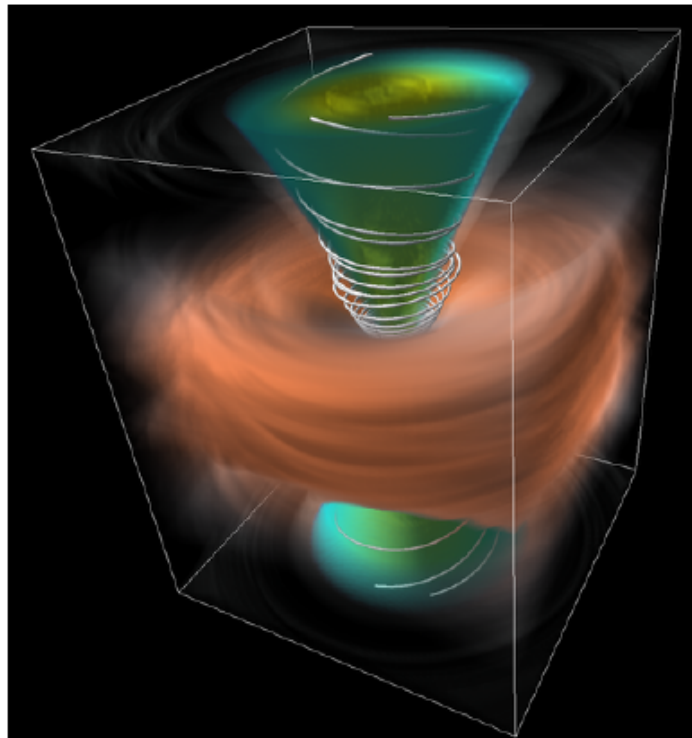


Disk modelling by global radiation-MHD simulations

~Confrontation of inflow & outflow~

Shin Mineshige (Kyoto) & Ken Ohsuga (NAOJ)



Magnetic tower jet
by RMHD simulation
(Takeuchi+11)

Outline

- Introduction

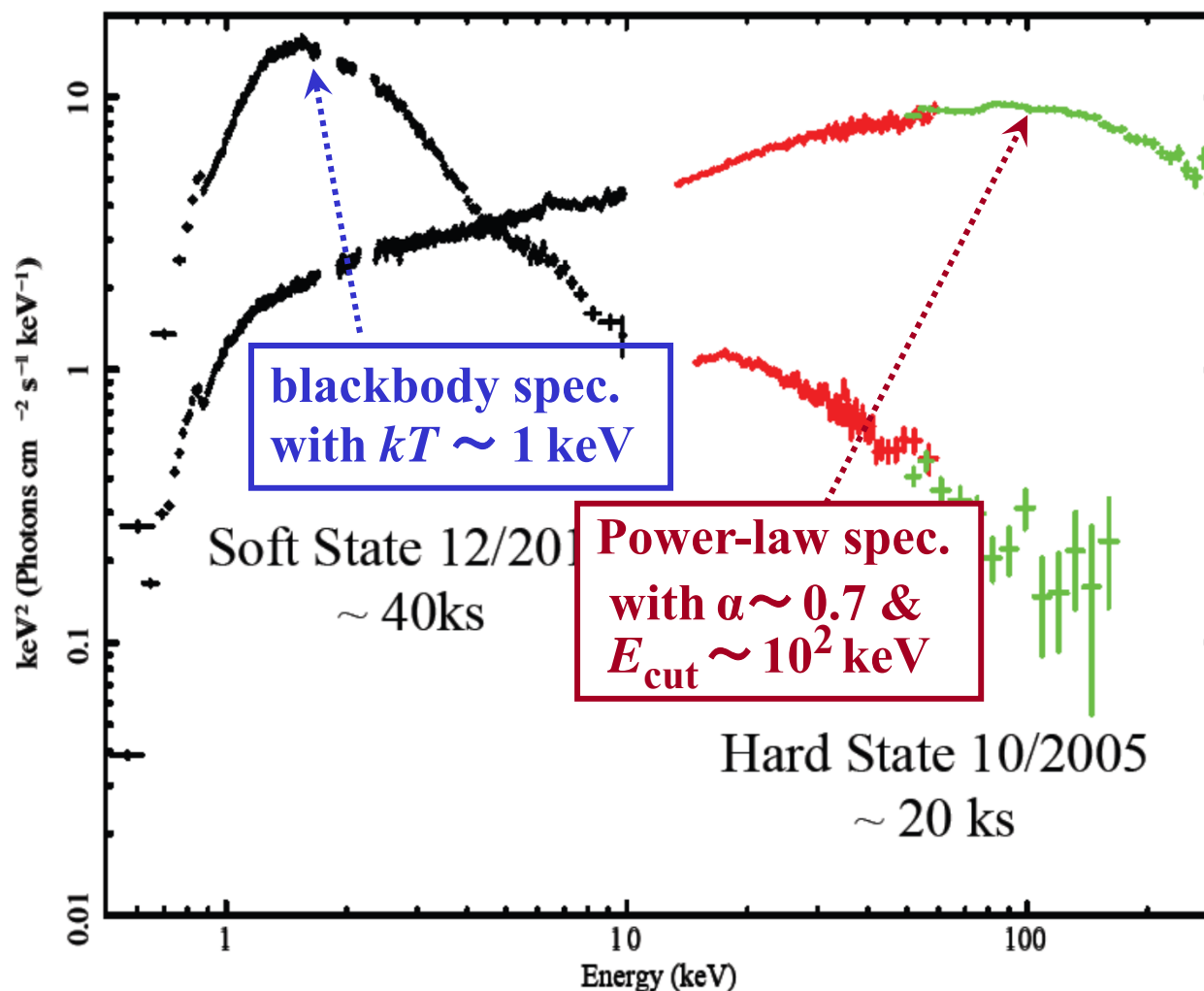
- Various states and/or spectral components
- Basic theory of black hole accretion

- Global R-MHD simulations of super-critical flow

- Motivation: why global radiation-MHD simulations?
- Three distinct regimes of accretion flow
- P_{rad} -driven jet & clumpy outflow
- Comptonization (?)

