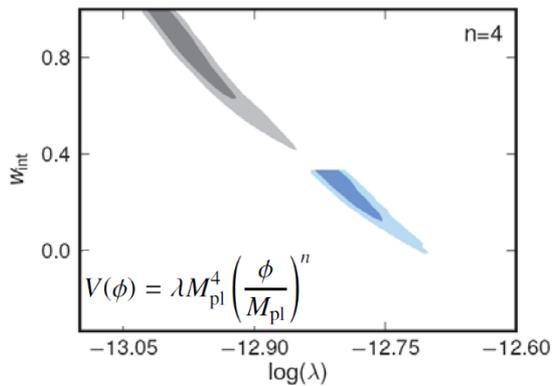
- Planck+WP
- Planck+WP+highL
- Planck+WP+BAO
- Natural Inflation
- Power law inflation
- Low Scale SSB SUSY
- R^2 Inflation (Higgs $\xi \gg 1$)
- $V \propto \phi^{2/3}$
- $V \propto \phi$
- $V \propto \phi^2$ Chaotic
- $V \propto \phi^3$
- $N_* = 50$
- $N_* = 60$

➔ Exponential potential models (power-law inf.), simplest hybrid inflationary models (SB SUSY), monomial potential models of degree $n > 2$ do not provide a good fit to the data.

$$N_* \approx 71.21 - \ln\left(\frac{k_*}{a_0 H_0}\right) + \frac{1}{4} \ln\left(\frac{V_{\text{hor}}}{M_{\text{pl}}^4}\right) + \frac{1}{4} \ln\left(\frac{V_{\text{hor}}}{\rho_{\text{end}}}\right) + \frac{1 - 3w_{\text{int}}}{12(1 + w_{\text{int}})} \ln\left(\frac{\rho_{\text{th}}}{\rho_{\text{end}}}\right)$$

$$2t \sim 3, 3t \sim 10, 4t \sim \mathcal{O}(1) \\ \rightarrow N_* \in [50, 60]$$

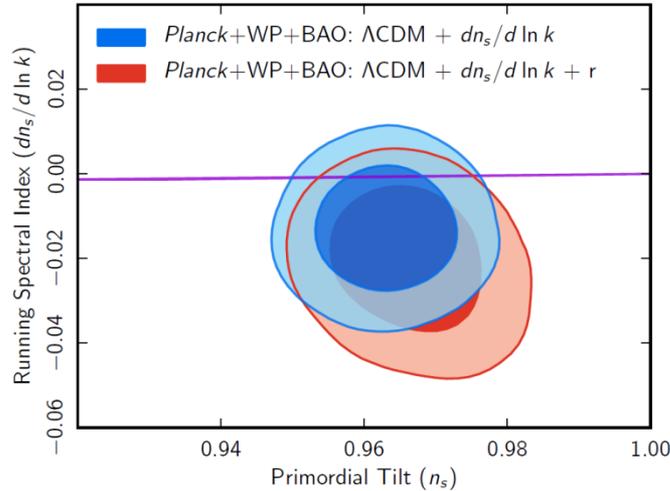
w_{int} = parameter of *effective* equation of state between end of inflation and thermalisation (instantaneous $\leftrightarrow w_{\text{int}}=1/3$)



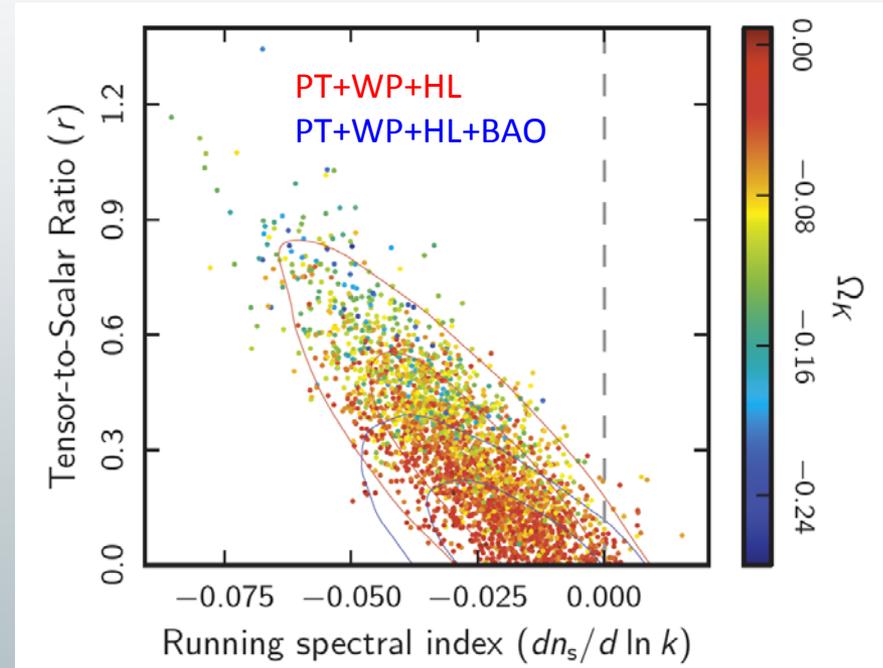
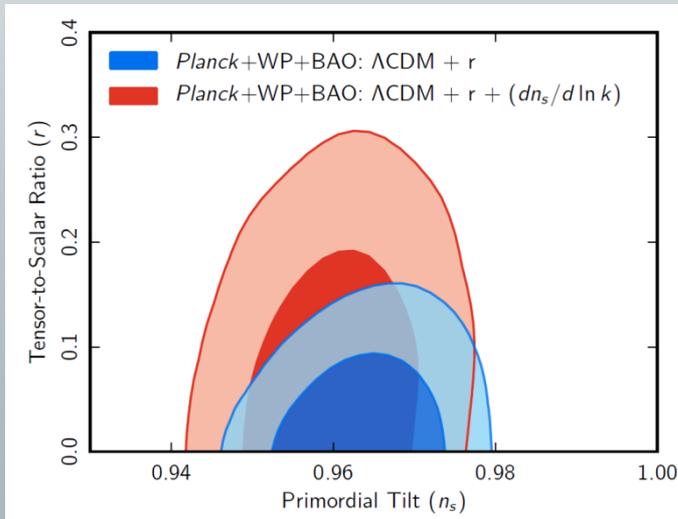
Blue and grey area respectively correspond to the restrictive ($\rho_{th}^{1/4} = 10^9 GeV, w_{int} \in [-\frac{1}{3}, \frac{1}{3}]$) and permissive ($\rho_{th}^{1/4} = 10^3 GeV, w_{int} \in [-\frac{1}{3}, 1]$) entropy generation scenario; red is instantaneous reheating for natural and hilltop.



$$\mathcal{P}_{\mathcal{R}}(k) = A_s \left(\frac{k}{k_*} \right)^{n_s - 1 + \frac{1}{2} \frac{dn_s}{d \ln k} \ln(k/k_*)}$$

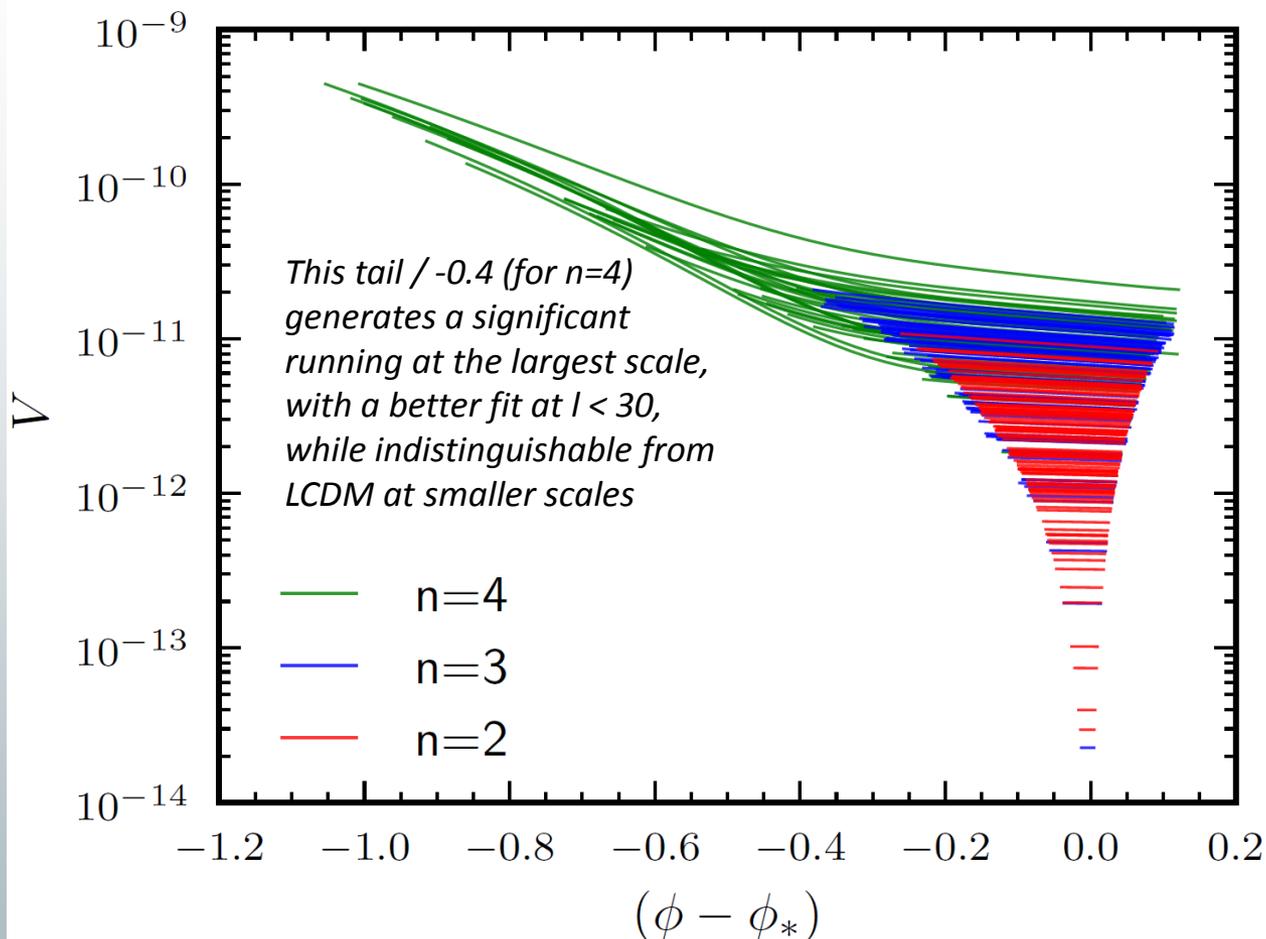


purple strip is the prediction for single monomial chaotic inflationary models with $50 < N_* < 60$.



→ Open inflation is not needed (to explain $|\Omega_k| \gtrsim 10^{-5}$)

$\Lambda\text{CDM} + dn_s/d \ln k = -0.013 \pm 0.009$ PT+WP
 $\Lambda\text{CDM} + r \rightarrow 0.021 \pm 0.012$ PT+WP



Best fitting potentials, when $V(\phi)$ is Taylor expanded at the n-th order around the pivot scale;

Planck-T+WP;
Flat priors on ϵ , η , ξ^2 ;

Φ_* in natural units / $(8\pi)^{1/2} M_{\text{pl}} = 1$.