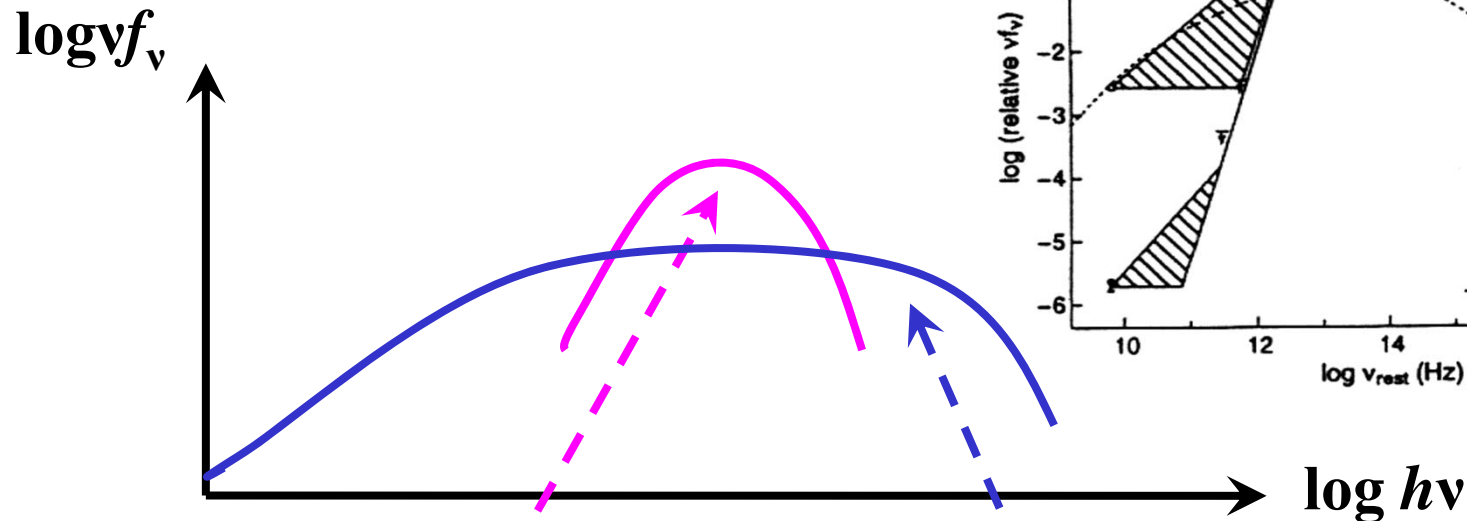
BH binaries: two (basic) spec. states

(S. Yamada/
Suzaku data)



\Rightarrow Combination of **Hot flow** ($T \sim 10^9\text{ K}$) & **Cool disk** ($\sim 10^7\text{ K}$)

AGN: two spectral components



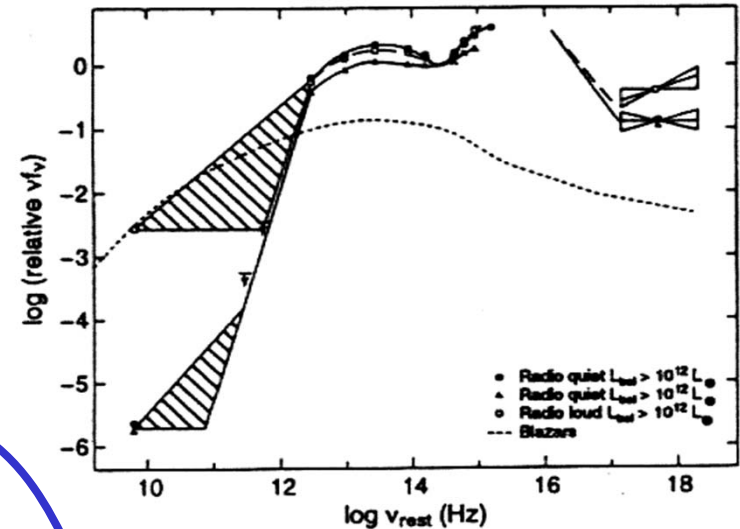
Big Blue Bump (UV)

blackbody with

$T_{\text{eff}} \sim 10^5 \text{ K}$ (10 eV)

broad-band comp. (radio $\sim \gamma$)

$f_\nu \propto \nu^{-\alpha}$ with $\alpha \sim 0.7$, cutoff
with $T_{\text{elec}} \sim 10^9 \text{ K}$ + self-abs.



Sanders + 89

\Rightarrow coexistence of hot ($\sim 10^9 \text{ K}$) and 'cool' ($\sim 10^5 \text{ K}$)
material (indicated by Fe fluorescence line)

Two basic solutions!

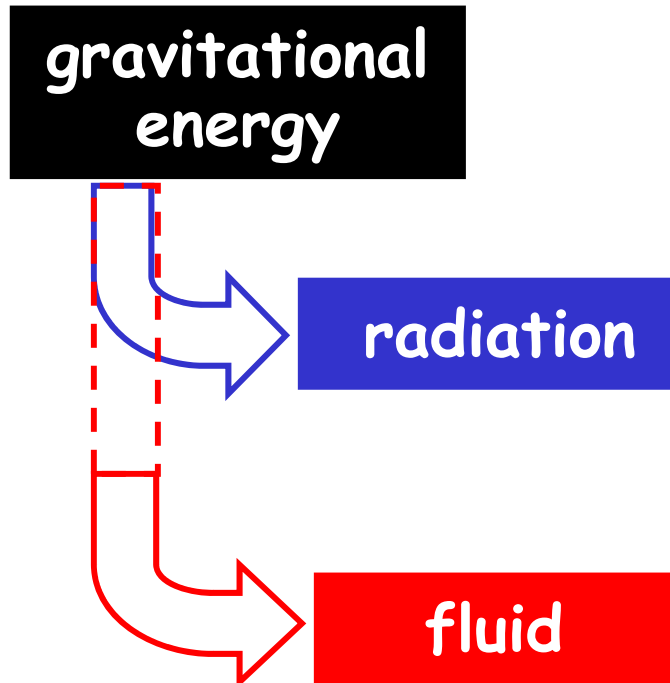
Kato, Fukue & Mineshige (2008)

Energy equation: $Q_{\text{adv}} = Q_{\text{vis}} - Q_{\text{rad}}$

$Q_{\text{adv}} \sim \Sigma T v_r (ds/dr) = \text{advection term}$

$Q_{\text{vis}} = \text{viscous heating}$

$Q_{\text{rad}} = \text{radiative cooling}$



standard disk: $Q_{\text{vis}} = Q_{\text{rad}} \gg Q_{\text{adv}}$

→ cool bright disk

→ $T \sim 10^7 \text{ K } (M_{\text{BH}}/M_{\text{sun}})^{-1/4}$

RIAF: $Q_{\text{vis}} = Q_{\text{adv}} \gg Q_{\text{rad}}$

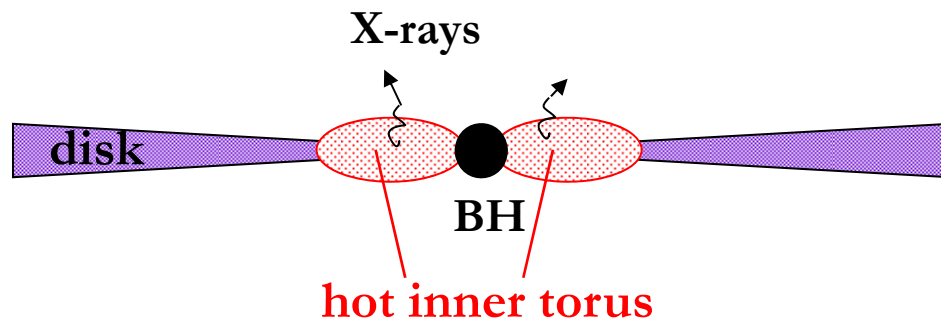
→ hot faint flow → $T \sim 10^9 \text{ K}$

(RIAF=Radiatively Inefficient Accretion Flow)

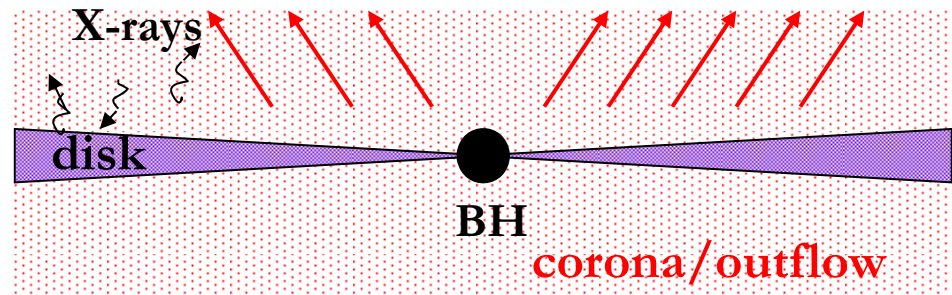
Situations may not be so simple...

Cool disks and hot corona/outflow

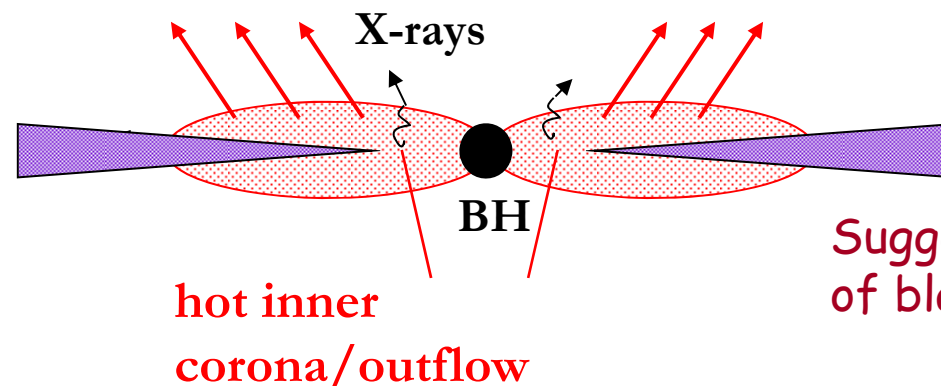
(a) horizontal separation



(b) vertical separation



(c) mixture type



Suggested by X-ray obs.
of black hole binaries.

Why global radiation-MHD simulations?