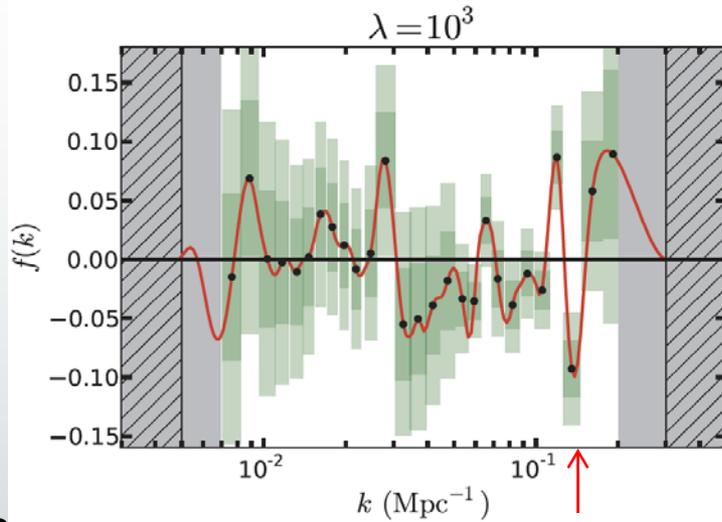
n	from $V(\phi)$		
	2	3	4
$\ln[10^{10} A_s]$	$3.087^{+0.050}_{-0.050}$	$3.115^{+0.066}_{-0.063}$	$3.130^{+0.071}_{-0.066}$
n_s	$0.961^{+0.015}_{-0.015}$	$0.958^{+0.017}_{-0.016}$	$0.954^{+0.018}_{-0.018}$
$100 \, dn_s/d \ln k$	$-0.05^{+0.13}_{-0.14}$	$-2.2^{+2.2}_{-2.3}$	$-0.61^{+3.1}_{-3.1}$
$100 \, d^2 n_s/d \ln k^2$	$-0.01^{+0.73}_{-0.75}$	$-0.3^{+1.0}_{-1.2}$	$6.3^{+8.6}_{-7.8}$
r	< 0.12	< 0.22	< 0.35



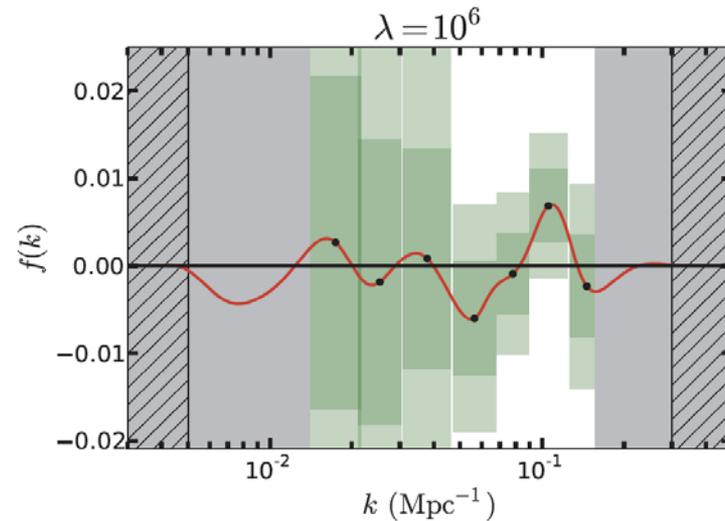
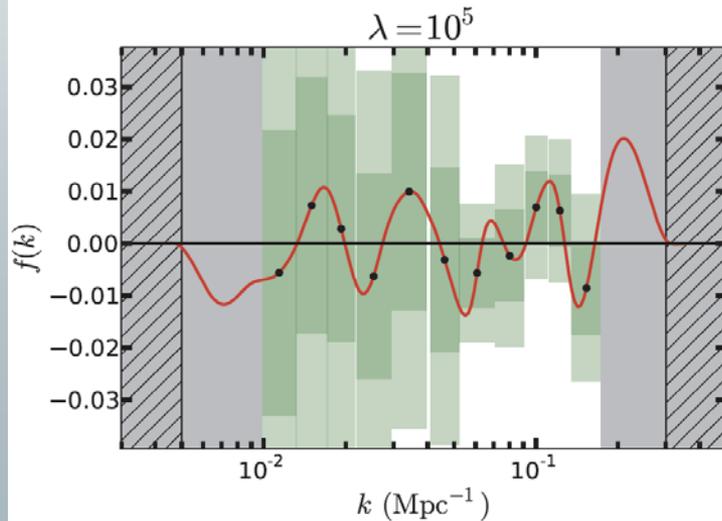
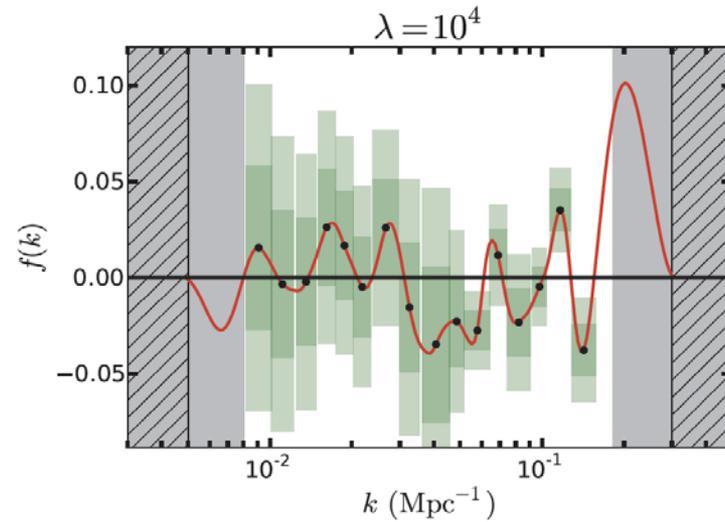
Features search in $P(k)$ - non-parametric



1&2 σ
boxes



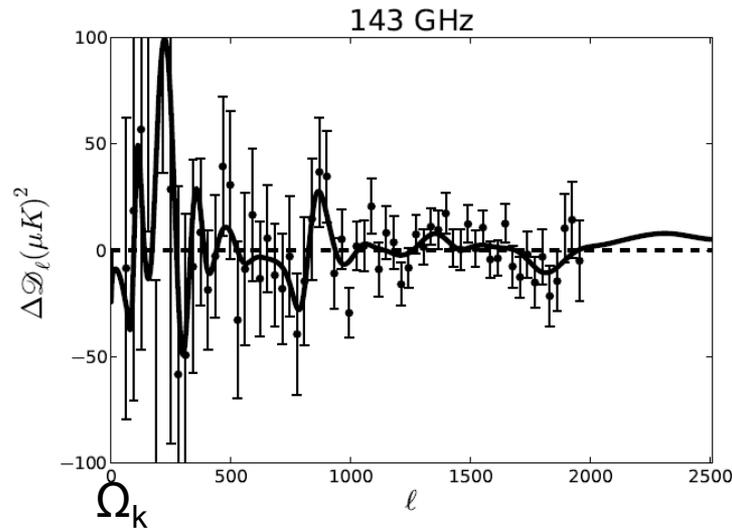
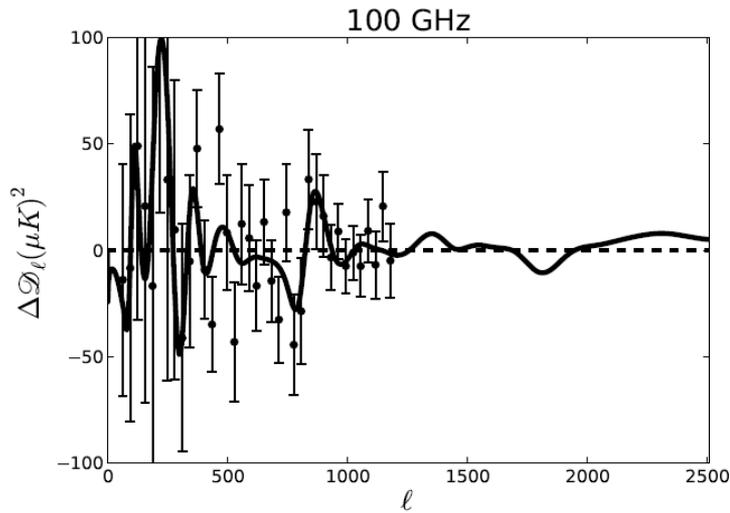
$\sim 3.1\sigma$ @ $k=0.13$



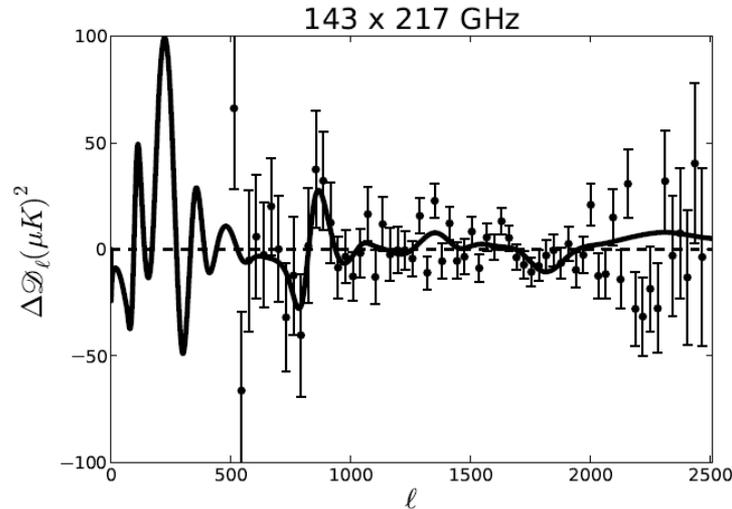
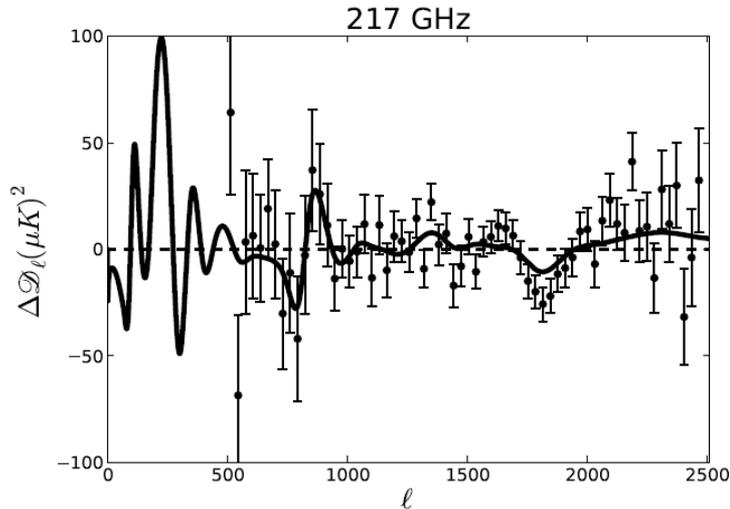
Red=ML
solution for
given
smoothing
parameter λ

Nuisance
parameters
(beams)
fixed to their
fiducial best
fit values

What is fitted in the CI's?



$$\lambda = 10^3$$



The features in the T power spectrum, (broad dip at 1800), cannot be explained by any of the known systematics which we propagated through the data analysis pipeline.

Analysis of the full mission data will help assessing whether these small departures from the best fit LCDM spectrum may possibly be due to unknown systematic effects, (or inaccurate propagation of known syst. effects), into the final power spectra.



Features search (parametric)

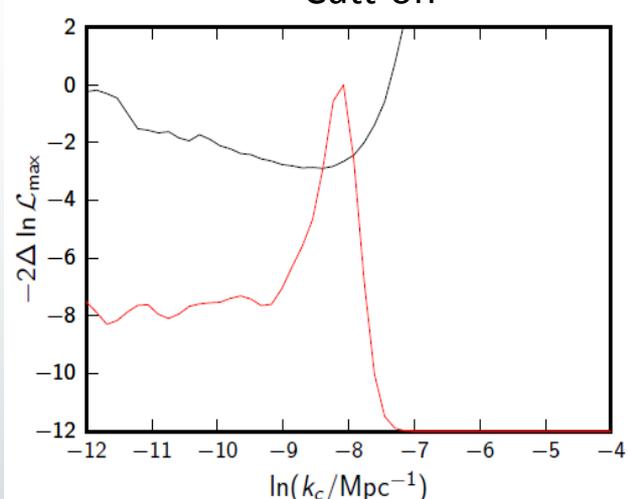
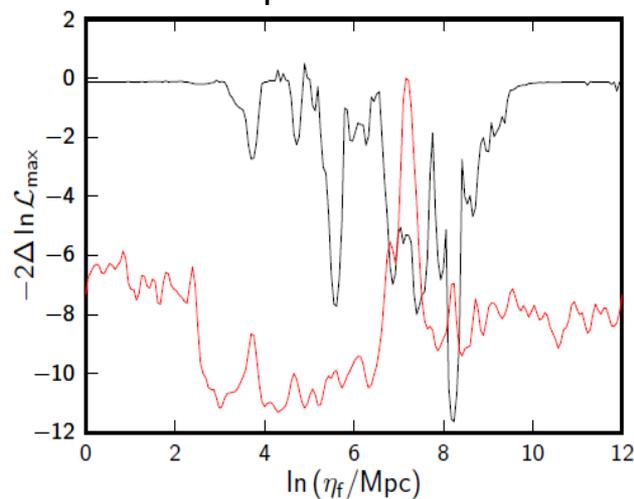
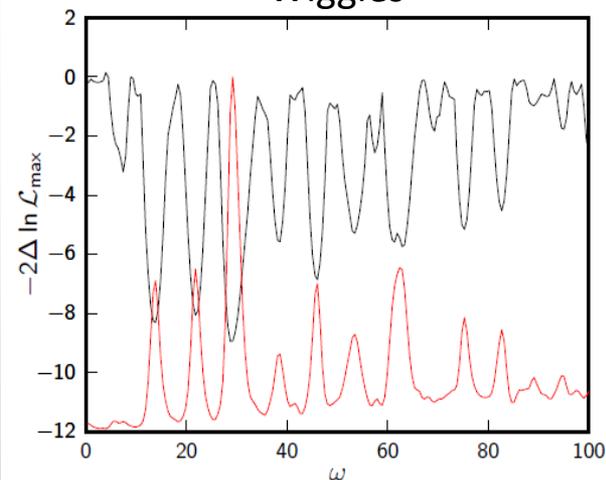