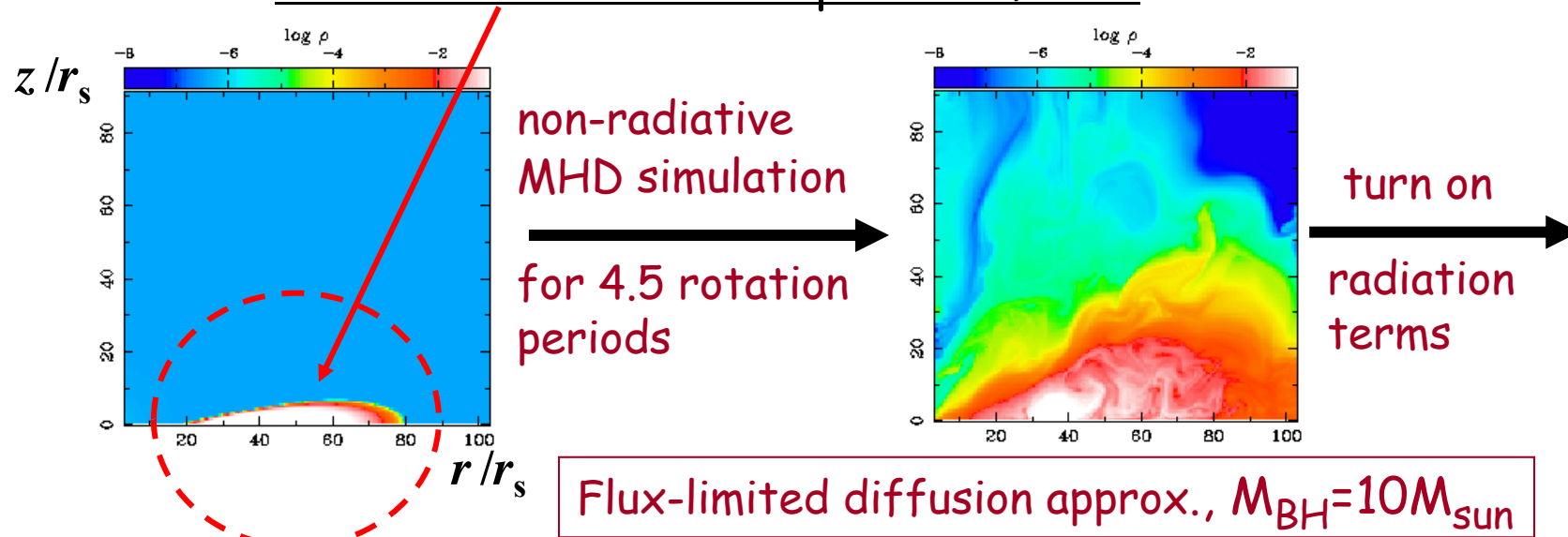
One-dimensional models (e.g. standard disk, slim disk, ADAF) are quite useful for understanding the basics of different modes of accretion. **However...**

- Multi-dimensional gas motions, such as circulation, convection, outflow... are totally neglected.
→ **Need 2D/3D simulations**
- Disk viscosity was treated by the phenomenological α model. Its validity needs to be examined. →
Need MHD simulations
- Strong radiation-matter interactions expected at high luminosities are not properly treated.
→ **Need radiation-hydrodynamical (RHD) simulations**

Our global 2D RMHD simulations

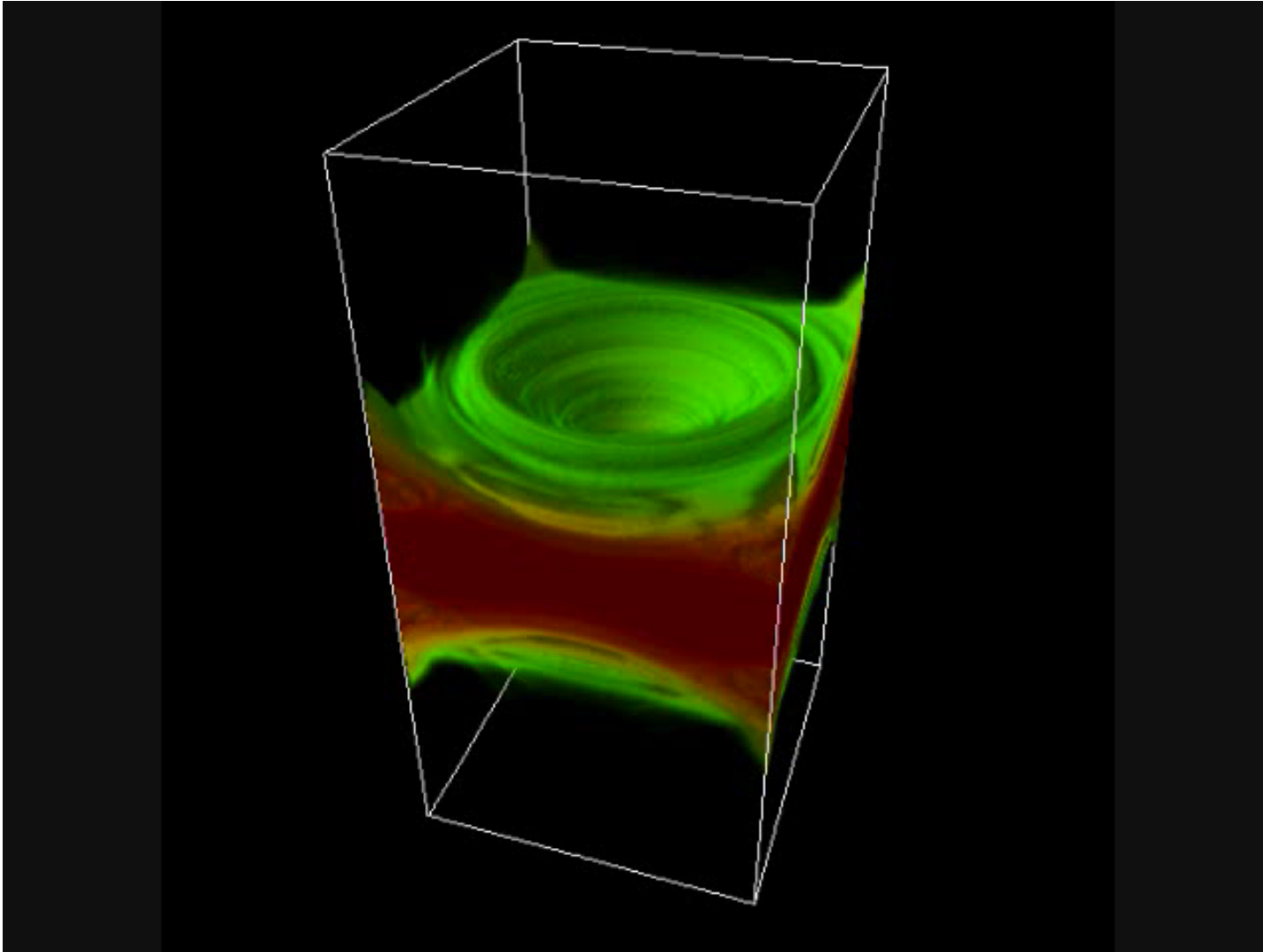
Ohsuga, SM, et al. (2009, PASJ 61, L7)

- Start with a torus threaded weak poloidal fields:



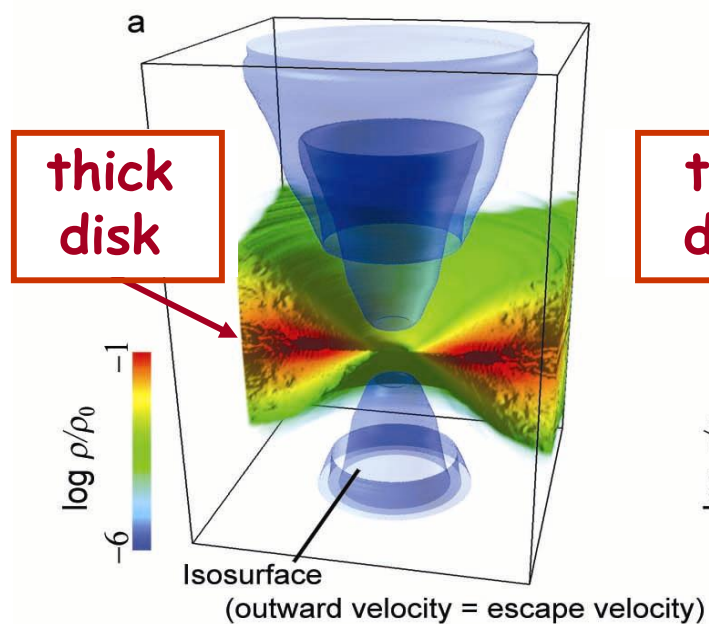
- Results depend on the initial density normalization (ρ_0)
 - Model A ($\rho_0=10^0 \text{ g/cm}^3$) : supercritical state (P_{rad} force)
 - Model B ($\rho_0=10^{-4} \text{ g/cm}^3$) : standard-disk state (rad. cooling)
 - Model C ($\rho_0=10^{-8} \text{ g/cm}^3$) : radiatively inefficient flow state
- ⇒ Three distinct states can be modelled by one code.

Radiation-MHD simulation: model A



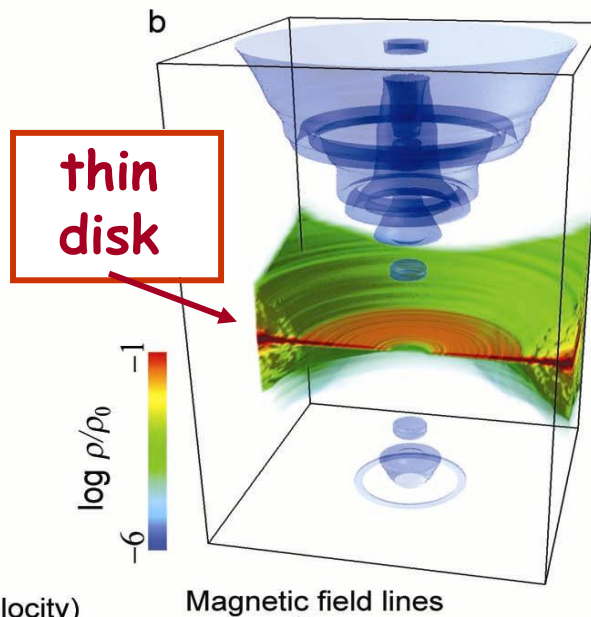
supercritical (high L)