planck

Wiggles

Step inflation

Cutt-off



Profile likelihood & **Posterior probability (marginalised)**

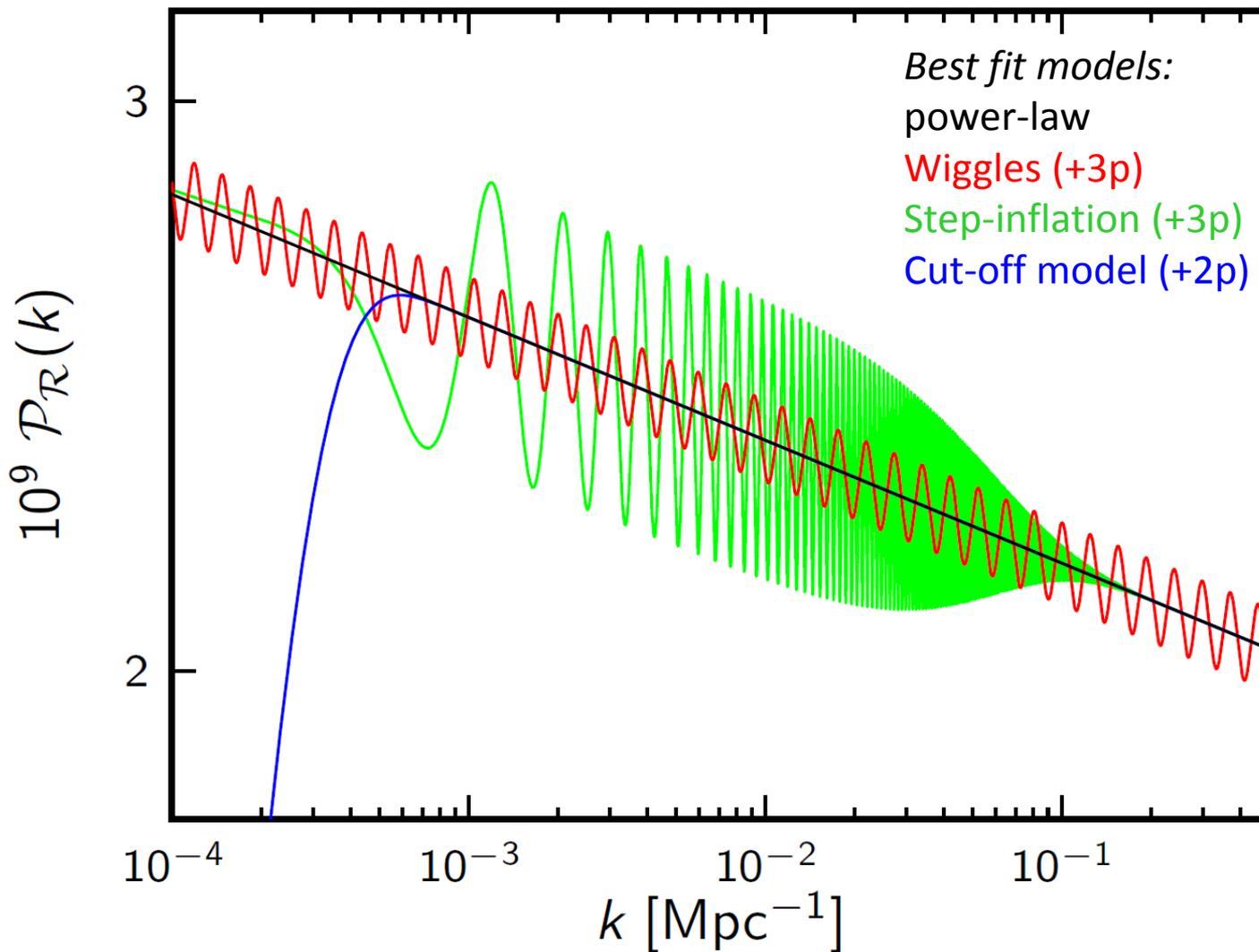
$$\mathcal{P}_0(k) = A_s \left(\frac{k}{k_*} \right)^{n_s-1} \quad \mathcal{P}_{\mathcal{R}}(k) = \mathcal{P}_0(k) \left\{ 1 + \alpha_w \sin \left[\omega \ln \left(\frac{k}{k_*} \right) + \varphi \right] \right\} \quad \text{wiggles}$$

$$\mathcal{P}_{\mathcal{R}}(k) = \exp \left[\ln \mathcal{P}_0(k) + \frac{\mathcal{A}_f}{3} \frac{k\eta_f/x_d}{\sinh(k\eta_f/x_d)} W'(k\eta_f) \right]$$

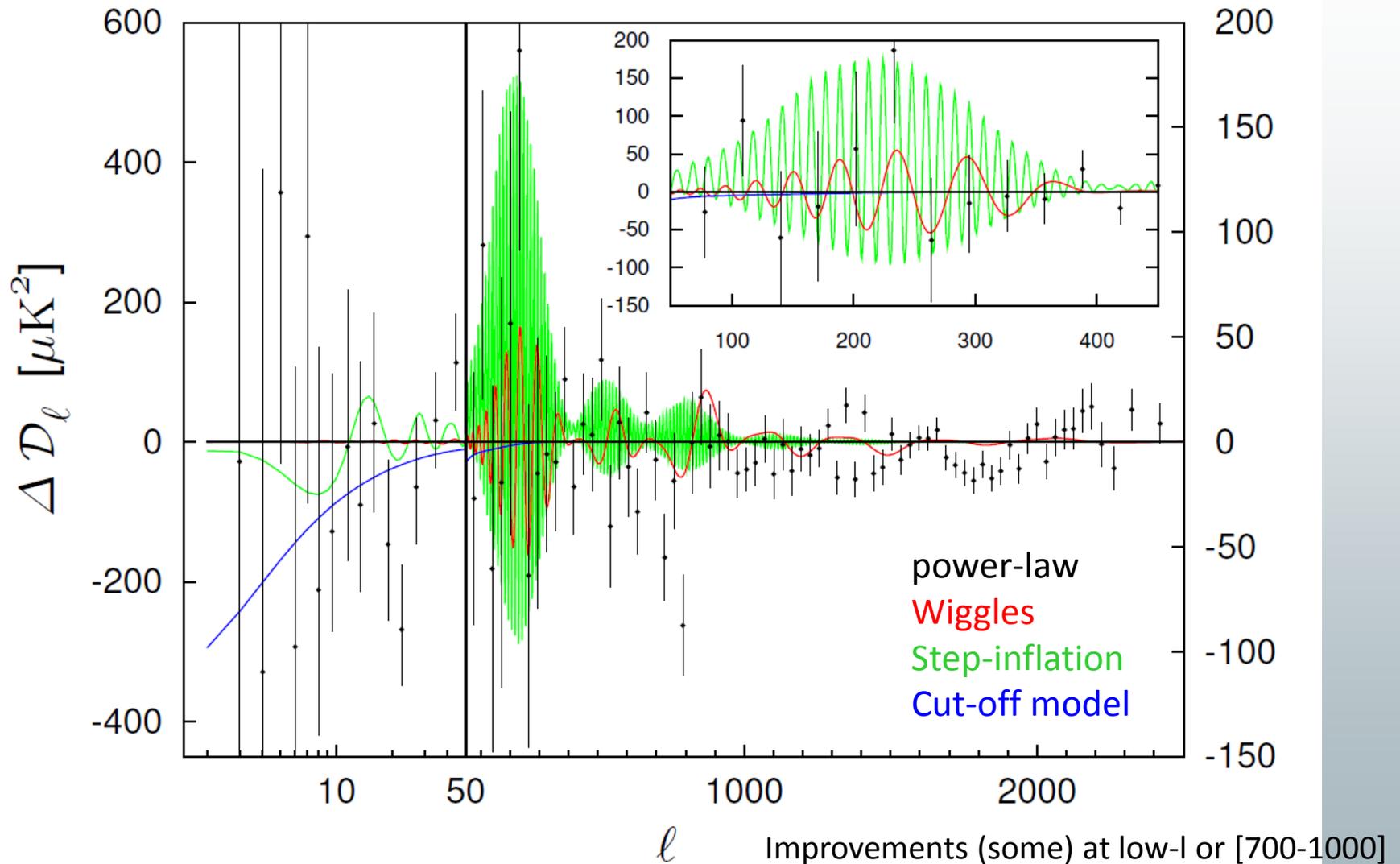
Step-inflation model

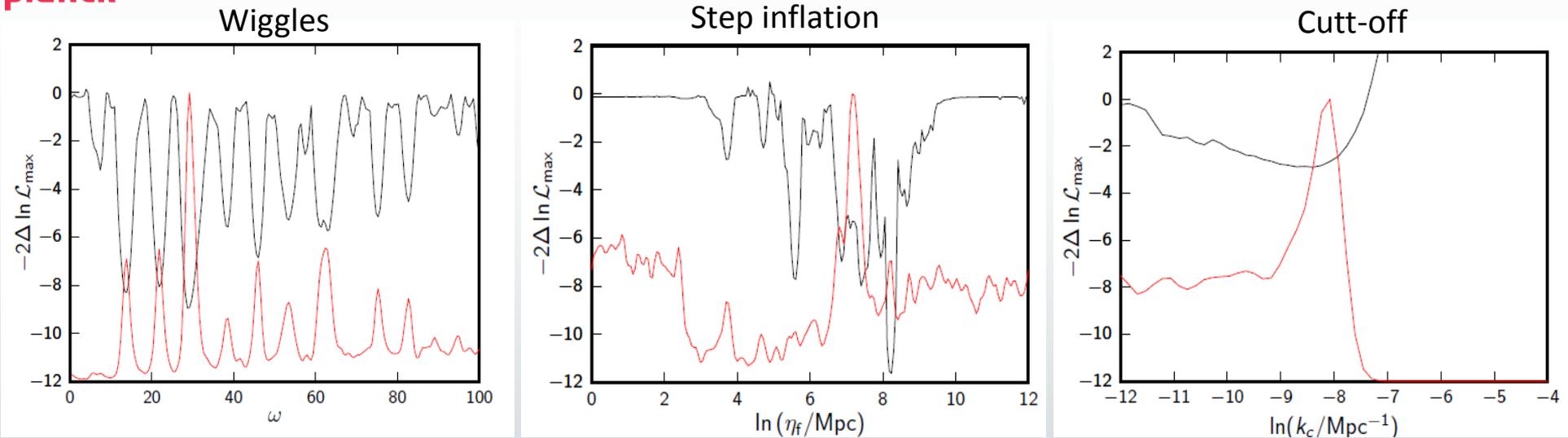
$$W'(x) = \left(-3 + \frac{9}{x^2} \right) \cos 2x + \left(15 - \frac{9}{x^2} \right) \frac{\sin 2x}{2x}$$

$$\mathcal{P}_{\mathcal{R}}(k) = \mathcal{P}_0(k) \left\{ 1 - \exp \left[- \left(\frac{k}{k_c} \right)^{\lambda_c} \right] \right\} \quad \text{Cut-off model}$$



Temperature residuals





Profile likelihood & Posterior probability (marginalised)

Are these improvements significant ?