Model A



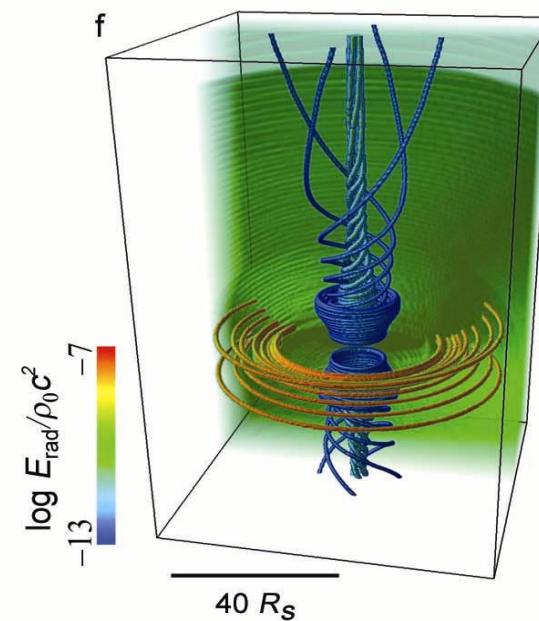
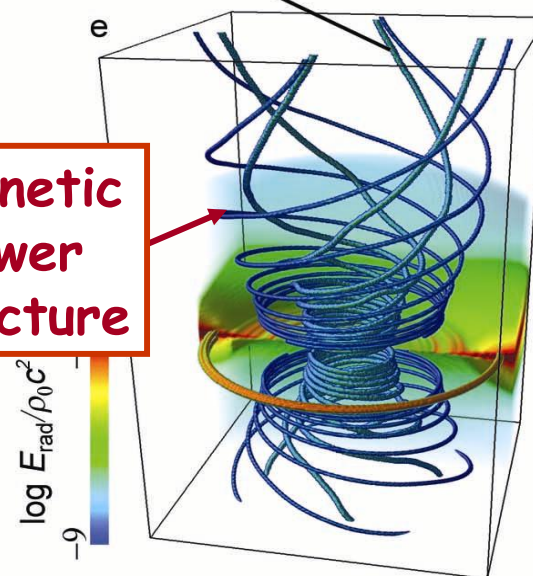
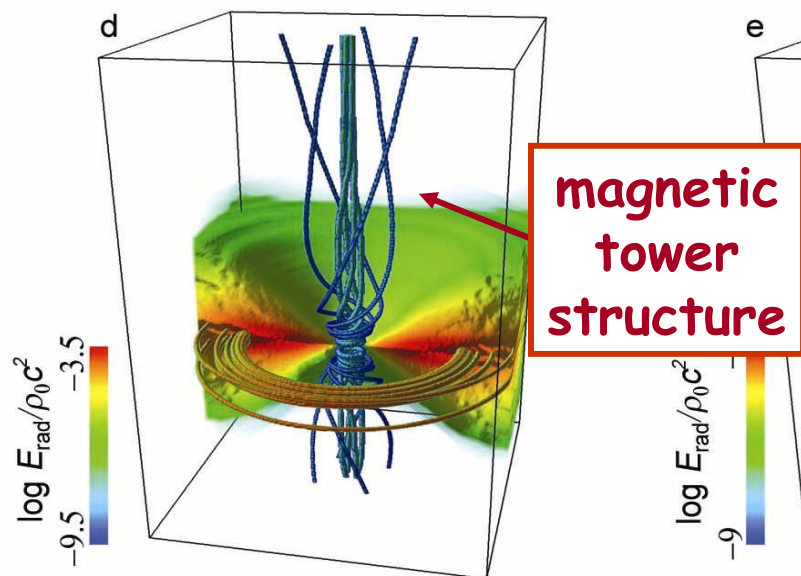
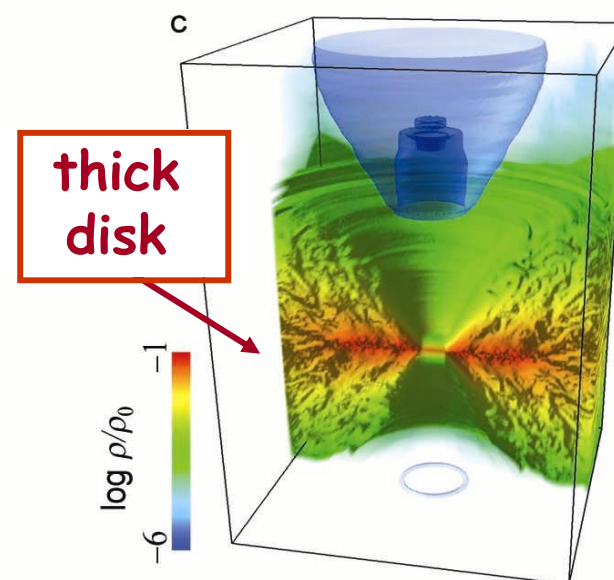
standard disk

Model B



RIAF (low L)

Model C



Summary of radiation-MHD simulations

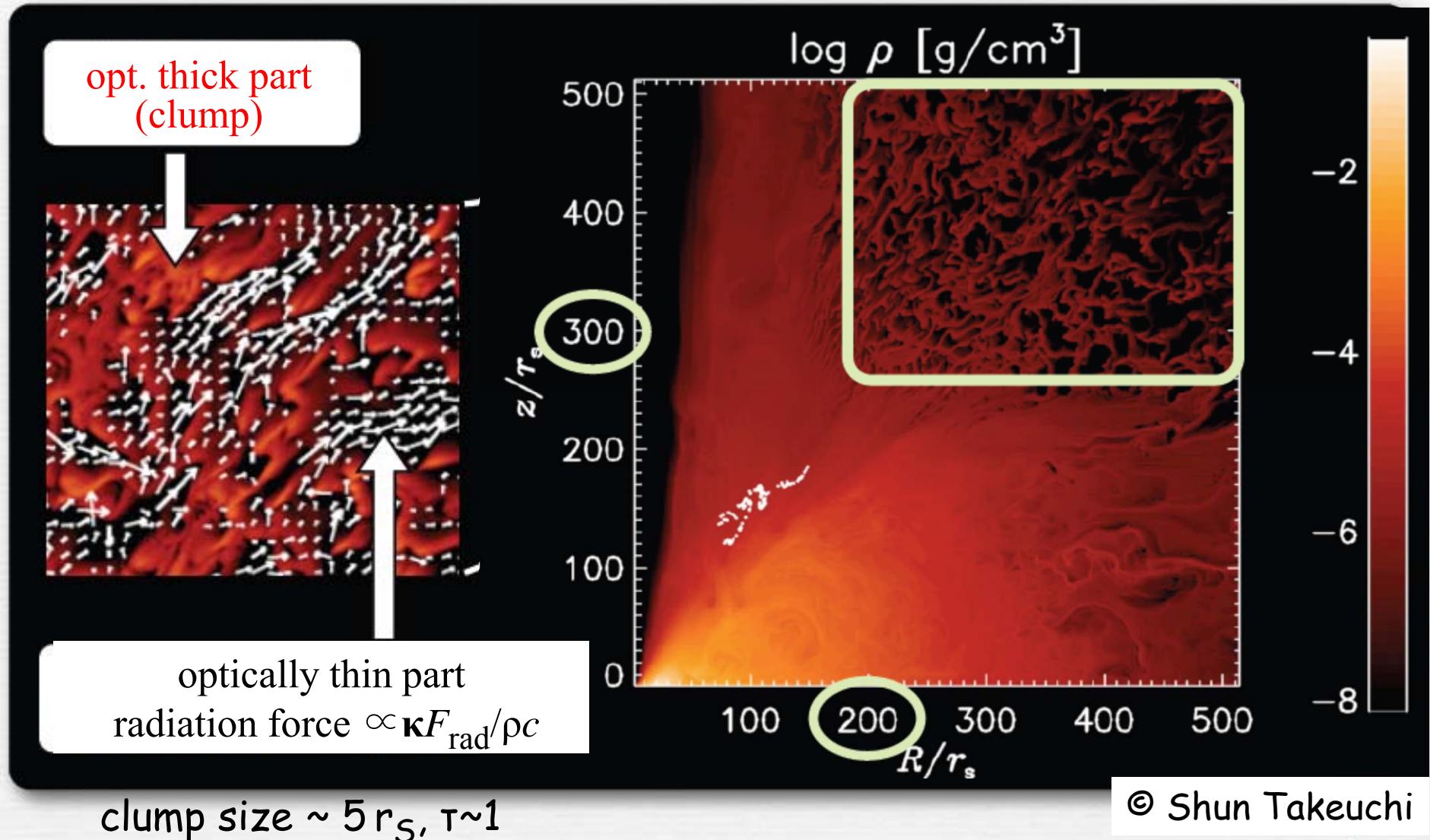
Ohsuga & SM (2011, ApJ 237)

Model	density & temperature	luminosity, L/L_E	Energetics	kin. luminosity, L_{kin}/L
Model A (supercritical)	$\rho \sim 10^{-2} \text{ g/cm}^3$ $T \sim 10^8 \text{ K}$	$\sim 10^0$	$E_{\text{rad}} > E_{\text{mag}} > E_{\text{gas}}$	~ 1
Model B (standard)	$\rho \sim 10^{-5} \text{ g/cm}^3$ $T \sim 10^6 \text{ K}$	$\sim 10^{-2}$	$E_{\text{gas}} \gtrsim E_{\text{mag}} \sim E_{\text{rad}}$	~ 0.003
Model C (RIAF)	$\rho \sim 10^{-9} \text{ g/cm}^3$ $T \sim 10^{10} \text{ K}$	$\sim 10^{-8}$	$E_{\text{gas}} > E_{\text{mag}} \gg E_{\text{rad}}$	~ 3

- ⇒ **Model A (supercritical): thick disk with P_{rad} -driven jet & outflow**
Model B (standard) : thin disk with weak outflow
Model C (RIAF) : thick disk with P_{mag} -driven jet & outflow

Remark 1. Clumpy outflows \rightarrow variability

(Outflow speed $\lesssim 0.1c$, $\dot{M}(\text{outflow}) \sim L_E/c^2$, opening angle = 20-50 deg)



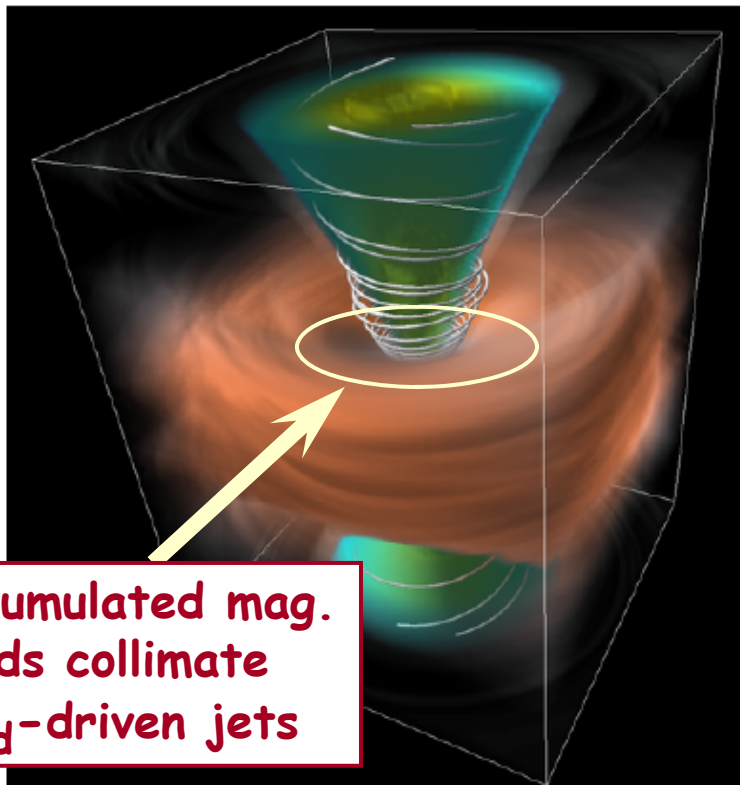
Remark 2. Radiation-MHD jets

(jet speed $\sim 1c$, opening angle $< 10^\circ$)

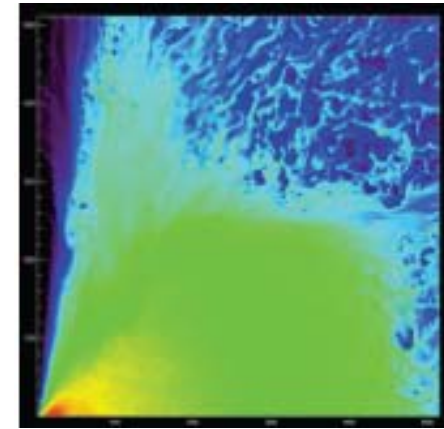
(Takeuchi+11)

Powerful P_{rad} -driven magnetic tower jets are produced at high luminosities ($L \gtrsim L_E$).