- We have not run a full frequentist analysis to account for the “look elsewhere” effect
- Bayes: these models are not predictive enough (for our choice of prior) -> **ΛCDM preferred**

- If real, likely to show up in polarisation and NG searches

Model	$-2\Delta \ln \mathcal{L}_{\max}$	$\ln B_{0X}$	Parameter	Best fit value
Wiggles	-9.0	1.5	α_w	0.0294
			ω	28.90
			φ	0.075 π
Step-inflation	-11.7	0.3	\mathcal{A}_f	0.102
			$\ln(\eta_f/\text{Mpc})$	8.214
			$\ln x_d$	4.47
Cutoff	-2.9	0.3	$\ln(k_c/\text{Mpc}^{-1})$	-8.493
			λ_c	0.474

NB: Some improvement had already been noted earlier



Isocurvature modes

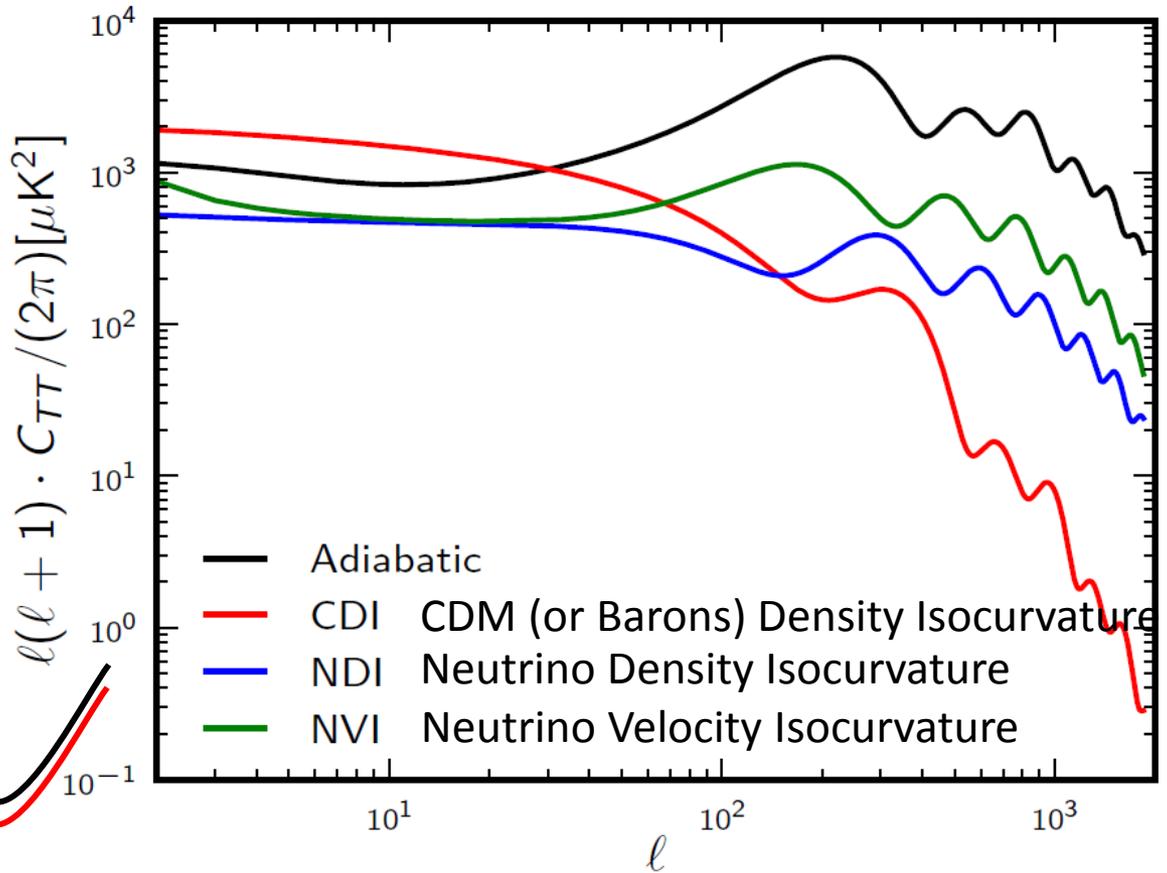


Isocurvature modes arise from spatial variations in the equation of state or from relative velocities between the components.

I.e. the density perturbations in the various components (e.g., baryons&leptons, CDM, photons, neutrinos) are not locked together.

These might be excited, e.g., with multi-component inflaton field.

Expect correlations between isocurvature and curvature degrees of freedom



$$\mathcal{P}(k) = \begin{pmatrix} \mathcal{P}_{\mathcal{R}} & \mathcal{P}_{\mathcal{R}\mathcal{I}_{\text{CDI}}} & \mathcal{P}_{\mathcal{R}\mathcal{I}_{\text{NDI}}} & \mathcal{P}_{\mathcal{R}\mathcal{I}_{\text{NVI}}} \\ \mathcal{P}_{\mathcal{I}_{\text{CDI}}\mathcal{R}} & \mathcal{P}_{\mathcal{I}_{\text{CDI}}\mathcal{I}_{\text{CDI}}} & \mathcal{P}_{\mathcal{I}_{\text{CDI}}\mathcal{I}_{\text{NDI}}} & \mathcal{P}_{\mathcal{I}_{\text{CDI}}\mathcal{I}_{\text{NVI}}} \\ \mathcal{P}_{\mathcal{I}_{\text{NDI}}\mathcal{R}} & \mathcal{P}_{\mathcal{I}_{\text{NDI}}\mathcal{I}_{\text{CDI}}} & \mathcal{P}_{\mathcal{I}_{\text{NDI}}\mathcal{I}_{\text{NDI}}} & \mathcal{P}_{\mathcal{I}_{\text{NDI}}\mathcal{I}_{\text{NVI}}} \\ \mathcal{P}_{\mathcal{I}_{\text{NVI}}\mathcal{R}} & \mathcal{P}_{\mathcal{I}_{\text{NVI}}\mathcal{I}_{\text{CDI}}} & \mathcal{P}_{\mathcal{I}_{\text{NVI}}\mathcal{I}_{\text{NDI}}} & \mathcal{P}_{\mathcal{I}_{\text{NVI}}\mathcal{I}_{\text{NVI}}} \end{pmatrix}$$

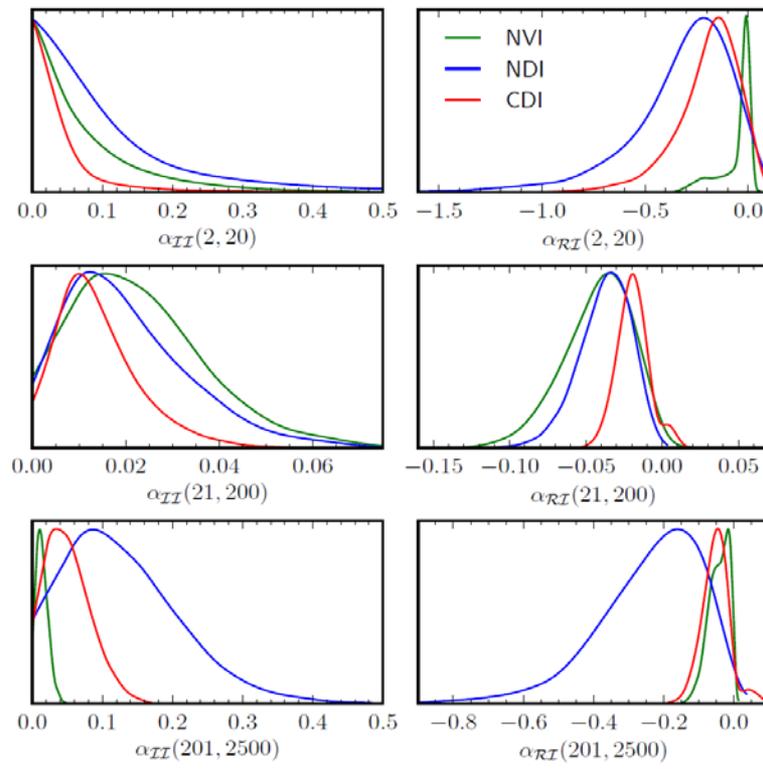
Only weak fractional contribution in various l -ranges allowed:

$$\alpha_{II}(\ell_{\min}, \ell_{\max}) = \frac{(\Delta T)_{II}^2(\ell_{\min}, \ell_{\max})}{(\Delta T)_{\text{tot}}^2(\ell_{\min}, \ell_{\max})}$$

$$\alpha_{RI}(\ell_{\min}, \ell_{\max}) = \frac{(\Delta T)_{RI}^2(\ell_{\min}, \ell_{\max})}{(\Delta T)_{\text{tot}}^2(\ell_{\min}, \ell_{\max})}$$

with

$$(\Delta T)_X^2(\ell_{\min}, \ell_{\max}) = \sum_{\ell=\ell_{\min}}^{\ell_{\max}} (2\ell + 1) C_{X,\ell}^{TT}$$



Low- l
[2,20]

Mid- l
[21,200]

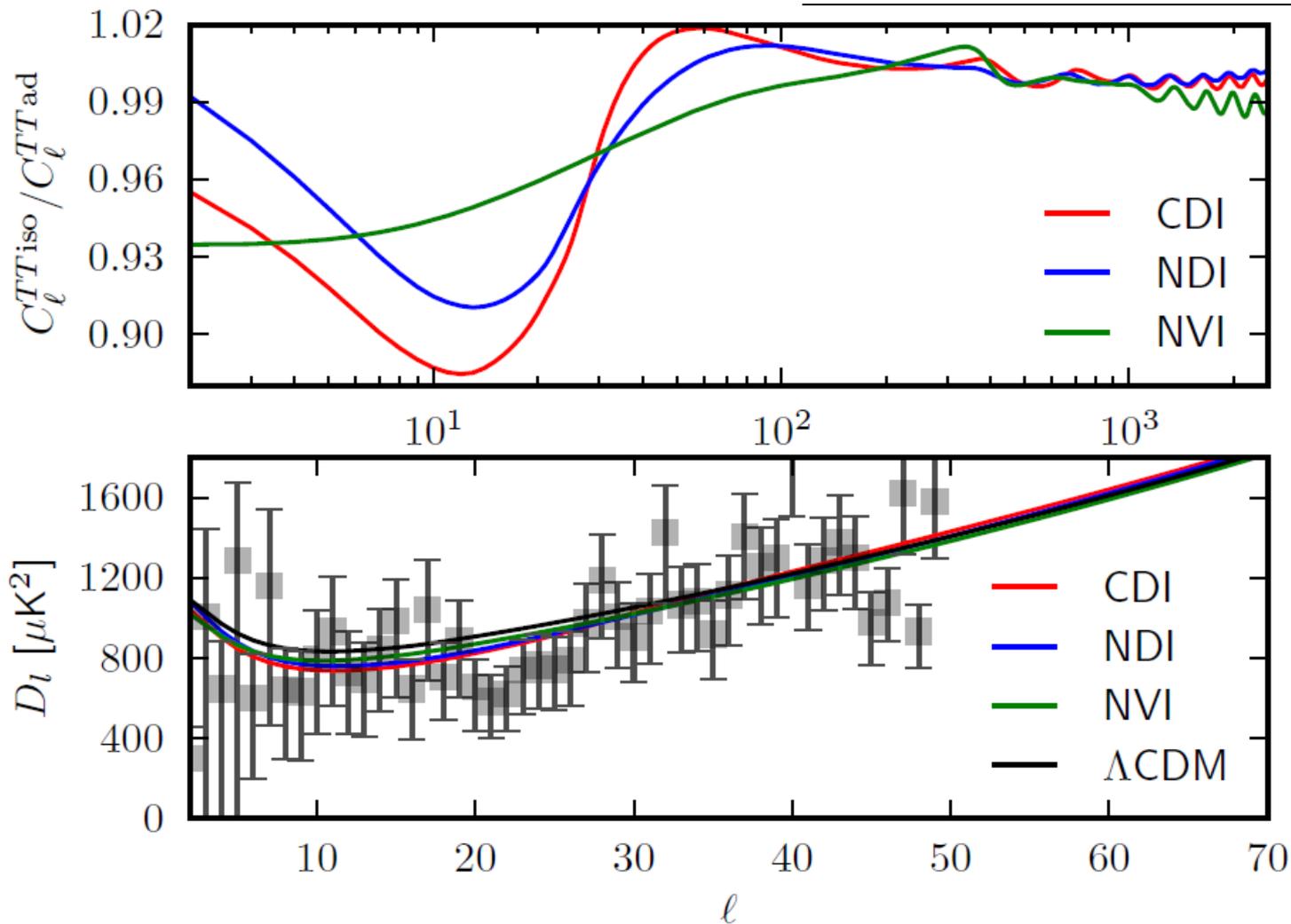
High- l
[201,2500]

$r_D > 0.98$

Model	95%CL bounds	$\alpha_{RR}^{(2,2500)}$	$\alpha_{II}^{(2,2500)}$	$\alpha_{RI}^{(2,2500)}$	Δn	$-2\Delta \ln \mathcal{L}_{\max}$
General model:						
CDM isocurvature		[0.98:1.07]	0.039	[-0.093:0.014]	4	-4.6
ND isocurvature		[0.99:1.09]	0.093	[-0.18:0]	4	-4.2
NV isocurvature		[0.96:1.05]	0.068	[-0.090:0.026]	4	-2.5
Special CDM isocurvature cases:						
Uncorrelated, $n_{II} = 1$, ("axion")		[0.98:1]	0.016	—	1	0
Fully correlated, $n_{II} = n_{RR}$, ("curvaton")		[0.97:1]	0.0011	[0:0.028]	1	0
Fully anti-correlated, $n_{II} = n_{RR}$		[1:1.06]	0.0046	[-0.067:0]	1	-1.3

Only weak evidence

Assuming a single isocurvature mode



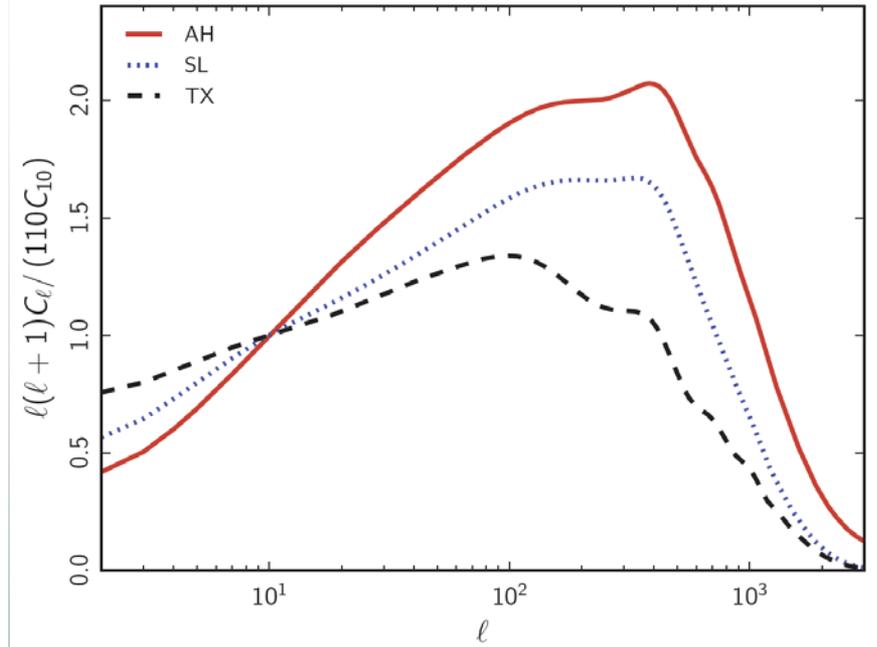
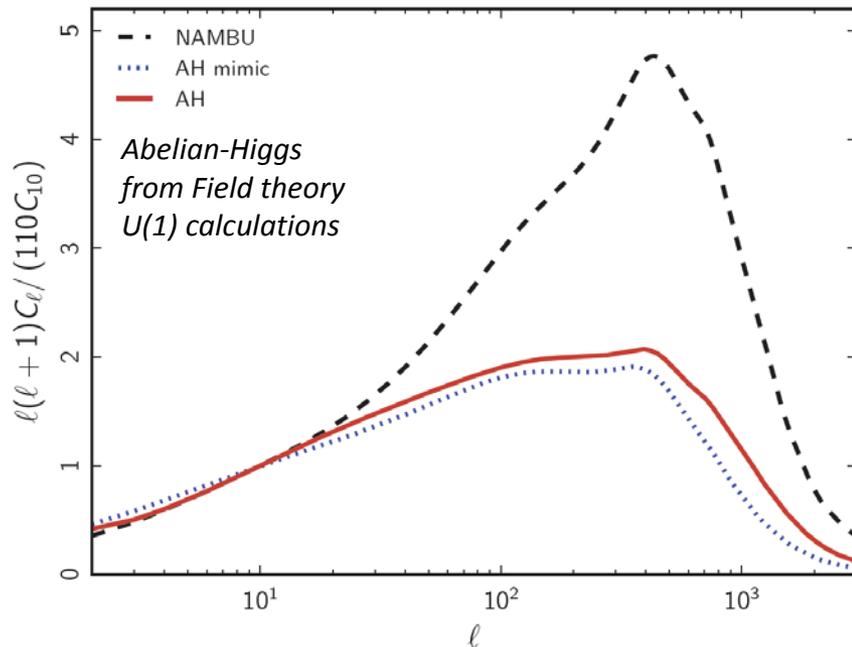
... helps (somewhat) at low- l (again!)



Cosmic defects



- Cosmic defects are a generic outcome of symmetry-breaking phase transitions in the early Universe.
 - *Strings appear in a variety of supersymmetric and other grand unified theories, forming at the end of inflation.*
 - *Cosmic (super-)strings also emerge in higher-dimensional theories for LSS, as in brane inflation.*
 - *Comparable effects can also be caused by other types of cosmic defects, notably semi-local strings and global textures.*
- The discovery of any of these objects would identify GUT-scale symmetry breaking patterns, perhaps even providing direct evidence for extra dimensions. Conversely, their absence tightly constrains symmetry breaking schemes.
- The most stringent constraints on the string tension arise from predicted backgrounds of gravitational waves from decaying loops, thus rely on most uncertain part of string physics (loop production scale and nature of string radiation from cusps)
- Defects affect both CMB power spectrum and NG properties.



Cosmic strings, normalised to equal power at $l=10$ (around $2 \cdot 10^{-6}$ with W7) \rightarrow less than 3% of overall spectrum

Comparison between global texture (black dashed) and semilocal (blue dotted) string power spectra and the AH field theory strings (red solid)

From power spectrum analysis: (other parameters, inc. n_s , unaffected)

Defect type	<i>Planck</i> +WP		<i>Planck</i> +WP+highL	
	f_{10}	$G\mu/c^2$	f_{10}	$G\mu/c^2$
NAMBU	0.015	1.5×10^{-7}	0.010	1.3×10^{-7}
AH-mimic	0.033	3.6×10^{-7}	0.034	3.7×10^{-7}
AH	0.028	3.2×10^{-7}	0.024	3.0×10^{-7}
SL Semilocal	0.043	11.0×10^{-7}	0.041	10.7×10^{-7}
TX Global texture	0.055	10.6×10^{-7}	0.054	10.5×10^{-7}

Cosmological parameters are unaffected

f_{10} = fractional contribution at $l=10$

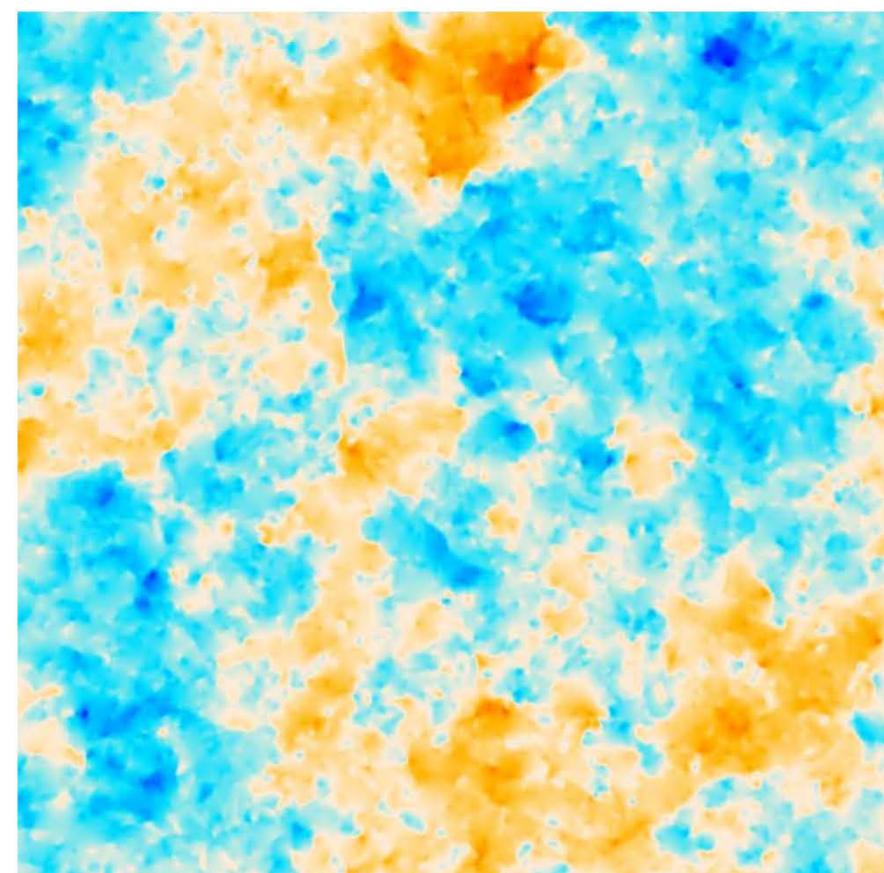
Mass per unit length, $G\mu/c^2 = (\eta/m_p)^2$, η = energy scale of symmetry breaking
(ie $\eta < 4.7 \times 10^{15}$ GeV for Nambu-Goto strings)



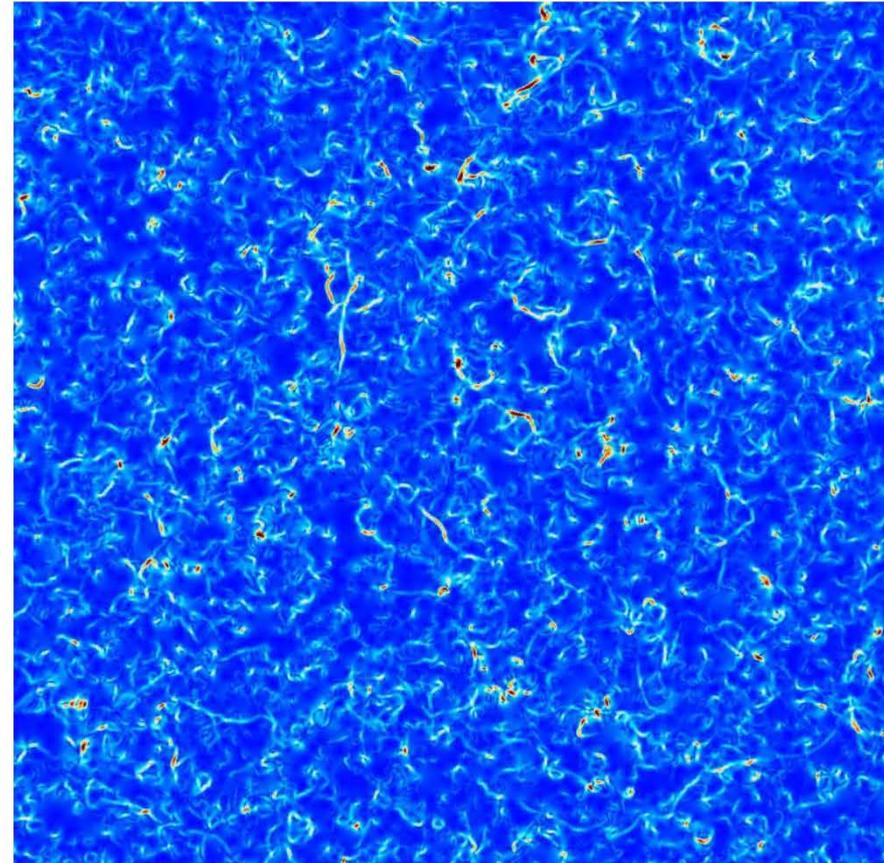
Zooming in reveals the NG aspect



A zoom on a 20 degree patch extracted from a full sky simulated makes temperature steps visible. Applying the spherical gradient magnitude operator enhances the steps, and thus the string locations.