Bird's eye view \downarrow



Accumulated mag. fields collimate P_{rad} -driven jets



jet ($v \lesssim c$)

outflow ($v \sim 0.1c$)

clumpy outflow

optically thick disk

500 R_s

Remark 3. Comptonization (?)

(RHD simulations by Kawashima et al. 2009, 2012)

Compton cooling

γ -parameter > 1

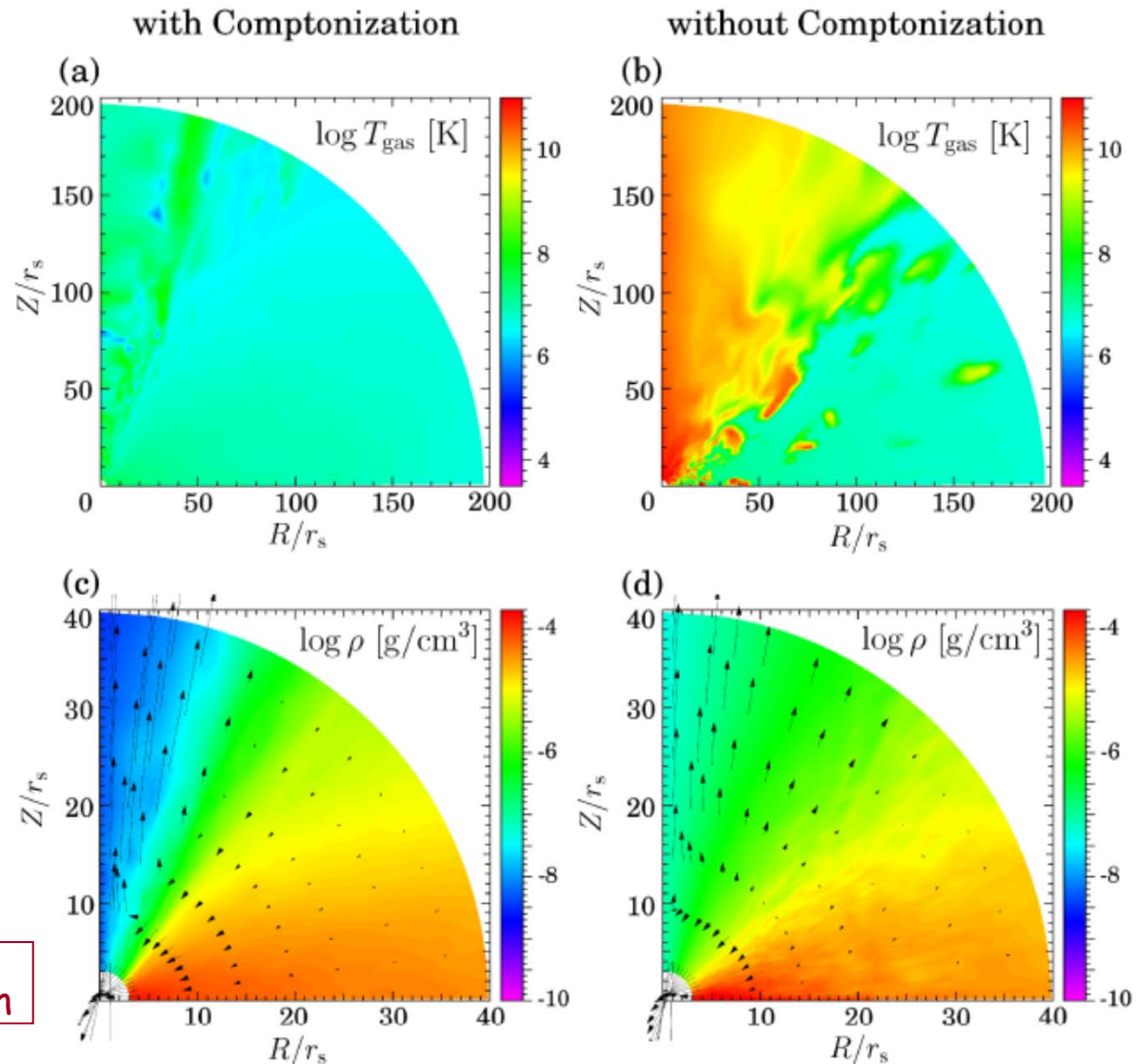
→ significant cooling

→ also bulk Compton

Should be important
for BHBs (& ULXs).

How about the cases
with supermassive
BHs?

$$M_{\text{BH}} = 10 M_{\text{sun}}$$



Summary

- Variety of accretion modes are observed in BH objects. The flow is composed of hot ($\sim 10^9\text{K}$) plasmas and cool disk (10^{5-7}K) but their locations are unknown.
- For making a unified view of various accretion modes, we are performing global radiation-MHD simulations.
- By controlling a density normalization we could for the first time reproduce three distinct modes of accretion flow and outflow by one code.
- Outflow is quite ubiquitous. At high luminosities ($\gtrsim L_E$), in particular, powerful P_{rad} -driven jets and clumpy outflows are produced. Comptonization may be important.