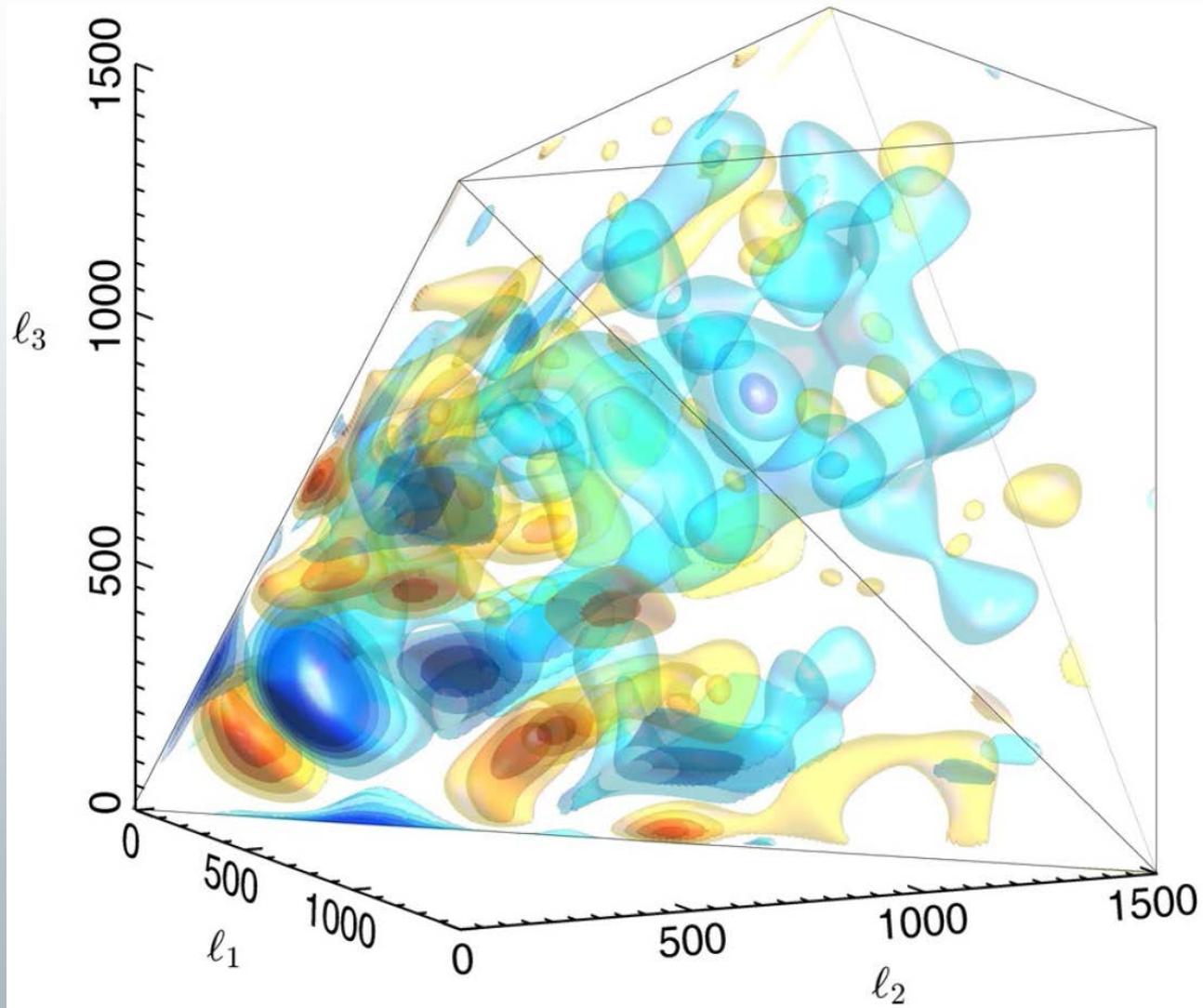
-80.0  80.0
 $\Delta T/T / (G\mu/c^2)$



0  25000
 $|\nabla(\Delta T/T)| / (G\mu/c^2)$

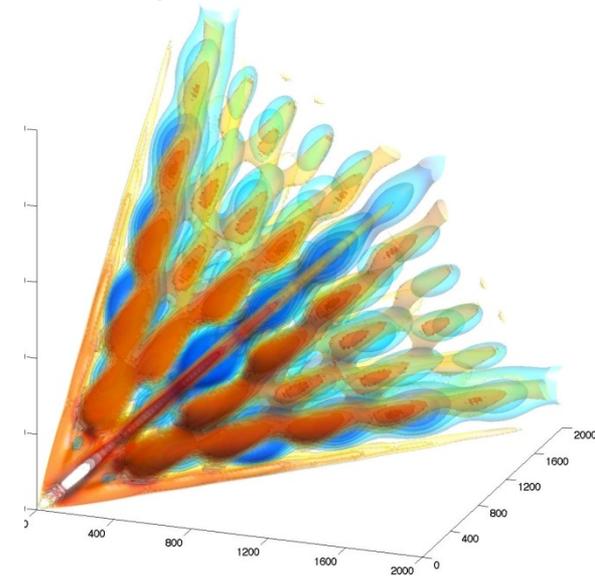
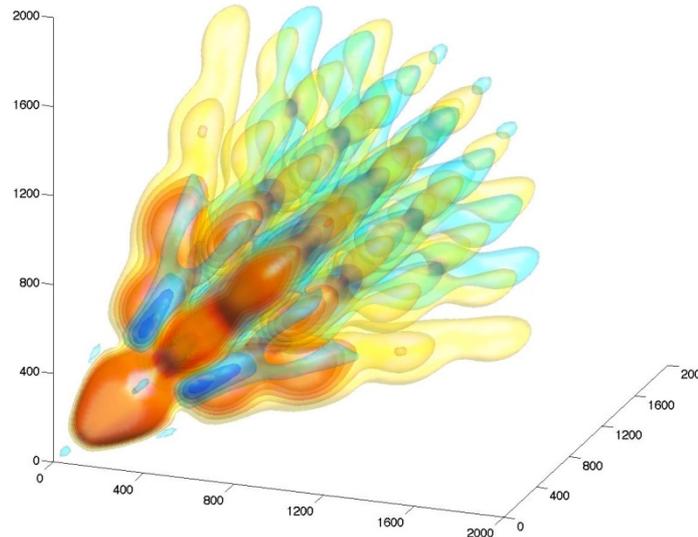
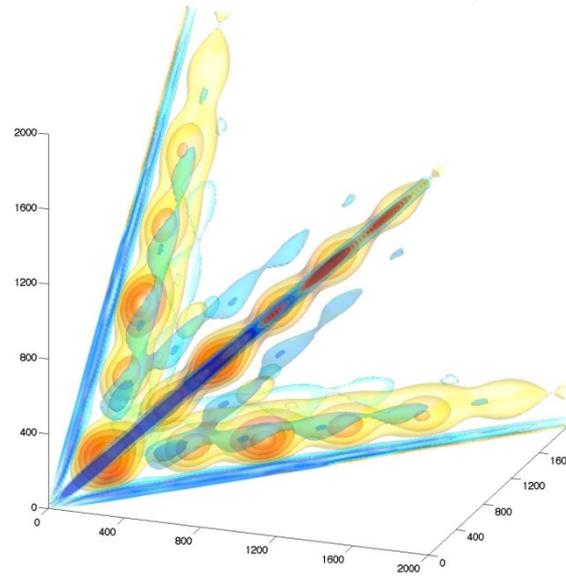
→ Modal bispectrum: $G\mu/c^2 < 8.8 \times 10^{-7}$ (95%CL); Minkowski functionals : $G\mu/c^2 < 7.8 \times 10^{-7}$ (95%CL); (post rec)

The Planck CMB bispectrum



(modal decomposition on SMICA)

LEO (local, Equilateral, Orthogonal) are common outputs



NG of **local** type:

- Multi-field models
- Curvaton
- **Ekpyrotic/cyclic models**

(Also NG of Folded type

- Non Bunch-Davis
- Higher derivative)

NG of **equilateral** type:

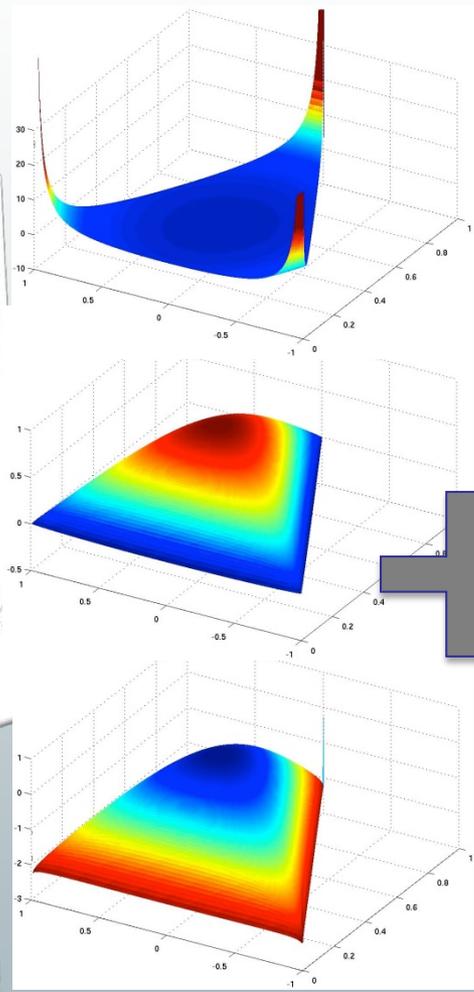
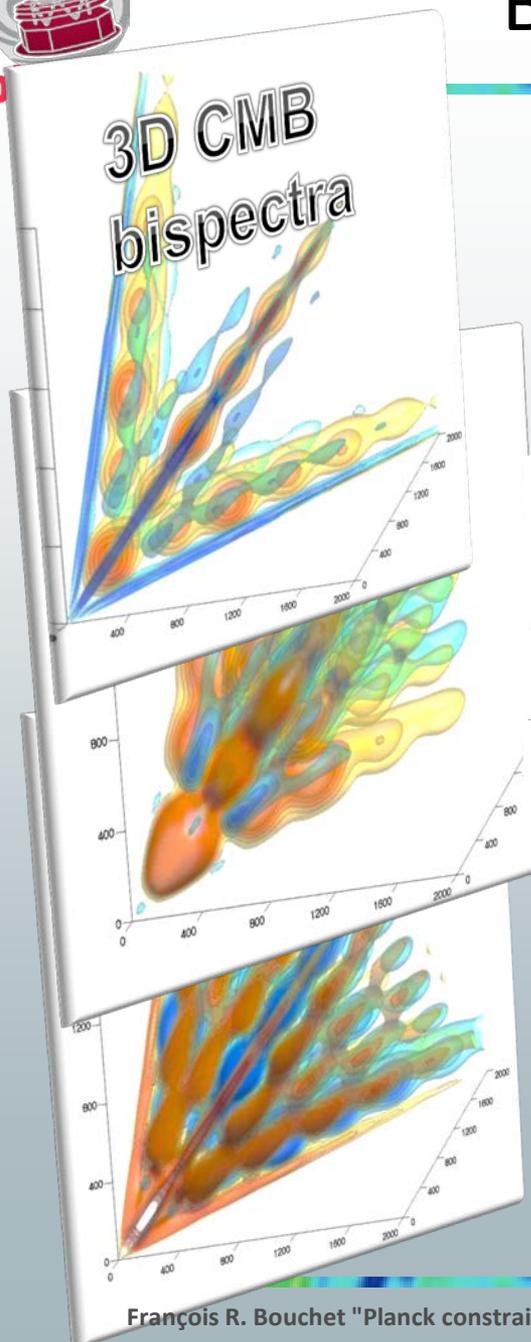
- Non-canonical kinetic term
 - K-inflation
 - DBI inflation
- Higher-derivate terms in Lagrangian
 - Ghost inflation
- Effective field theory

NG of **orthogonal** type:

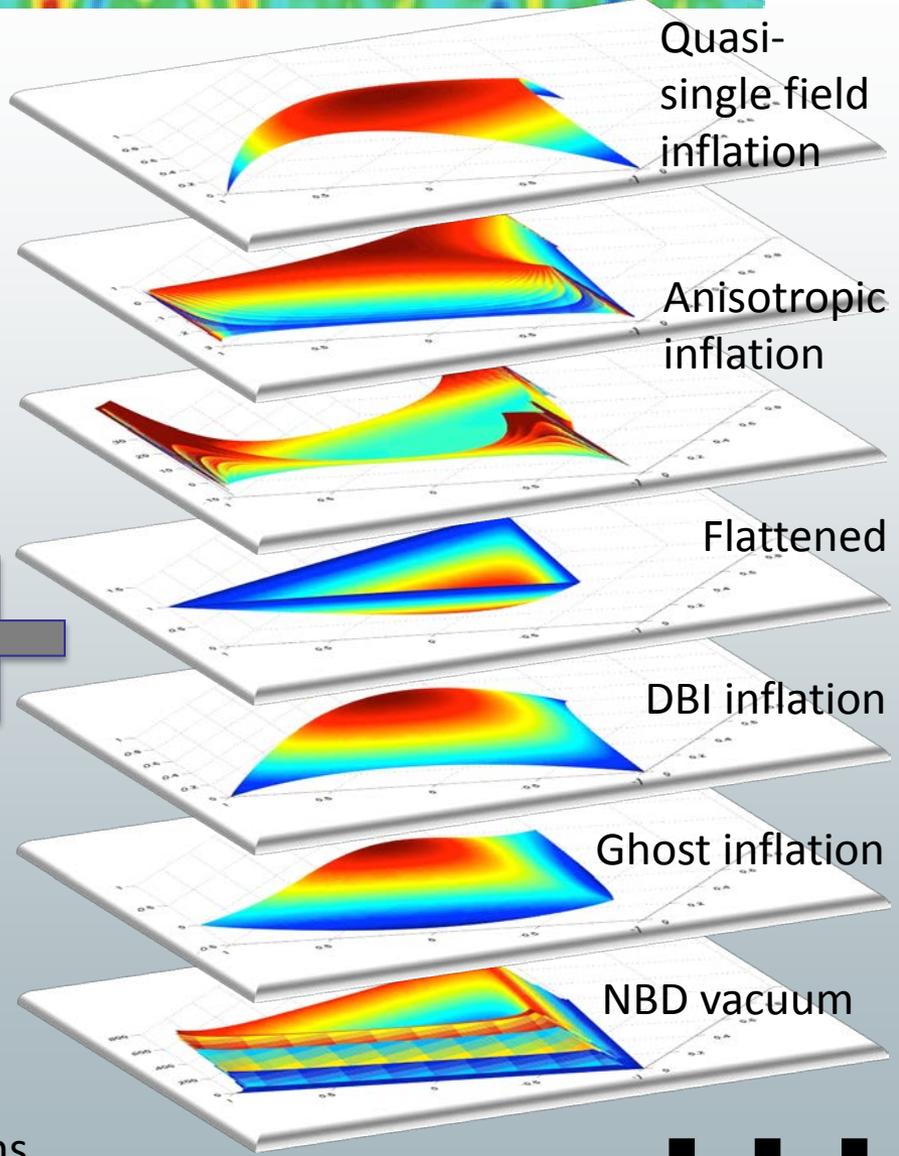
- Distinguishes between different variants of
 - Non-canonical kinetic term
 - Higher derivative interactions
- Galileon inflation



Bispectrum fingerprinting



Slices through bispectra of primordial fluctuations



Quasi-single field inflation

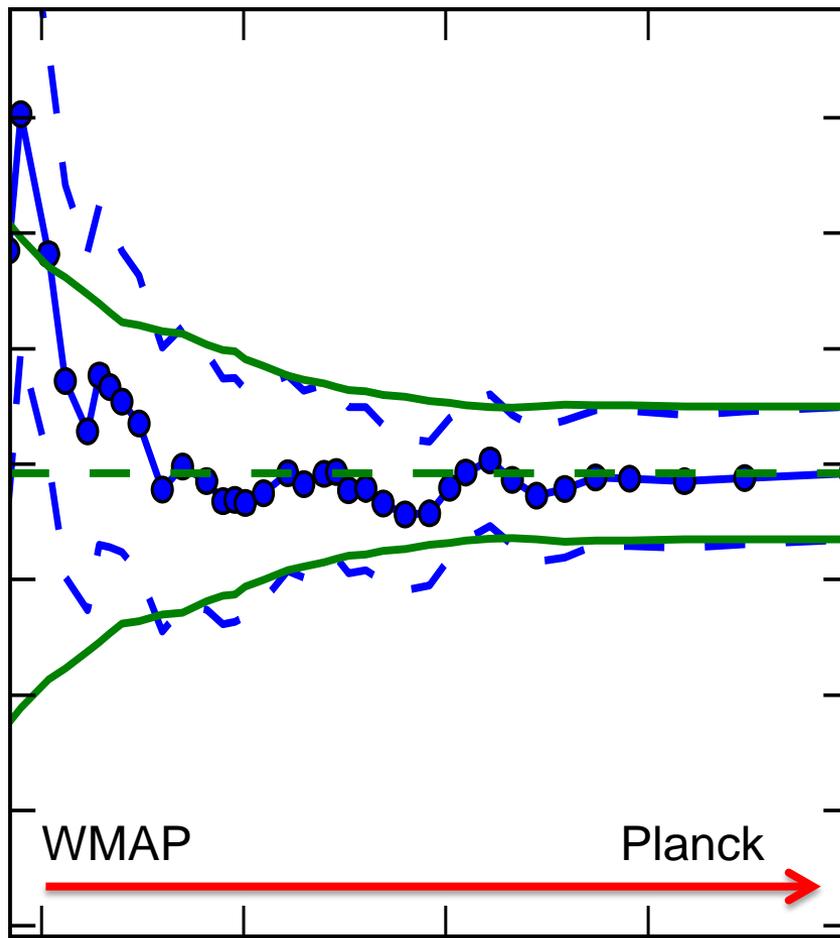
Anisotropic inflation

Flattened

DBI inflation

Ghost inflation

NBD vacuum



Limiting the analysis to large scales (low l), we make contact with WMAP9 (which gave $f_{NL}^{\text{local}} = 37.2 \pm 20$)

Planck now rules out the WMAP central value by ~ 6 sigma.

this figure is before subtraction of ISW X lensing bias, which is clearly visible



Planck results on standard LEO shapes



	Independent			ISW-lensing subtracted		
	KSW	Binned	Modal	KSW	Binned	Modal
SMICA						
Local	9.8 ± 5.8	9.2 ± 5.9	8.3 ± 5.9	2.7 ± 5.8	2.2 ± 5.9	1.6 ± 6.0
Equilateral	-37 ± 75	-20 ± 73	-20 ± 77	-42 ± 75	-25 ± 73	-20 ± 77
Orthogonal	-46 ± 39	-39 ± 41	-36 ± 41	-25 ± 39	-17 ± 41	-14 ± 42
NILC						
Local	11.6 ± 5.8	10.5 ± 5.8	9.4 ± 5.9	4.5 ± 5.8	3.6 ± 5.8	2.7 ± 6.0
Equilateral	-41 ± 76	-31 ± 73	-20 ± 76	-48 ± 76	-38 ± 73	-20 ± 78
Orthogonal	-74 ± 40	-62 ± 41	-60 ± 40	-53 ± 40	-41 ± 41	-37 ± 43
SEVEM						
Local	10.5 ± 5.9	10.1 ± 6.2	9.4 ± 6.0	3.4 ± 5.9	3.2 ± 6.2	2.6 ± 6.0
Equilateral	-32 ± 76	-21 ± 73	-13 ± 77	-36 ± 76	-25 ± 73	-13 ± 78
Orthogonal	-34 ± 40	-30 ± 42	-24 ± 42	-14 ± 40	-9 ± 42	-2 ± 42
C-R						
Local	12.4 ± 6.0	11.3 ± 5.9	10.9 ± 5.9	6.4 ± 6.0	5.5 ± 5.9	5.1 ± 5.9
Equilateral	-60 ± 79	-52 ± 74	-33 ± 78	-62 ± 79	-55 ± 74	-32 ± 78
Orthogonal	-76 ± 42	-60 ± 42	-63 ± 42	-57 ± 42	-41 ± 42	-42 ± 42

Estimators agree

- KSW (*Komatsu, Spergel, Wandelt 2003*)
- Modal (*Fergusson, Liguori, Shellard, 2009*)
- Binned (*Bucher, v. Tent, Carvalho 2009*)
- Skew-Cl (*Munshi, Heavens 2010*)
- Minkowski functionals (*Ducout et al. 2012*)

Foreground cleaning methods agree

- SMICA
- NILC
- SEVEM
- (C-R)

Null tests pass

Consistency test pass

- Resolution dependence
- Frequency dependence
- Mask dependence

Negligible impact of FG residuals

→ To no avail for flattened shapes...

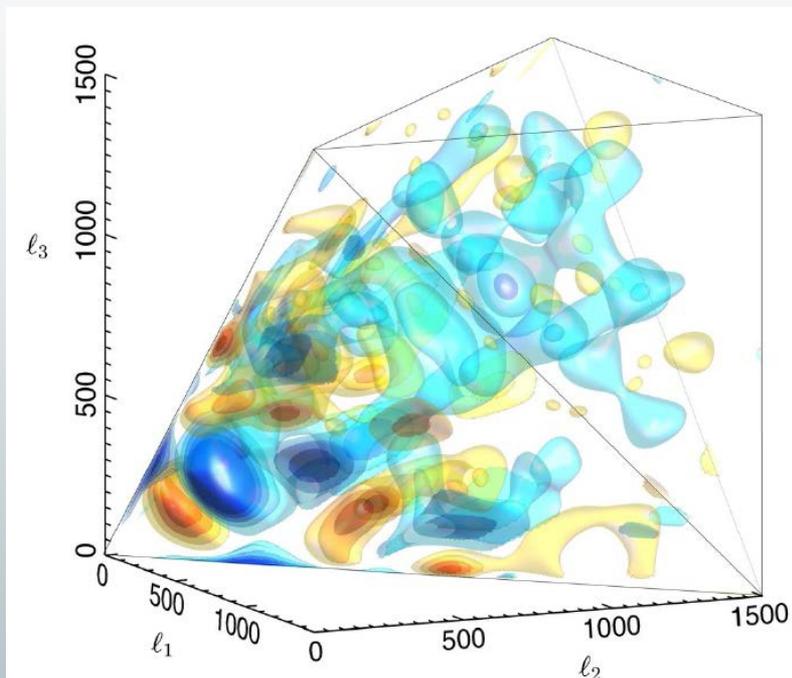


Flattened model (Eq. number)	Raw f_{NL}	Clean f_{NL}	Δf_{NL}	σ	Clean σ
Flat model (13)	70	37	77	0.9	0.5
Non-Bunch-Davies (NBD)	178	155	78	2.2	2.0
Single-field NBD1 flattened (14)	31	19	13	2.4	1.4
Single-field NBD2 squeezed (14)	0.8	0.2	0.4	1.8	0.5
Non-canonical NBD3 (15)	13	9.6	9.7	1.3	1.0
Vector model $L = 1$ (19)	-18	-4.6	47	-0.4	-0.1
Vector model $L = 2$ (19)	2.8	-0.4	2.9	1.0	-0.1

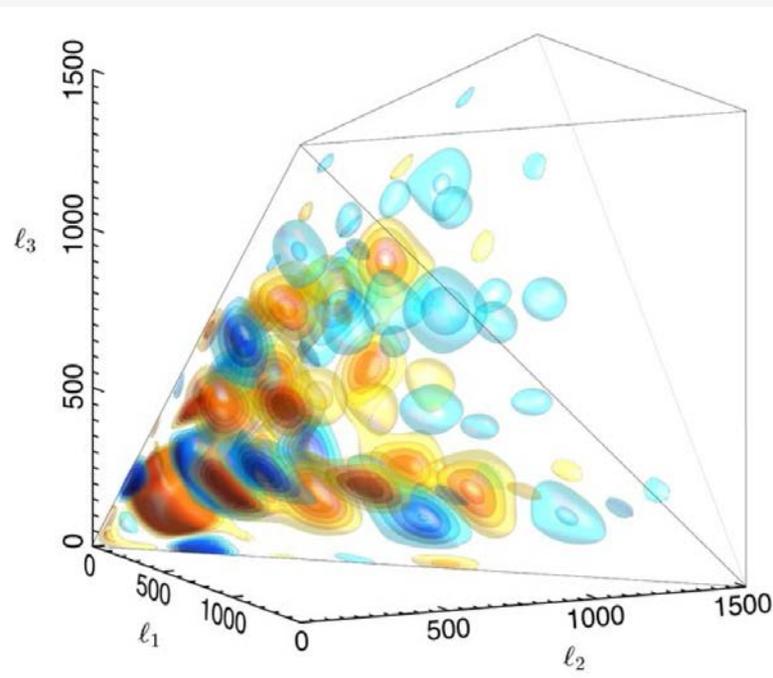
→ And with tantalising hints for some feature models...

Width	$\Delta k = 0.015$	$\Delta k = 0.03$	$\Delta k = 0.045$	Full
Model	$f_{\text{NL}} \pm \Delta f_{\text{NL}} (\sigma)$			
$k_c = 0.01125; \phi = 0$	$765 \pm 275 (2.8)$	$703 \pm 241 (2.9)$	$648 \pm 218 (3.0)$	$434 \pm 170 (2.6)$
$k_c = 0.01750; \phi = 0$	$-661 \pm 234 (-2.8)$	$-494 \pm 192 (-2.6)$	$-425 \pm 171 (-2.5)$	$-335 \pm 137 (-2.4)$
$k_c = 0.01750; \phi = 3\pi/4$	$399 \pm 207 (1.9)$	$438 \pm 183 (2.4)$	$442 \pm 165 (2.7)$	$366 \pm 126 (2.9)$
$k_c = 0.01875; \phi = 0$	$-562 \pm 211 (-2.7)$	$-559 \pm 180 (-3.1)$	$-515 \pm 159 (-3.2)$	$-348 \pm 118 (-3.0)$
$k_c = 0.01875; \phi = \pi/4$	$-646 \pm 240 (-2.7)$	$-525 \pm 189 (-2.8)$	$-468 \pm 164 (-2.9)$	$-323 \pm 120 (-2.7)$
$k_c = 0.02000; \phi = \pi/4$	$-665 \pm 229 (-2.9)$	$-593 \pm 185 (-3.2)$	$-500 \pm 160 (-3.1)$	$-298 \pm 119 (-2.5)$

Data



Best fit feature model bispectrum



Not (yet?) highly significant if “look-elsewhere” effect is taken into account.

➔ Warrants further analysis.



Cases of joint $C(l)$ - f_{NL} constraints



- **Case 1:** $c_s = \text{constant}$

NG $\rightarrow c_s > 0.02$ (95%CL)

This reduce degeneracies

\rightarrow one get the constraint (adjacent) between the 1st two Hubble Flow Functions

- **Case 2:** IR-DBI

$V = V_0 - \frac{1}{2} \beta H^2 \phi^2, .1 < \beta < 10^9$

Planck constraint on $n_s + f_{NL}^{DBI}$

$\rightarrow \beta < 0.7$ (95%CL)

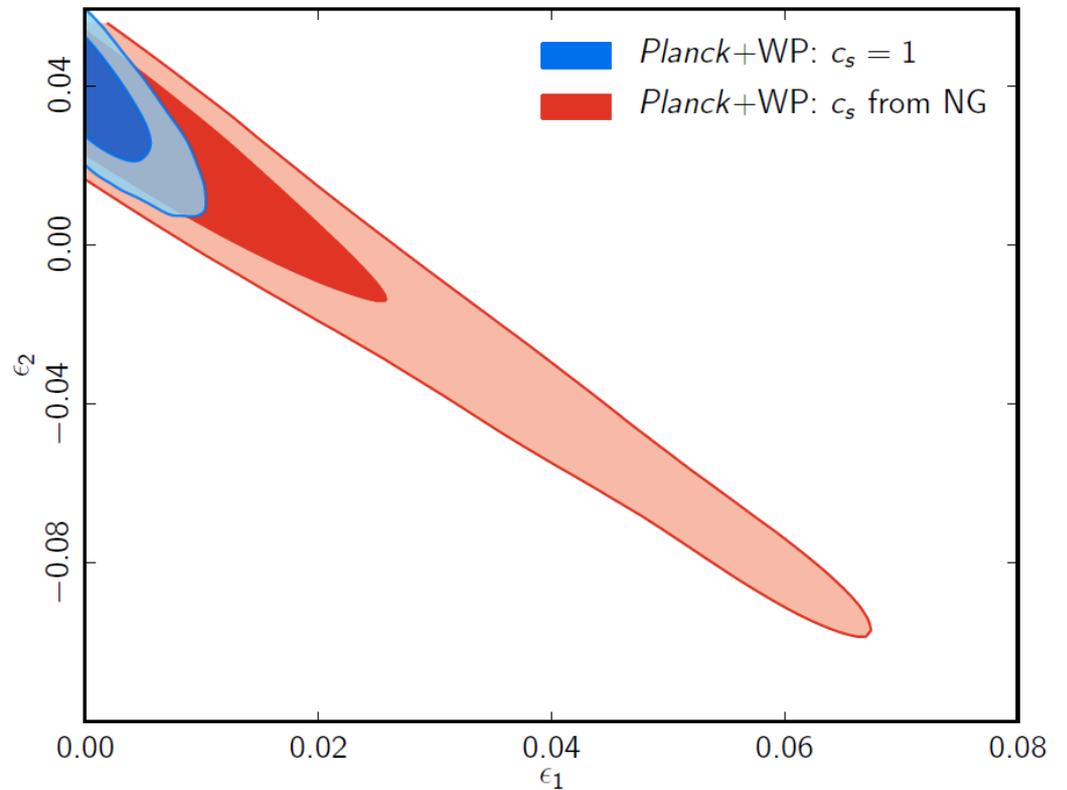
- **Case 3:** k-inflation

One class depends on single param. γ (Amendariz-Picon et al 99). Planck:

$f_{NL}^{equi} \rightarrow \gamma > 0.05$ (95%CL)

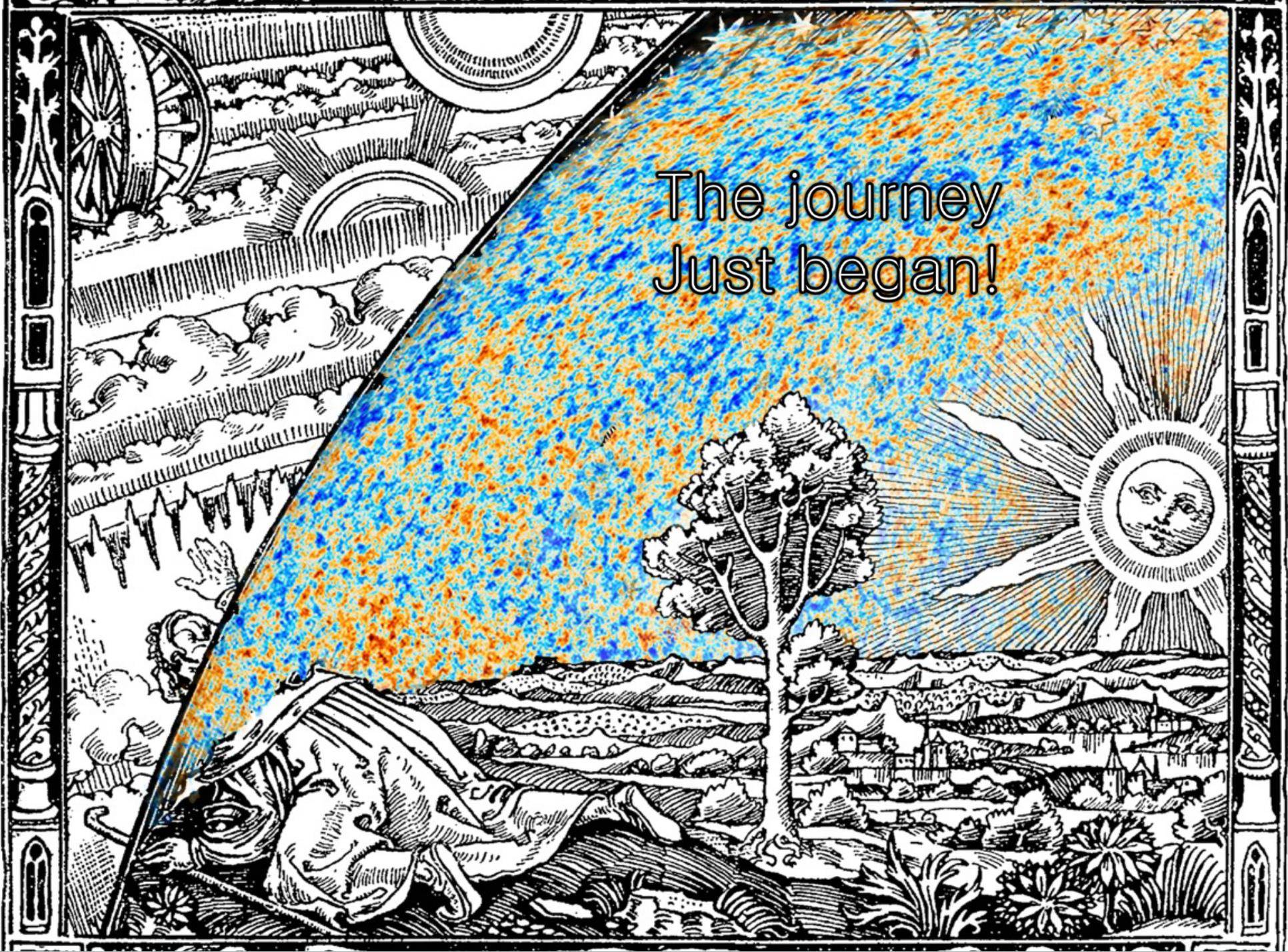
$n_s \rightarrow 0.01 < \gamma < 0.02$ 95%CL !

- **Case 4:** ekpyrotic alternative to inf. « conversion mechanism » decisely ruled out. Viable parameter space for « kinetic conversion » dramatically restricted



NB: the absence of isocurvature modes (the curvaton fully correlated case $n_{II} = n_R$) yields the stronger constraints $R_D > 0.98$ than $r_D > 0.15$ from the f_{NL}^{local} limit

The journey
Just began!





A full cryogenic life



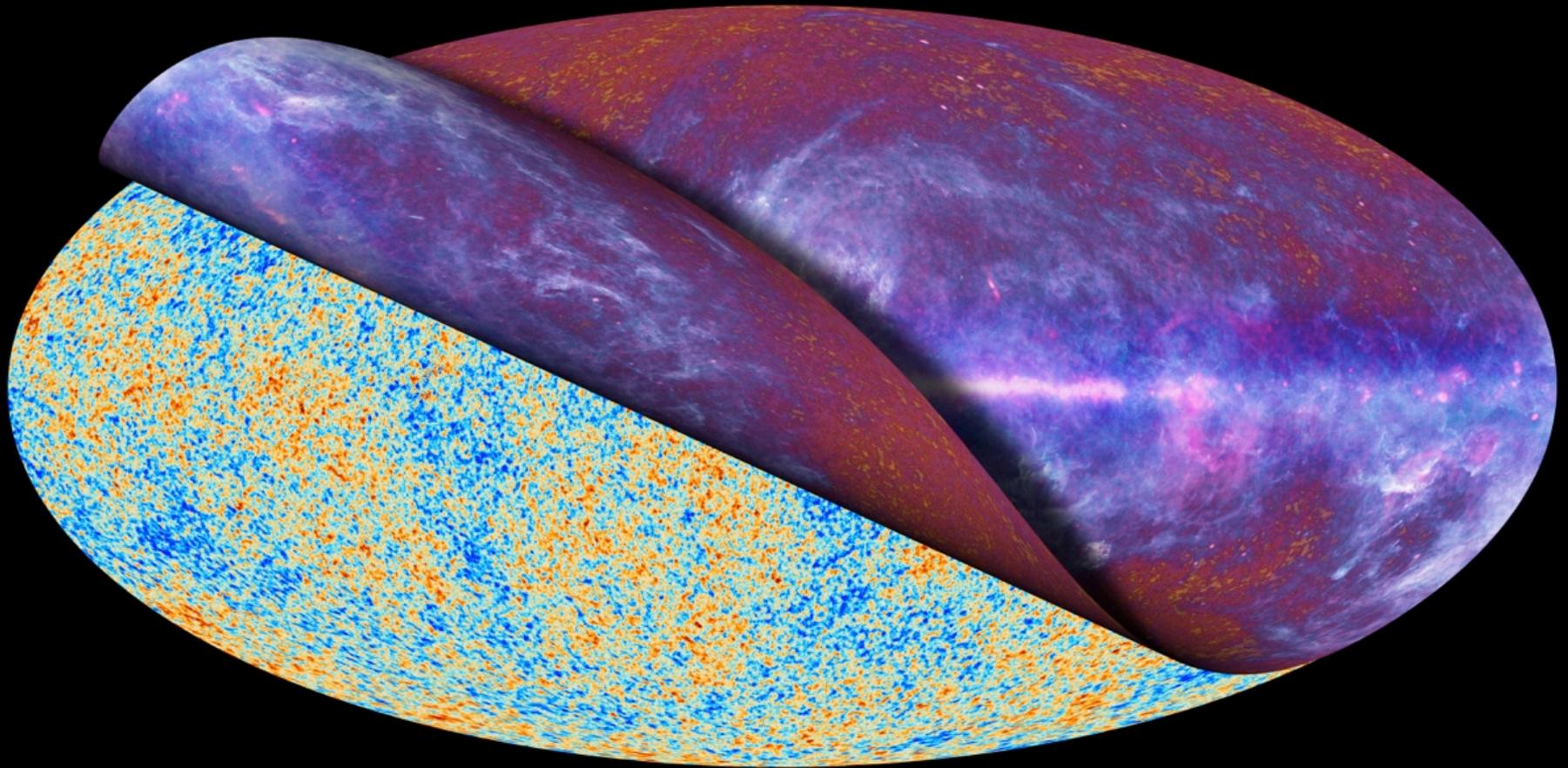
HFI starts cooling,
Hotel des roches,
Kourou
14 mai 2009

HFI starts warming up
HFI Core team meeting
Institut d'Astrophysique de Paris,
14 janvier 2012





planck



Planck unveils the Cosmic Microwave Background